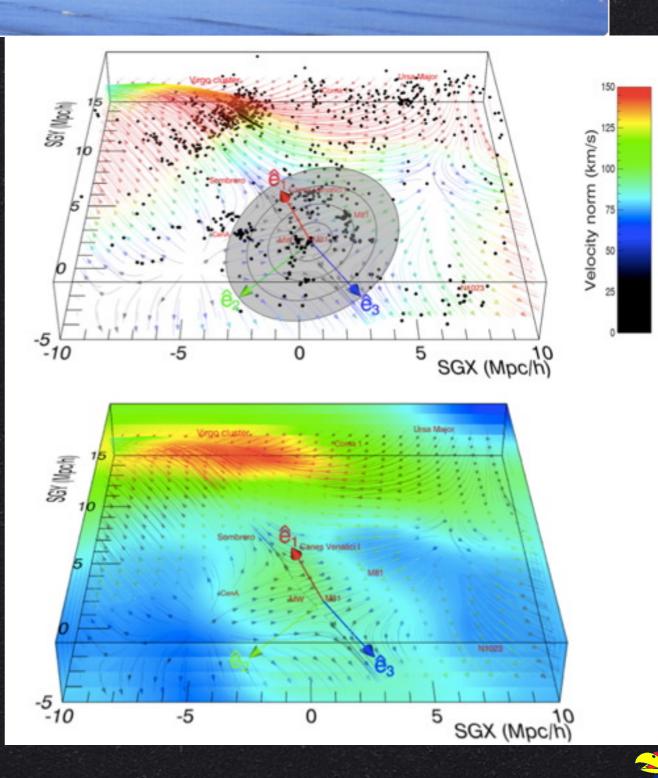
LARGE SCALE STRUCTURE AND GALAXY FLOWS July 3rd - 9th, 2016 • Quy Nhon, Vietnam

velacity correlation Problems and Opportunities

Hume A. Feldman

Physics & Astronomy University of Kansas



Likelihood Methods for Peculiar Velocities

A catalog of peculiar velocities galaxies, labeled by an index n

Positions r_n

Line-of-sight peculiar velocity estimators S_n

Uncertainties σ_n

Assume that observational errors are Gaussian distributed.



Estimators

Determination of the line-of-site peculiar (local) motion requires the measurement of the galaxy's distance

 $v = cz - H_o r$

At large distances, we can include the effects of cosmic acceleration

 $z_{\rm mod} = z[1+0.5(1-q_o)z - (1/6)(1-q_o-3q_o^2+1)z^2]$

Since redshift is not an additive quantity

See also Davis & Scrimgeour 2014; Springob et al. 2014

 $(1 + z_{\text{mod}}) = (1 + H_o r/c)(1 + v/c)$

The peculiar velocity is

 $v = \frac{cz_{\text{mod}} - H_o r}{1 + H_o r/c} \approx \frac{cz_{\text{mod}} - H_o r}{1 + z_{\text{mod}}}$

Estimators

Determination of the line-of-site peculiar (local) motion requires the measurement of the galaxy's distance

Distance estimators give the distance moduli (µ), that is, log distances with Gaussian distributed errors.

Distance errors are skewed, not Gaussian and with non-zero average.

$$\langle r_e \rangle \neq r \quad \Rightarrow \quad \langle v_e \rangle \neq v$$

These undesirable features can lead to biases and invalidate our statistical assumptions about the errors in peculiar velocities.



Estimators Proposal: estimate peculiar velocities using $v_e = cz \log(cz/H_o r_e)$

Watkins & Feldman, 2015, MNRAS 450, 1868

log distance \Rightarrow Gaussian distributed errors.

The uncertainty in the peculiar velocity

 $\delta v_e = cz \delta l_e$

. Uncertainty in the log distance

and

 $\langle \log(r_e) \rangle = \langle \log r \rangle$

 $\langle v_e \rangle \approx v$

as long as the true $v \ll cz$

Estimators

A more accurate estimator at large redshift

 $v_e = \frac{cz_{mod}}{(1+z_{mod})} \log(cz_{mod}/H_o r_e)$

with uncertainty

$$\delta v_e = c z_{mod} \delta l_e / (1 + z_{mod})$$

We assume that actual velocity (v) is small compared to the redshift, not the estimated velocity (v_e) .

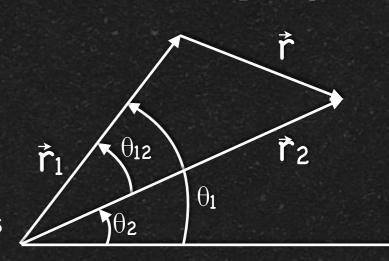
 $v_e = {
m few} imes 10^3 \, {
m km/s}$ whereas $v = {
m few} imes 10^2 \, {
m km/s}$ Should hold quite well for galaxies at distances $\gtrsim 20 \, {
m Mpc}$

 $\Psi_{ij}(\vec{r}) = \langle v_i(\vec{r}_1)v_j(\vec{r}_2) \rangle$

 $\vec{u}_1 = \hat{r}_1 u_1 = \hat{r}_1 (\hat{r}_1 \cdot \vec{v}_1)$

Correlation Tensor

The observable quantity line-of-site peculiar velocities



 $\vec{r} = \vec{r}_2 - \vec{r}_1$

 $\cos \theta_1 = \hat{r}_1 \cdot \hat{r} \qquad \cos \theta_2 = \hat{r}_2 \cdot \hat{r} \qquad \cos \theta_{12} = \hat{r}_1 \cdot \hat{r}_2$

The Line-of-site velocity correlations

$$\begin{split} \psi_1(r) &\equiv \frac{\Sigma_{\text{pairs}(r)} u_1 u_2 \cos \theta_{12}}{\Sigma_{\text{pairs}(r)} \cos^2 \theta_{12}} \,, \quad \text{The} \\ \psi_2(r) &\equiv \frac{\Sigma_{\text{pairs}(r)} u_1 u_2 \cos \theta_1 \cos \theta_2}{\Sigma_{\text{pairs}(r)} \cos \theta_{12} \cos \theta_1 \cos \theta_2} \end{split}$$

The dot product of the radial peculiar velocities

The product of the components of the Line-of-site velocities along the vector separating the two galaxies

Górski, 1988, ApJL, 332, L7 Górski, Davis, Strauss, White & Yahil, 1989, ApJ 344: 1–19 Torment etal 1993, ApJ 411, 16–33 Brogani, etal, 2000, AJ, 119, 102



July 5th, 2016

The line-of-site velocity correlations are related to the peculiar velocity correlations

$$\langle u_{1i}u_{2j}\rangle = \hat{r}_{1i}\hat{r}_{2j}\langle v_m v_n\rangle\hat{r}_{1m}\hat{r}_{2n}$$

where

$$\langle v_m v_n \rangle = \left[\phi_{\parallel}(r) - \phi_{\perp}(r) \right] \hat{r}_m \hat{r}_n + \phi_{\perp}(r) \delta_{mn}$$

Monin and Yaglom, 1975 Górski, 1988, ApJL, 332, L7

We can express ψ_1 and ψ_2 in terms of ϕ_{\parallel} and ϕ_{\perp} assuming homogeneity and isotropy

 $\psi_1(r) = A(r)\phi_{\parallel} + [1 - A(r)]\phi_{\perp}$ $\psi_2(r) = B(r)\phi_{\parallel} + [1 - B(r)]\phi_{\perp}$

Where A and B are the moments of the selection function

 $A(r) = \frac{\Sigma_{\text{pairs}(r)} \cos \theta_1 \cos \theta_2 \cos \theta_{12}}{\Sigma_{\text{pairs}(r)} \cos^2 \theta_{12}}$ $B(r) = \frac{\Sigma_{\text{pairs}(r)} \cos^2 \theta_1 \cos^2 \theta_2}{\Sigma_{\text{pairs}(r)} \cos \theta_{12} \cos \theta_1 \cos \theta_2}$

Inverting to get the Line-of-site Radial and Transverse Velocity Correlations

$$\phi_{\parallel}(r) = \frac{[1 - B(r)]\psi_1(r) - [1 - A(r)]\psi_2(r)}{A(r) - B(r)}$$
$$\phi_{\perp}(r) = \frac{B(r)\psi_1(r) - A(r)\psi_2(r)}{B(r) - A(r)}$$

In general, use the distance estimator as r.

At large distances the redshift is a better estimator of the distance.

Use redshift measurements as the distances in the velocity correlations analysis.





Working equations

The Line-of-site Velocity Correlations Velocity Correlations $\psi_1(r) \equiv \frac{\Sigma_{\text{pairs}(r)} u_1 u_2 \cos \theta_{12}}{\Sigma_{\text{pairs}(r)} \cos^2 \theta_{12}},$ $\psi_2(r) \equiv \frac{\sum_{\text{pairs}(r)} u_1 u_2 \cos \theta_1 \cos \theta_2}{\sum_{\text{pairs}(r)} \cos \theta_{12} \cos \theta_1 \cos \theta_2} \,.$

The Radial and Transverse Velocity Correlations

$$\phi_{\parallel}(r) = \frac{[1 - B(r)]\psi_1(r) - [1 - A(r)]\psi_2(r)}{A(r) - B(r)}$$
$$\phi_{\perp}(r) = \frac{B(r)\psi_1(r) - A(r)\psi_2(r)}{B(r) - A(r)}$$

The Moments of the Selection Functions

 $A(r) = \frac{\sum_{\text{pairs}(r)} \cos \theta_1 \cos \theta_2 \cos \theta_{12}}{\sum_{\text{pairs}(r)} \cos^2 \theta_{12}}$ $B(r) = \frac{\Sigma_{\text{pairs}(r)} \cos^2 \theta_1 \cos^2 \theta_2}{\Sigma_{\text{pairs}(r)} \cos \theta_{12} \cos \theta_1 \cos \theta_2}$



In linear theory the radial and transverse correlation functions of the threedimensional peculiar velocity field are a convolution of the power spectrum of density fluctuations with the window functions

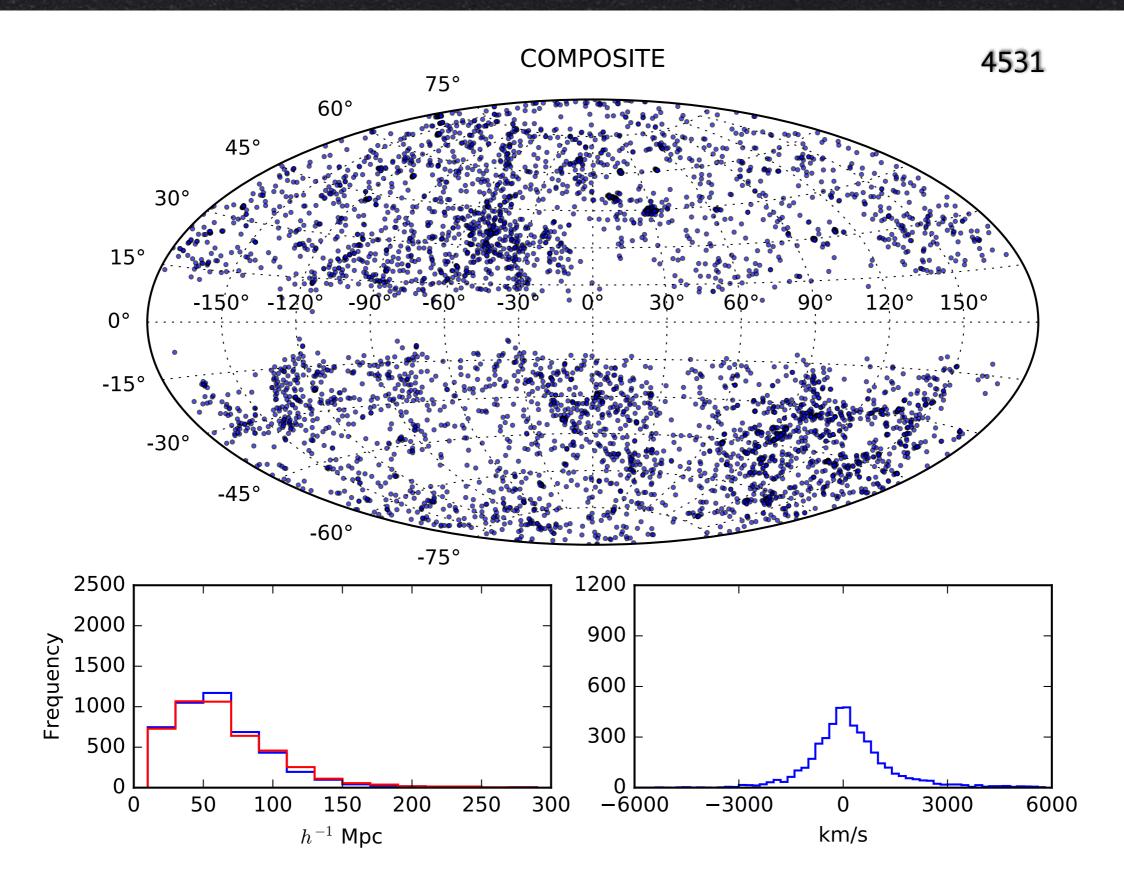
Górski, 1988, ApJL, 332, L7 Brogani, etal, 2000, AJ, 119, 102

Where $j_i(x)$ is the ith order spherical Bessel function

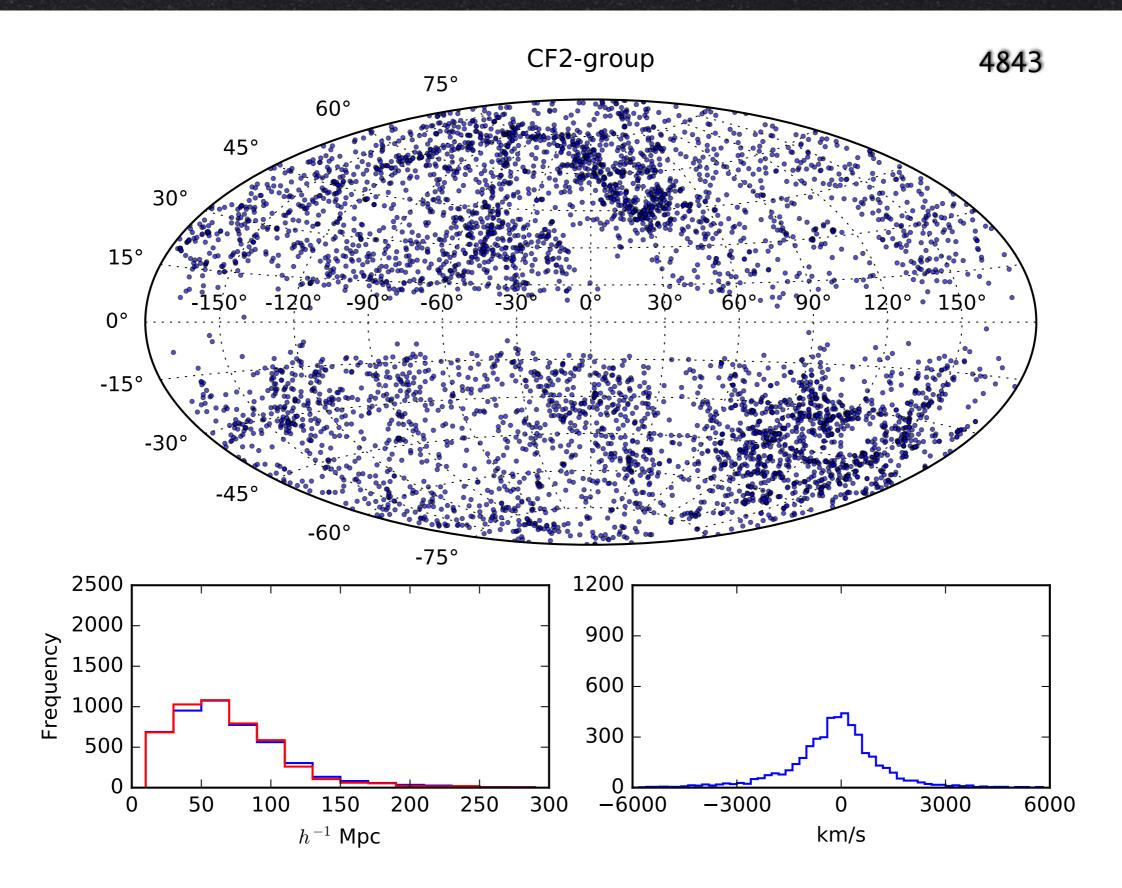
and

$$f(\Omega_m) \approx \Omega_m^{0.55}$$

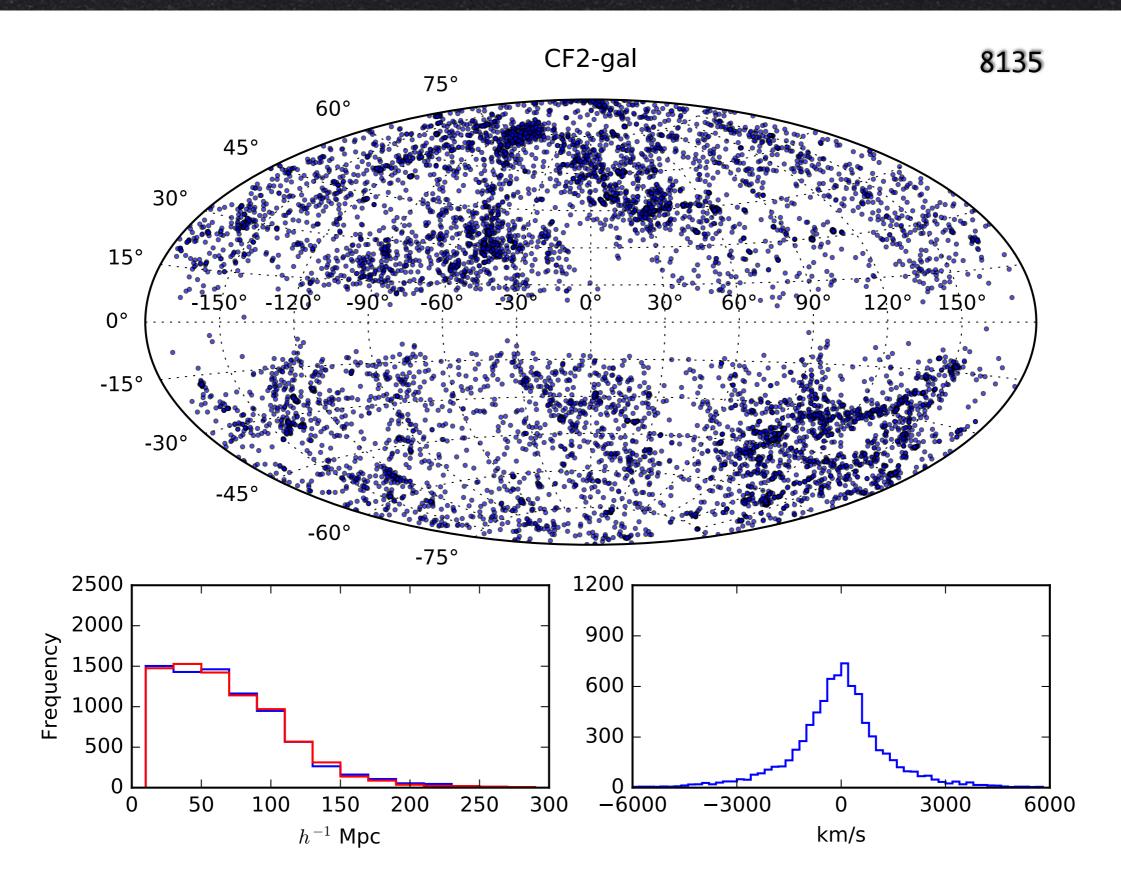
Linder, 2005, PRD, 72, 043529



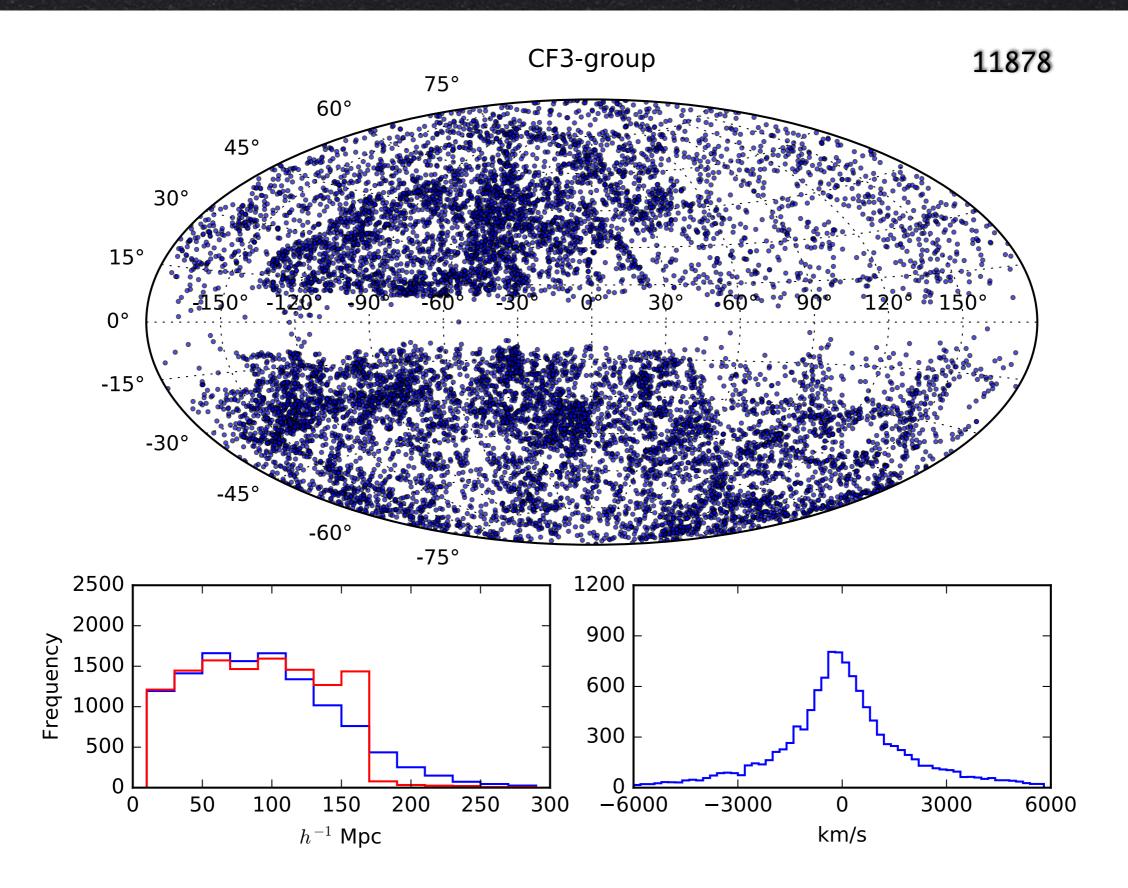
July 5th, 2016



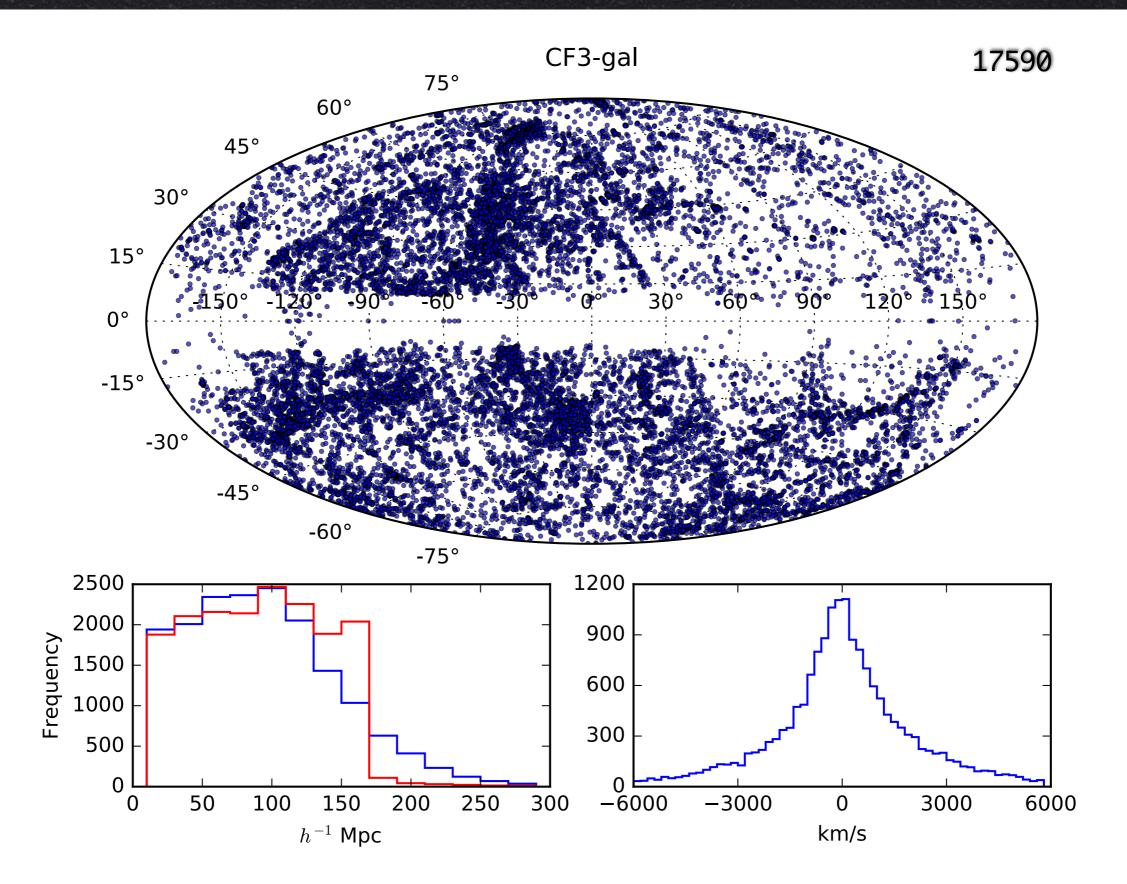
July 5th, 2016



July 5th, 2016

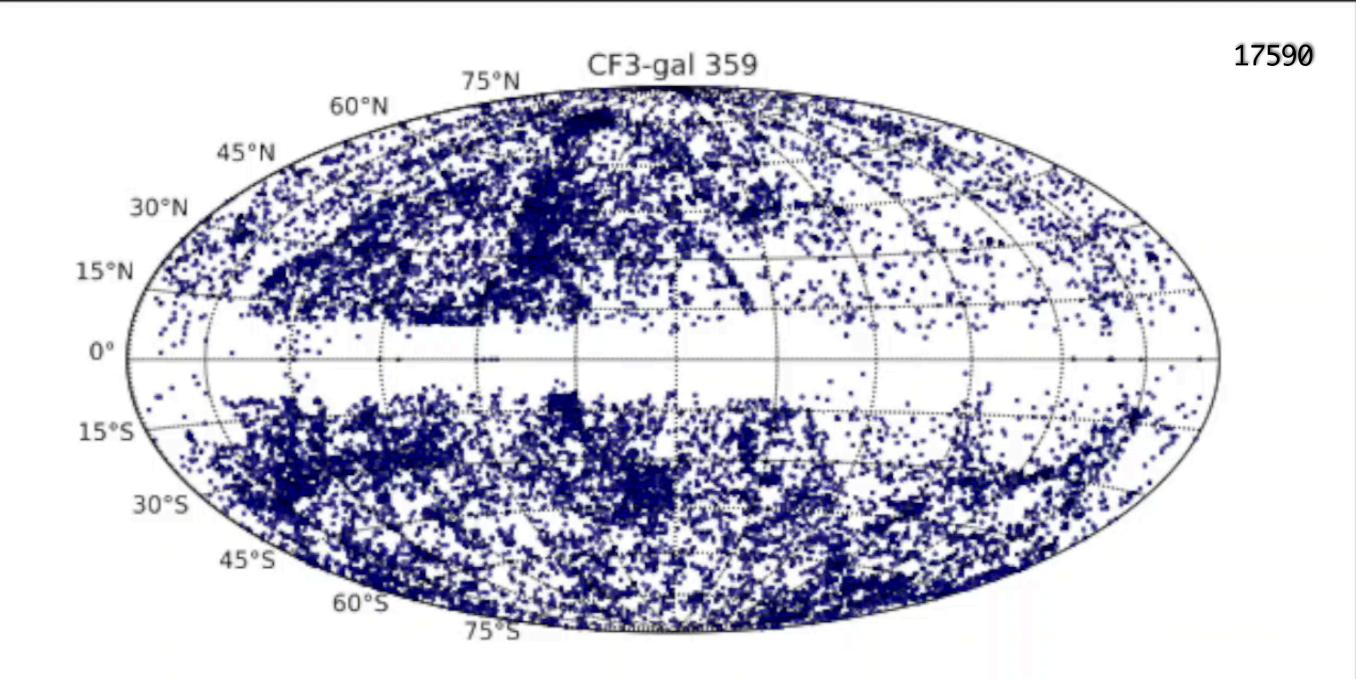


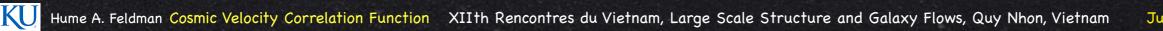
July 5th, 2016



July 5th, 2016

Data

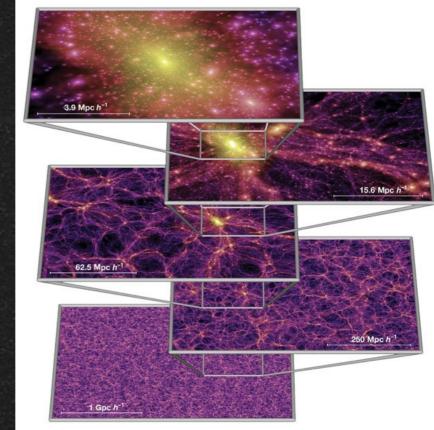






Use the Millennium Simulation, a dark matter only simulation using GADGET-2 simulation code

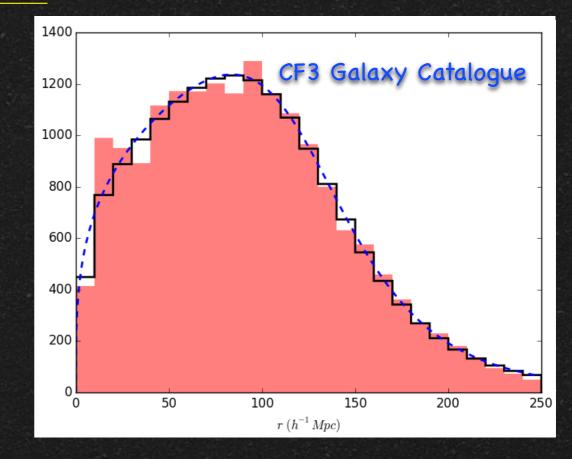
Matter density, Ω_m	0.25
Cosmological constant density, Ω_{Λ}	0.75
Baryon density, Ω_b	0.04
Hubble parameter, $h (100 km s^{-1} Mp c^{-1})$	0.73
Amplitude of matter density fluctuations, σ_8	0.9
Primordial scalar spectral index, n_s	1.0
Box size $(h^{-1}Mpc)$	500
Number of particles	2160
Particle mass, $m_p \ (10^8 h^{-1} M_{\odot})$	8.61
Softening, $f_c \ (h^{-1}kpc)$	5



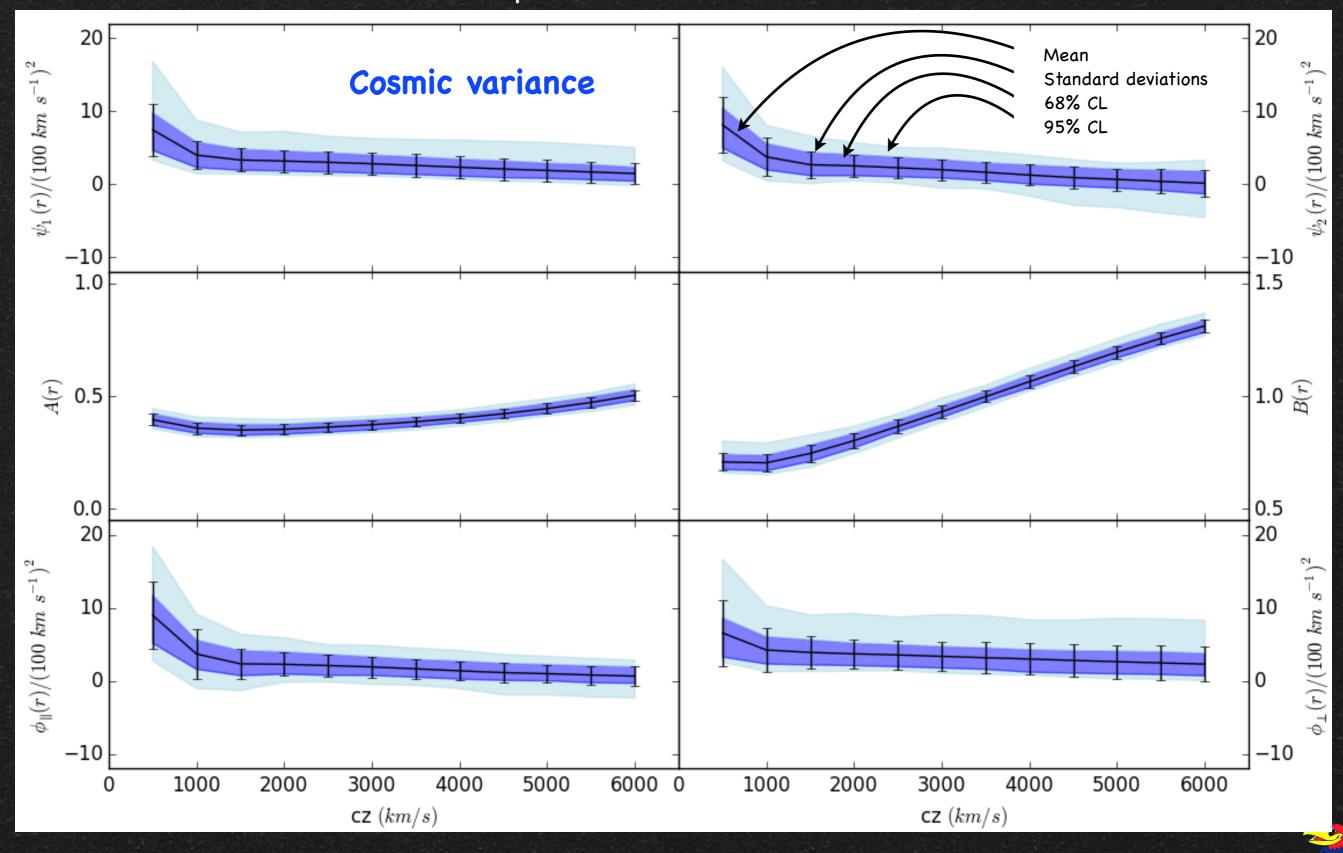
Springel etal, 2005, Nature, 435, 629 De Lucia G., Blaizot J., 2007, MNRAS, 375, 2 http://www.mpa.mpa-garching.mpg.de/millennium/http://www.mpa.mpa-garching.mpg.de/millennium/ https://www.mpa.mpa-garching.mpg.de/galform/virgo/millennium/

We select the galaxies for each mock survey by ensuring a best fit to the radial selection function fitted with

$$f(r) = \mathscr{A}\left(\frac{r}{r_0}\right)^{n_1} \left[1 + \left(\frac{r}{r_0}\right)^{n_1 + n_2}\right]$$



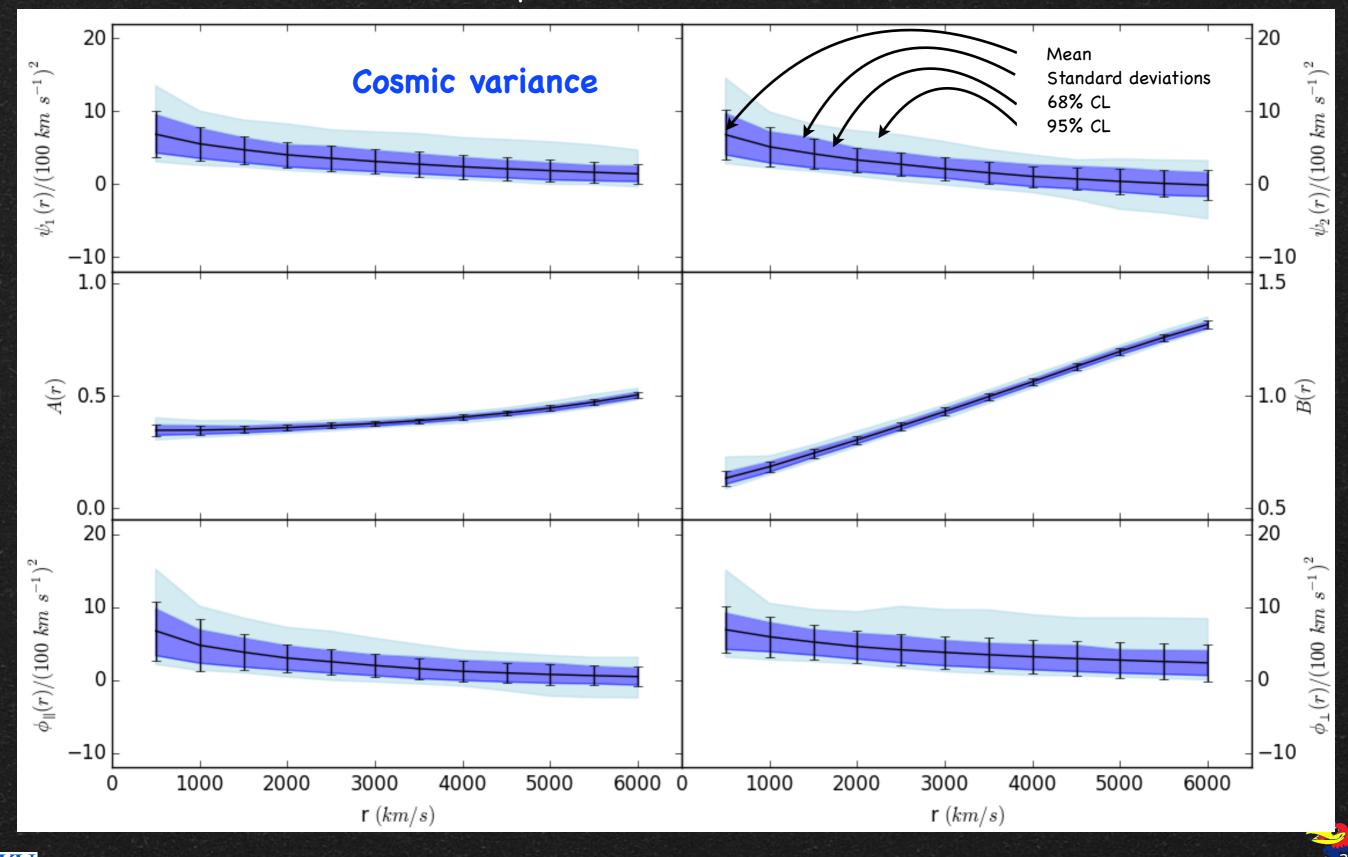
We chose 100 independent mock surveys and calculated the correlation functions with exact positions from the simulations



Hume A. Feldman Cosmic Velocity Correlation Function XIIth Rencontres du Vietnam, Large Scale Structure and Galaxy Flows, Quy Nhon, Vietnam July

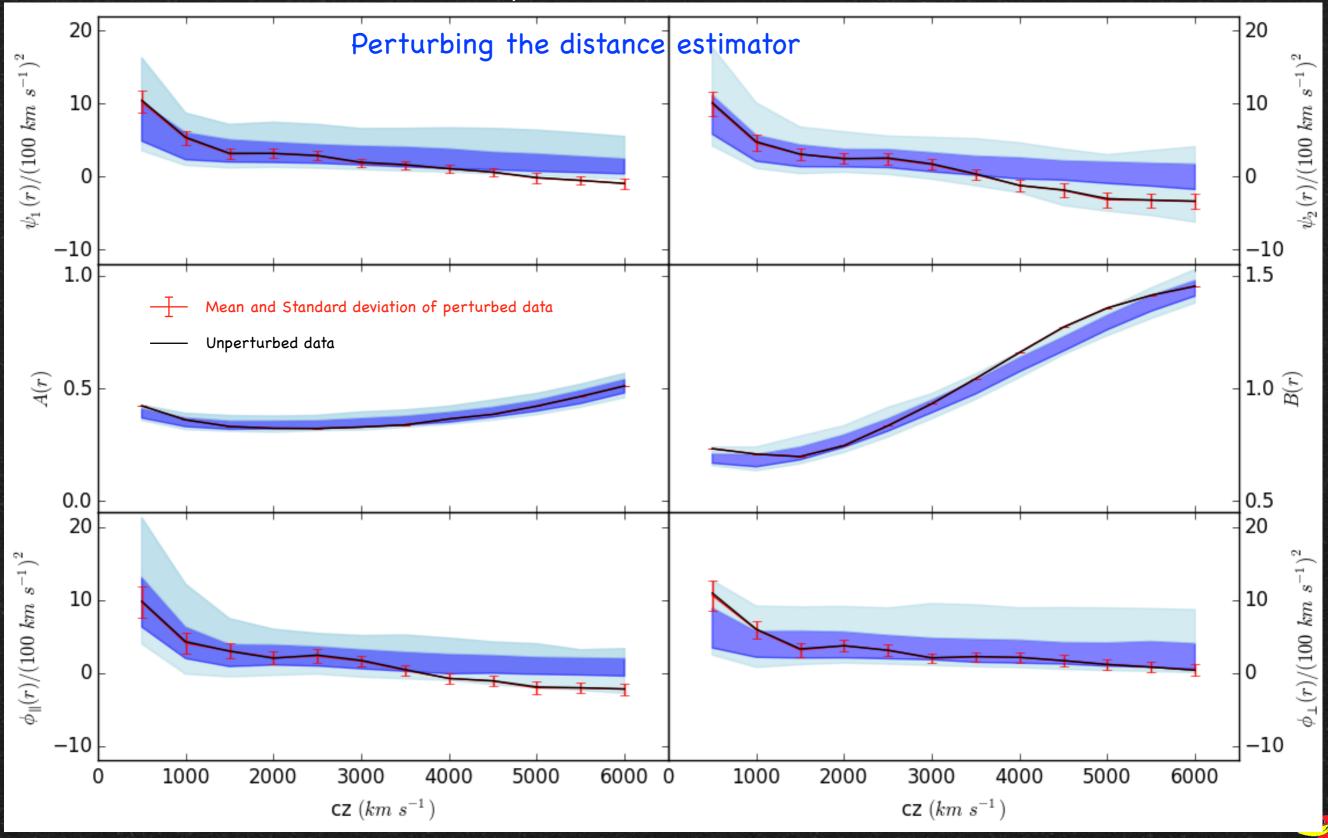
July 5th, 2016

We chose 100 independent mock surveys and calculated the correlation functions with exact positions from the simulations

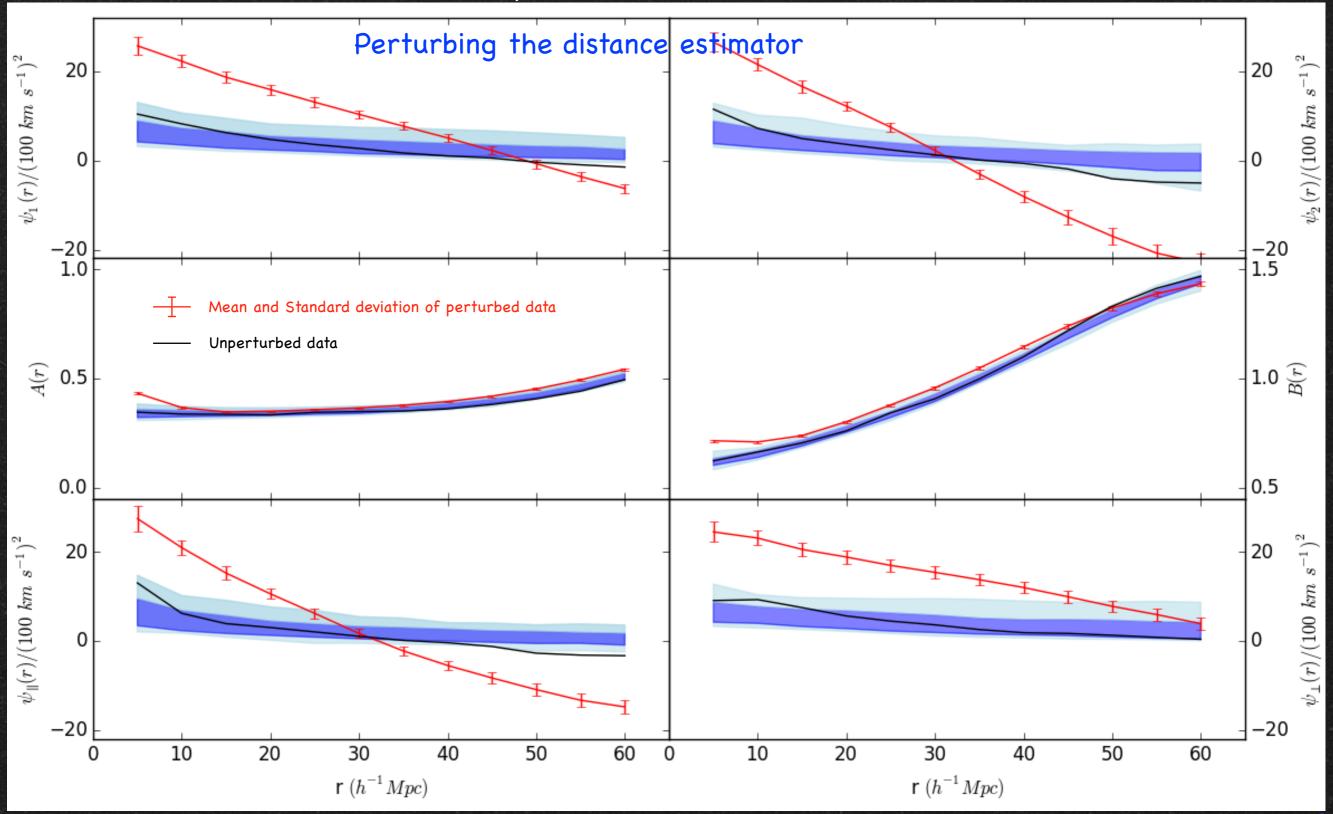


Hume A. Feldman Cosmic Velocity Correlation Function XIIth Rencontres du Vietnam, Large Scale Structure and Galaxy Flows, Quy Nhon, Vietnam July 5th, 2016

The other main source of uncertainty in the correlation functions comes from the the uncertainty of the peculiar velocity measurements

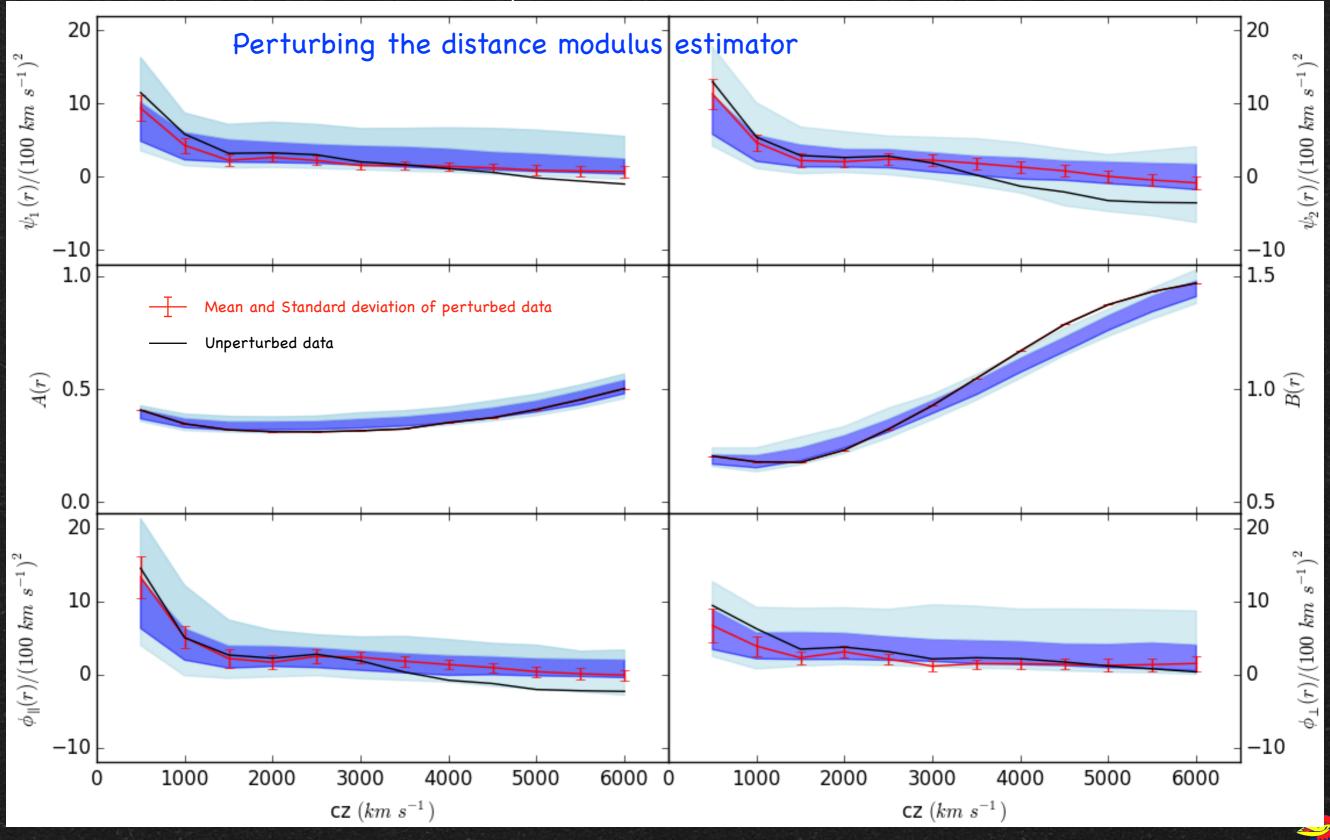


The other main source of uncertainty in the correlation functions comes from the the uncertainty of the peculiar velocity measurements



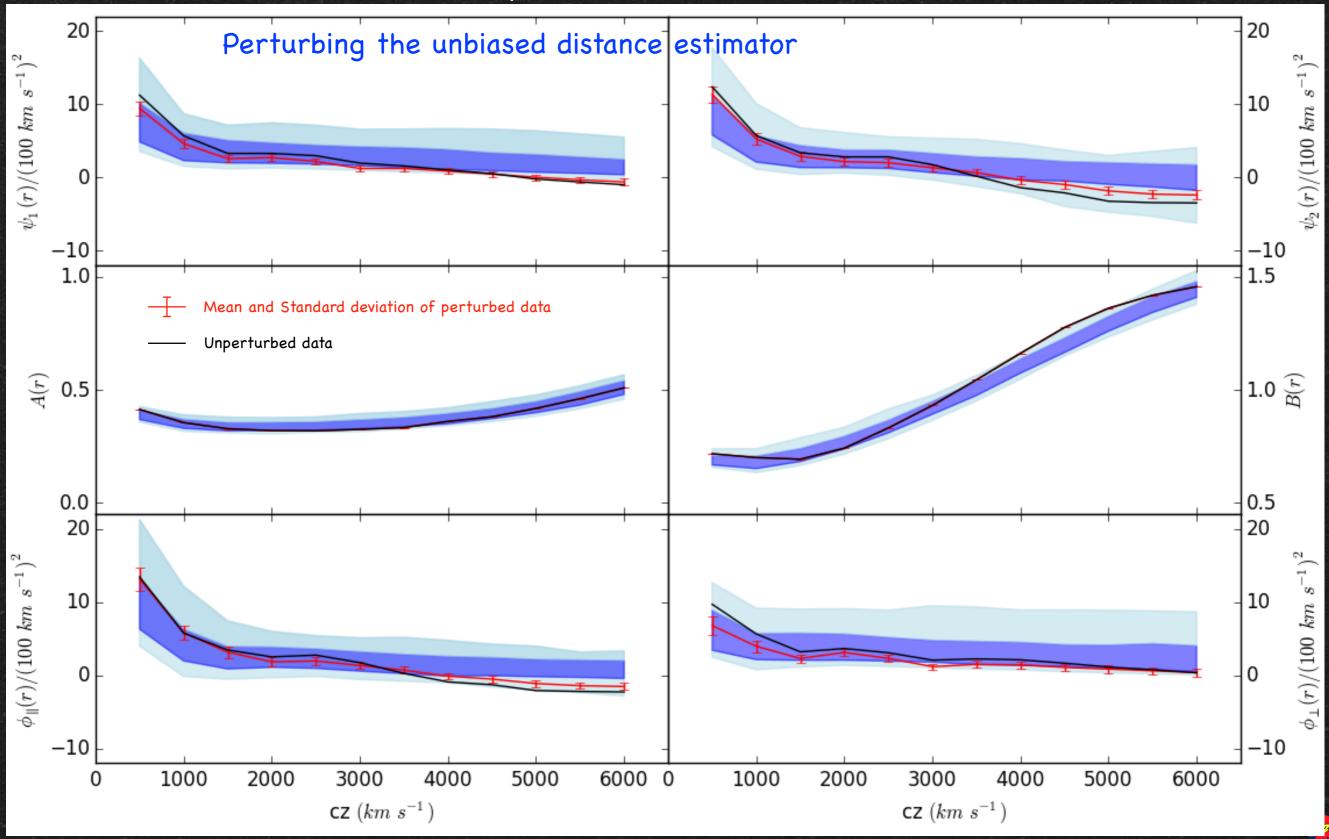
July 5th, 2016

The other main source of uncertainty in the correlation functions comes from the the uncertainty of the peculiar velocity measurements



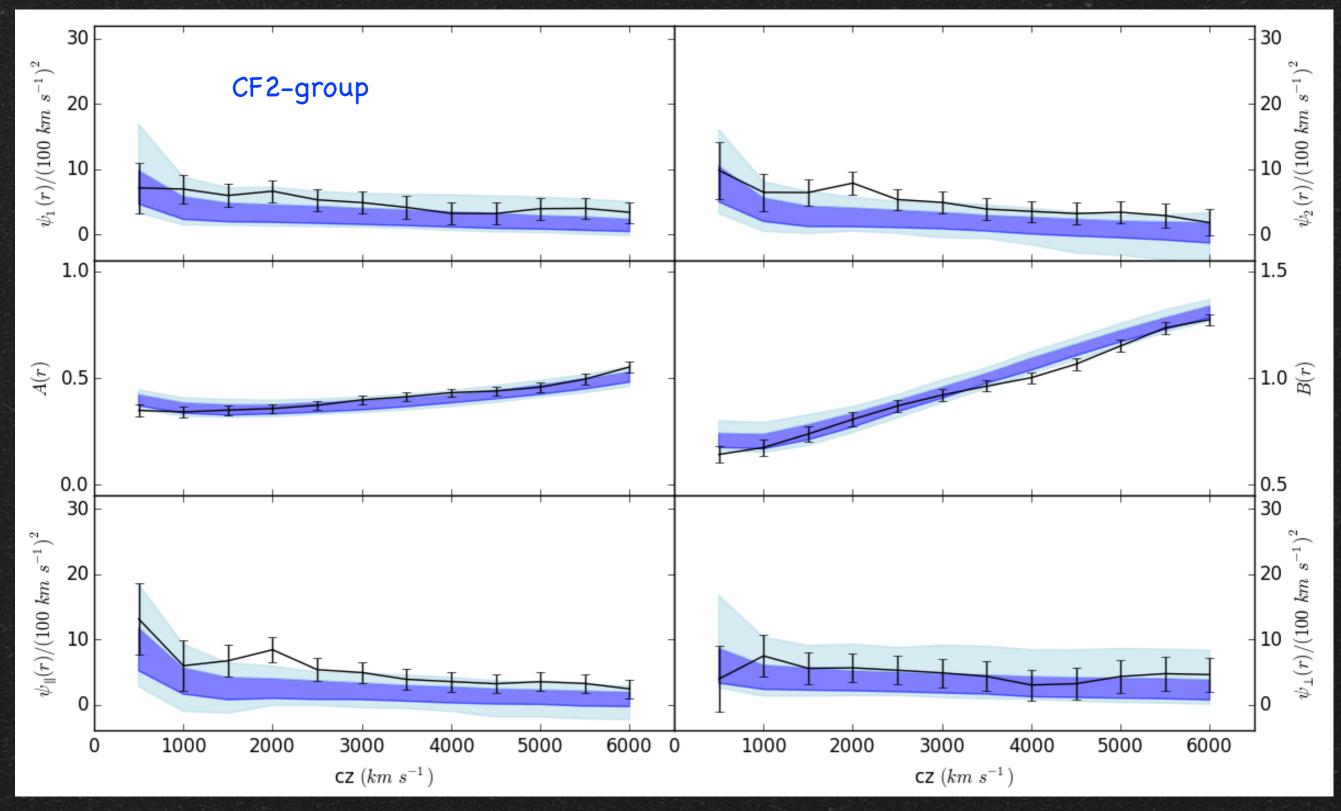
KU Hume A. Feldman Cosmic Velocity Correlation Function XIIth Rencontres du Vietnam, Large Scale Structure and Galaxy Flows, Quy Nhon, Vietnam July 5th, 2016

The other main source of uncertainty in the correlation functions comes from the the uncertainty of the peculiar velocity measurements

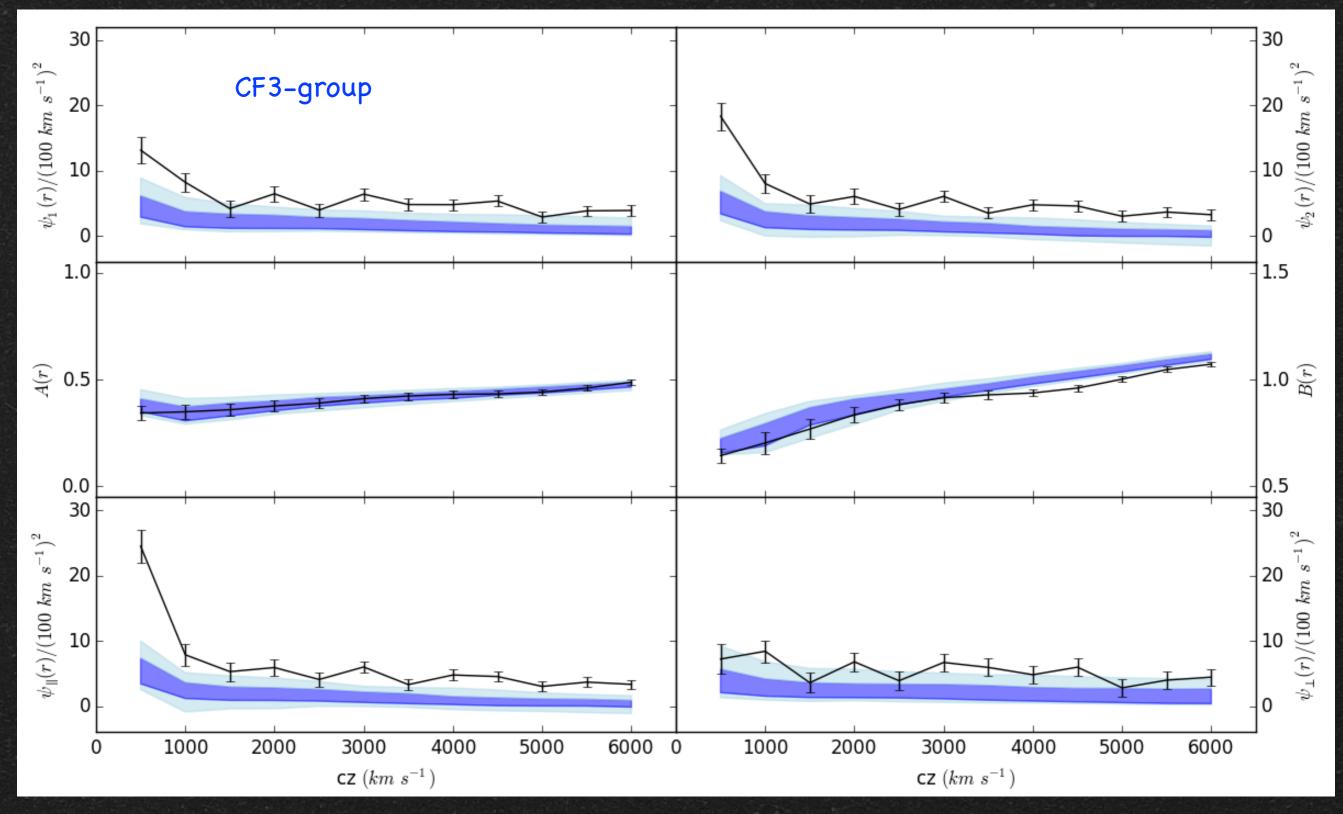


KU Hume A. Feldman Cosmic Velocity Correlation Function XIIth Rencontres du Vietnam, Large Scale Structure and Galaxy Flows, Quy Nhon, Vietnam July 5

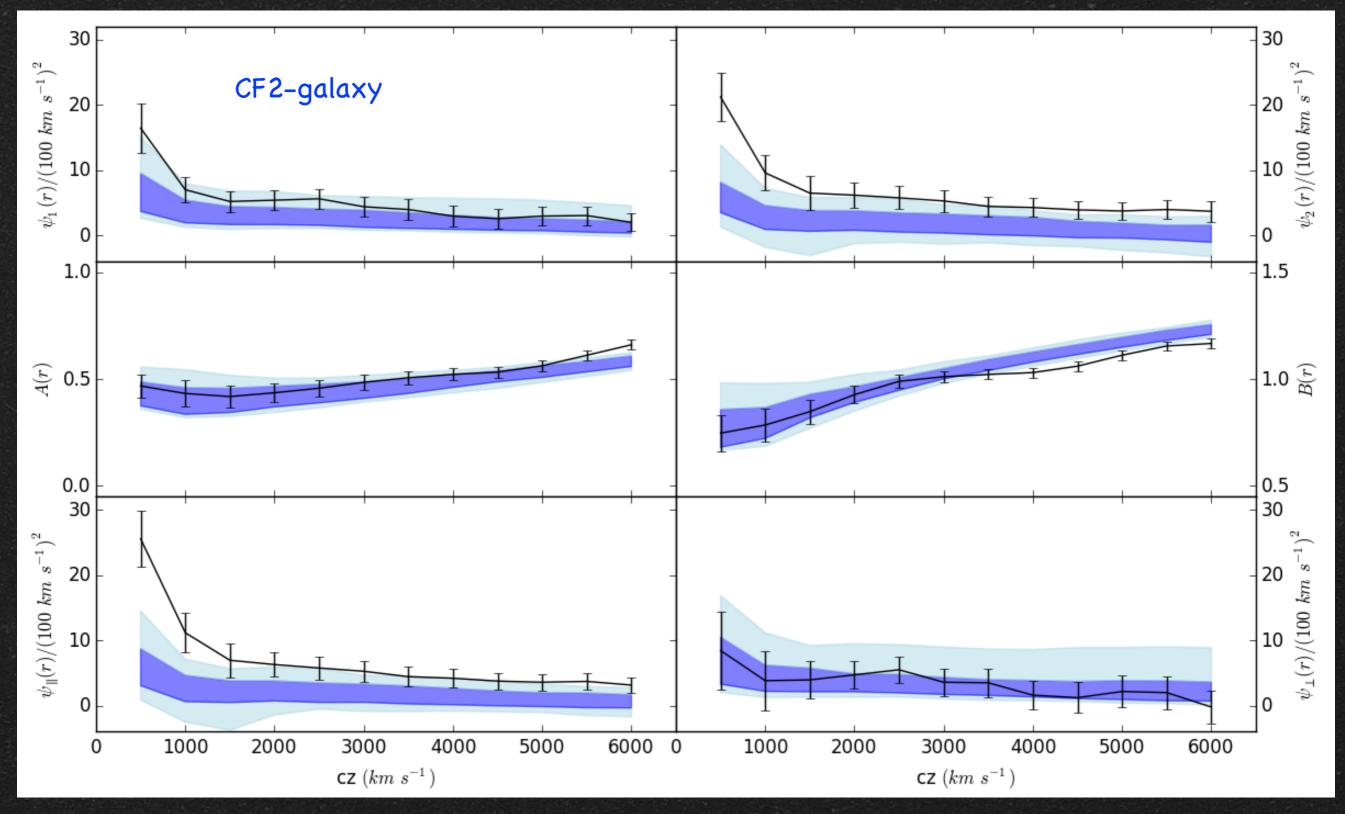
Observational Data



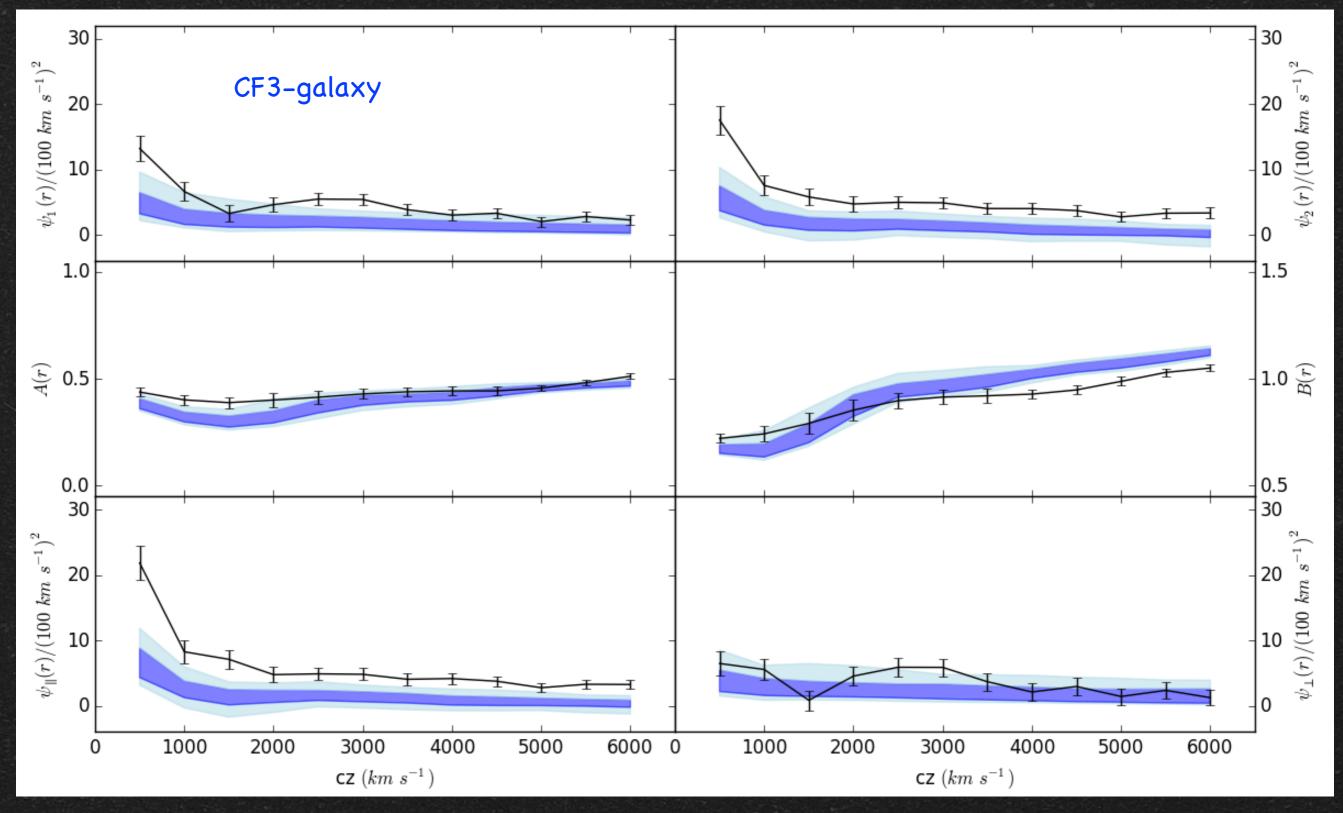
Observational Data



Observational Data



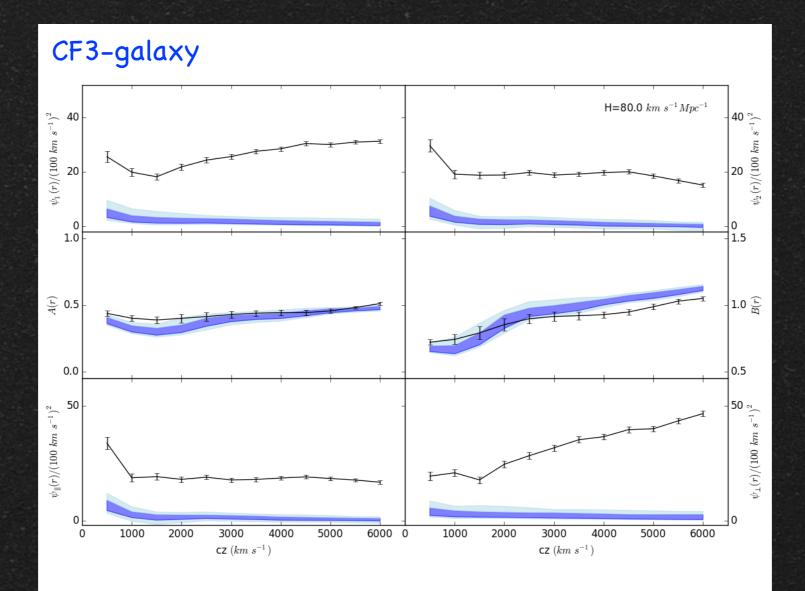
Observational Data



Use the correlation function to constrain H.

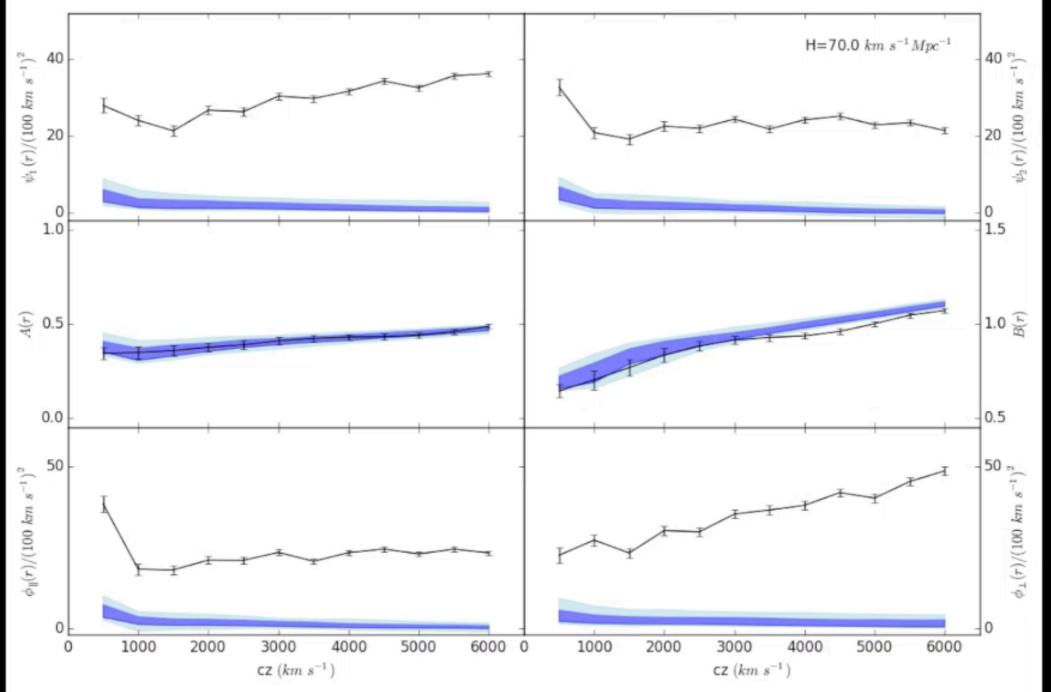
Create catalogues with different H.

Since we are using cz as the distance, only v_{p} changes

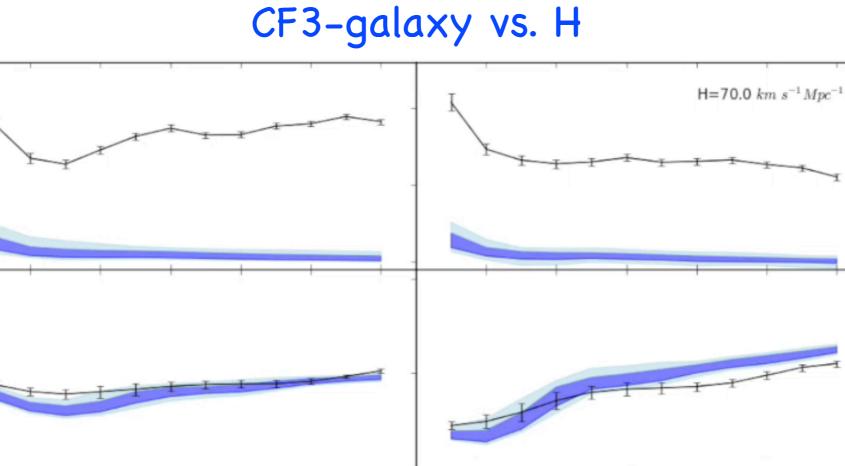


July 5th, 2016





KU Hume A. Feldman Cosmic Velocity Correlation Function XIIth Rencontres du Vietnam, Large Scale Structure and Galaxy Flows, Quy Nhon, Vietnam

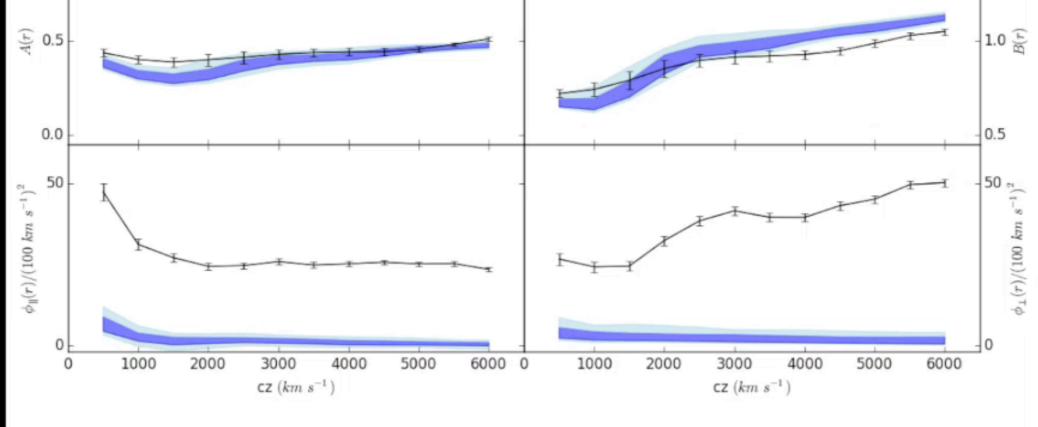


 $\psi_1\left(r\right)/(100 \; km \; s^{-1} \,)^2$

40

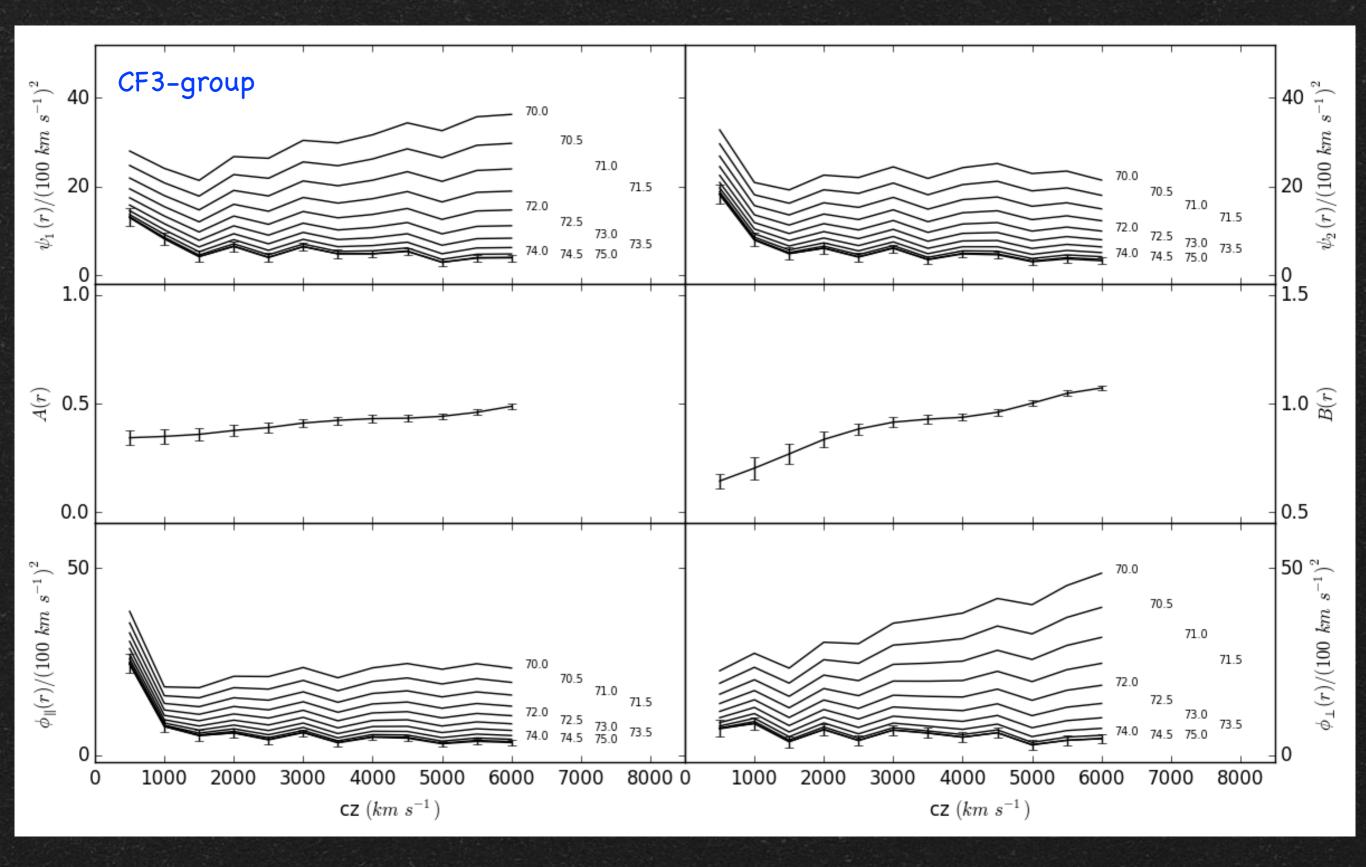
20

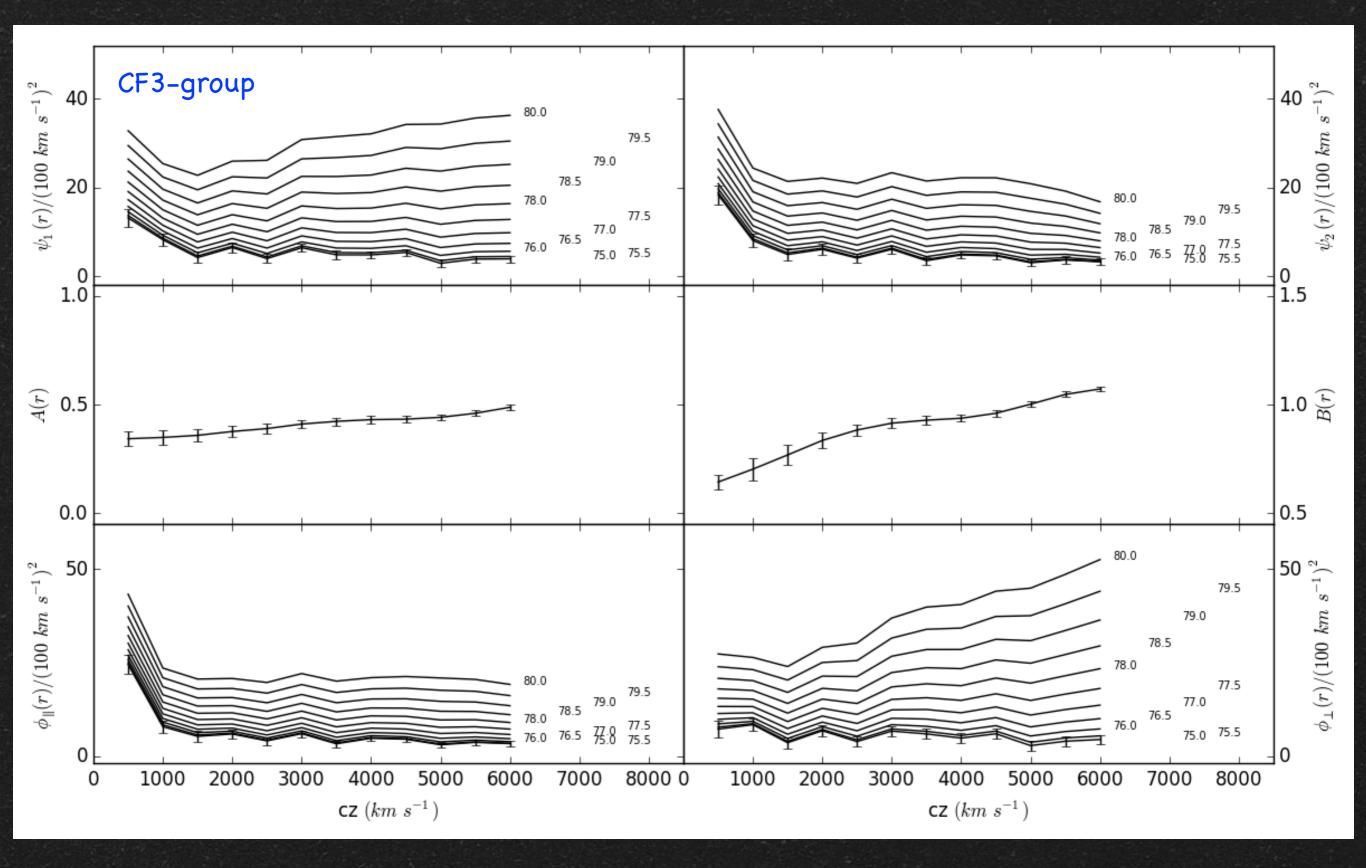
0



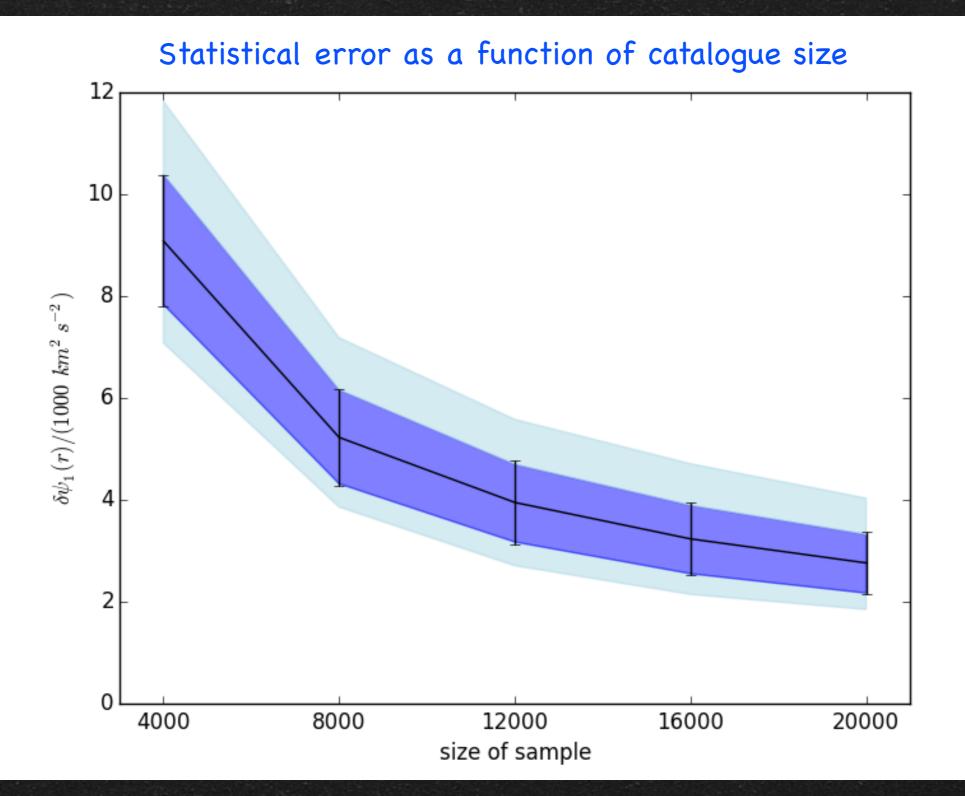
31

0 1.5





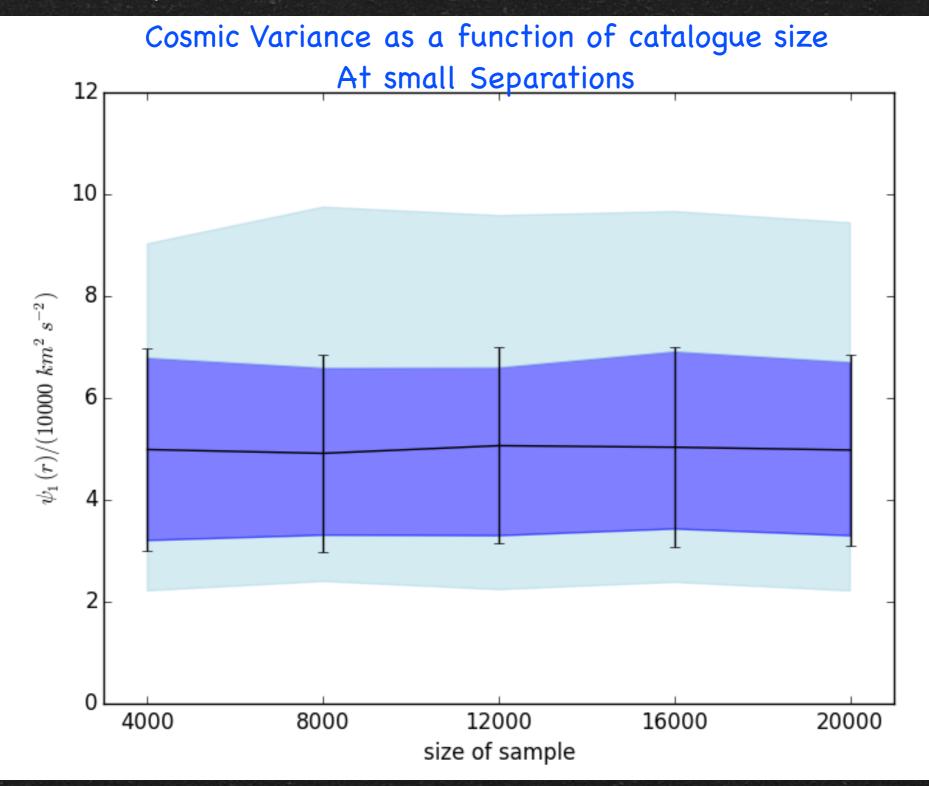
Velocity Correlations Errors as a function of sample size The statistical errors get smaller with sample size



Velocity Correlations Errors as a function of sample size

The statistical errors get smaller with sample size

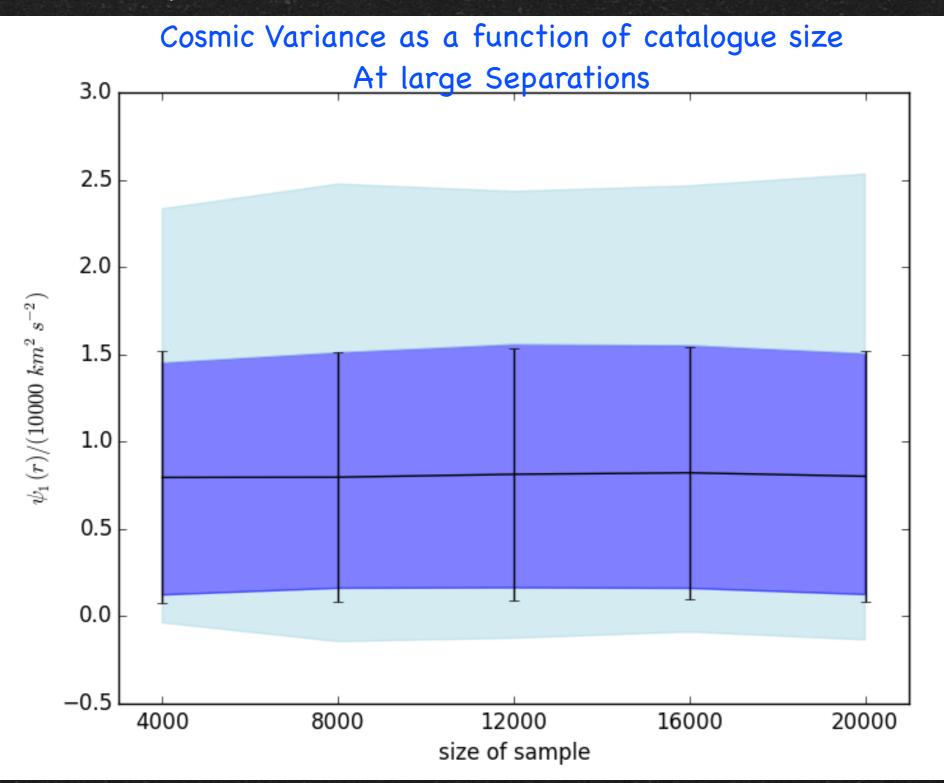
Cosmic Variance stays the same



Velocity Correlations Errors as a function of sample size

The statistical errors get smaller with sample size

Cosmic Variance stays the same



Velocity Correlations Conclusions

Correlation function statistic is robust and stable across different samples.

Provides a constraint on H = 75 \pm 2 km s⁻¹ Mpc⁻¹

Coherence length of ~1,500 \pm 500 km s⁻¹

Amplitude 5-10 $(10^2 \text{ km s}^{-1})^2$

Statistical errors get smaller with sample size

Cosmic variance is independent of sample size

Agrees with Λ CDM expectations with Planck/WMAP parameters

 $\Gamma = h \ \Omega = 0.3 \pm 0.15$

 $\sigma_8 = 0.85 \pm 0.6$

preliminary results

 $\eta_8 = \Omega^{0.55} \sigma_8 = 0.42 \pm 0.2$



