The luminosity distance - redshift relation in cosmologies beyond FLRW

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Background and motivation

- Standard Cosmological model ΛCDM
- Type Ia supernovae allow you to probe the geometry of the Universe
- Data is usually fitted with a d_L(z) relation based on FLRW cosmology
- Departures from FLRW due to:
 - inhomogeneities from cosmic structure
 - alternative cosmological models
 - \implies Need an improved formula for $d_L!$
- Main idea: use data from type Ia supernovae and an improved d_L(z) formula to analyse alternative models and possible deviations from FLRW

Main papers:

- arXiv:1802.06550v2
- arXiv:1803.08766

Cosmography:

arXiv:gr-qc/0411131

van Vleck determinant:

arXiv:hep-th/9303020

CFLRW cosmologies:

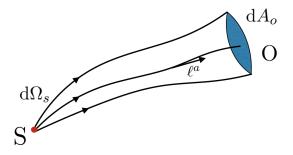
arXiv:1502.02758

Other approaches:

arXiv:astro-ph/0511183v5

arXiv:1209.4326v2

The luminosity distance and the redshift



 L_s - intrinsic luminosity of S F_o - measured flux at O

$$d_L := \sqrt{\frac{L_s}{4\pi F_o}} = (1+z)\sqrt{\frac{\mathrm{d}A_o}{\mathrm{d}\Omega_s}}$$
$$1+z := \frac{\omega_s}{\omega_o} = \frac{(g_{\mu\nu}U^{\mu}\ell^{\nu})_s}{(g_{\mu\nu}U^{\mu}\ell^{\nu})_o}$$

The van Vleck determinant



Definition:
$$\Delta_{vV}(x_s, x_o) := -\frac{\det\{\nabla^{x_s}_{\mu} \nabla^{x_o}_{\nu} \sigma_{\gamma}(x_s, x_o)\}}{\sqrt{g(x_s)g(x_o)}}$$

$$\sigma_\gamma(x_s,x_o):=\pmrac{1}{2}[s_\gamma(x_s,x_o)]^2$$
 geodesic interval

$$s_{\gamma}(x_s,x_o) = egin{cases} \int_{\gamma} \mathrm{d} au & ext{timelike} \ \int_{\gamma} \mathrm{d}s & ext{spacelike} \ 0 & ext{null} \end{cases}$$

The van Vleck determinant and the luminosity distance

Evolution of van Vleck determinant along the null geodesic congruence:

$$\frac{d\Delta_{vV}}{d\lambda} = \left(\frac{2}{\lambda - \lambda_s} - \theta\right) \Delta_{vV}$$

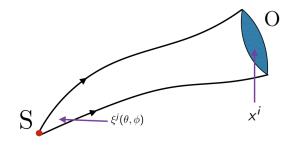
The expansion:
$$heta=rac{1}{dA/d\Omega_s}rac{d(dA/d\Omega_s)}{d\lambda}$$

Solution:
$$\Delta_{vV}(\lambda_o, \lambda_s) = \frac{(\lambda_o - \lambda_s)^2}{(dA_o/d\Omega_s)}$$

$$\implies d_L = (1+z) (\lambda_o - \lambda_s) \Delta_{vV}^{-1/2}$$

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The Jacobi determinant



$$J_j^i = rac{\partial x'}{\partial \xi^j}$$
 Jacobi map $rac{dA_o}{d\Omega_s} = |\det J|$ Jacobi determinant

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$$\implies d_L = (1+z)\sqrt{|\det J|}$$

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So far

$$egin{aligned} d_L &= (1+z) \left(\lambda_o - \lambda_s
ight) \Delta_{vV}^{-1/2} \ &= (1+z) |\det J|^rac{1}{2} \end{aligned}$$

$$|\det J| = (\lambda_o - \lambda_s)^2 \Delta_{vV}^{-1}$$

$$F_o = \frac{L_s}{4\pi d_L^2} = \frac{L_s \Delta_{vV}}{4\pi (1+z)^2 (\lambda_o - \lambda_s)^2}$$

$$\label{eq:lambda} \begin{split} \Delta_{\nu V} > 1 & \mbox{geodesic focussing} \\ \Delta_{\nu V} < 1 & \mbox{geodesic defocussing} \end{split}$$

FLRW: luminosity distance and redshift

$$ds_{FLRW}^{2} = a^{2}(\eta) \left\{ -d\eta^{2} + \left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta \ d\phi^{2}) \right] \right\}$$

$$1 + z = (1 + z_c)(1 + z_D) = \frac{a_o}{a_s} \frac{\gamma_s(1 - \vec{v_s}.\hat{n})}{\gamma_o(1 - \vec{v_o}.\hat{n})}$$
$$d_L = a_o(1 + z) \frac{\sin\left(\sqrt{k}\Delta\eta\right)}{\sqrt{k}}$$

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3 cases:

$$k = 0 \qquad d_L = a_o(1+z)\Delta\eta k = 1 \qquad d_L = a_o(1+z)\sin(\Delta\eta) k = -1 \qquad d_L = a_o(1+z)\sinh(\Delta\eta)$$

Simple cosmography

• Generally, express $\Delta \eta$ as a power series of z_c :

$$\Delta \eta = p(z_c)$$
 with $p(0) = 0$

• Cosmographic expansion ($z_D = 0$, k = 0 case):

$$d_L(z) = \frac{1}{H_o} \Big\{ z + \frac{1}{2} \Big[1 - q_o \Big] z^2 - \frac{1}{6} \Big[1 - q_o - 3q_o^2 + j_o \Big] z^3 + O(z^4) \Big\}$$

Cosmographic parameters:

$$H = \frac{\dot{a}}{a}$$
 $q = -\frac{1}{H^2}\frac{\ddot{a}}{a}$ $j = \frac{1}{H^3}\frac{\ddot{a}}{a}$

- From observational $d_L(z)$ curve one can constrain H_o , q_o , j_o .
- Want to investigate deviations from simple FLRW and simple cosmography.

Deviations from FLRW

- Nonperturbative vs. perturbative
- Conformally FLRW

$$ds_{CFLRW}^{2} = f^{2}(x) \left\{ -d\eta^{2} + \left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta \ d\phi^{2}) \right] \right\}$$

• Linearly perturbed FLRW, k = 0

$$\mathrm{d}s^{2} = a^{2}(\eta) \Big[-(1+2\Psi(x))\mathrm{d}\eta^{2} + (1+2\Phi(x))\delta_{ij}\mathrm{d}x^{i}\mathrm{d}x^{j} \Big]$$

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Conformally FLRW (CFLRW)

$$ds_{CFLRW}^{2} = f^{2}(x) \left\{ -d\eta^{2} + \left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta \ d\phi^{2}) \right] \right\}$$

The only way to introduce nonperturbative deviations from FLRW without distorting the CMB

$$f(x) =: \exp[heta(x)] a(\eta) \qquad < e^{ heta} >_{spatial} = 1 \qquad \Delta heta_{LS} \le 10^{-5}$$

$$1+z=\frac{f_o}{f_s}=\exp(\theta_o-\theta_s)\frac{a_o}{a_s}=(1+z_{loc})(1+z_c)$$

$$d_L = f_o(1+z) rac{\sin(\sqrt{k}\Delta\eta)}{\sqrt{k}}$$

Cosmography: $\Delta \eta = p(z_c)$

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Perturbed FLRW: the luminosity distance

2 scalar mode perturbations $\Psi, \Phi << 1$, Newtonian gauge, k = 0:

$$\mathrm{d}s^{2} = a^{2}(\eta) \Big[-(1+2\Psi(\vec{x},\eta))\mathrm{d}\eta^{2} + (1+2\Phi(\vec{x},\eta))\delta_{ij}\mathrm{d}x^{i}\mathrm{d}x^{j} \Big]$$

$$\begin{aligned} d_L &= \frac{a_o^2}{a_s} \Big[(\eta_o - \eta_s) + 2 \Phi_o (\eta_o - \eta_s) - \Psi_s (\eta_o - \eta_s) \\ &+ (\eta_o - \eta_s) \int_{\eta_s}^{\eta_o} \vec{\nabla} \xi \cdot \hat{n} \mathrm{d}\eta + \int_{\eta_s}^{\eta_o} \xi \mathrm{d}\eta + \int_{\eta_s}^{\eta_o} (\eta_o - \eta) (\vec{\nabla} \xi \cdot \hat{n}) \mathrm{d}\eta \\ &- \frac{1}{2} \int_{\eta_s}^{\eta_o} (\eta_o - \eta) (\nabla^2 \xi - n^i n^j \xi_{,ij}) (\eta - \eta_s) \mathrm{d}\eta \Big] \end{aligned}$$

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where $\xi = \Psi - \Phi$.

Perturbed FLRW: the redshift

$$1 + z = \frac{a_o}{a_s} \left(1 + \Psi_o - \Psi_s - \int_{\eta_s}^{\eta_o} \partial_\eta \xi \mathrm{d}\eta \right)$$

A product of 3 contributions:

1.
$$1 + z_c = \frac{a_o}{a_s}$$

2. $1 + z_{loc} = \sqrt{\frac{1+2\Psi_o}{1+2\Psi_s}} \approx 1 + \Psi_o - \Psi_s$
3. $1 + z_{ISW} = 1 - \int_{\eta_s}^{\eta_o} \partial_\eta \xi d\eta$

One can also add the Doppler redshift $1 + z_D = \frac{\gamma_s(1 - \vec{v_s}.\hat{n})}{\gamma_o(1 - \vec{v_o}.\hat{n})}$:

$$1 + z = (1 + z_D)(1 + z_c)(1 + z_{gr})(1 + z_{ISW})$$

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Summary and future work

So far:

- Need an improved formula for d_L to test alternative models
- Derived such a formula for CFLRW and linearly perturbed FLRW

Future:

- ► Within a particular model derive a d_L(z) relation (ideally as a cosmographic series)
- Use data from supernovae to constrain the parameter space of this model
- Look at the power spectrum of d_L at a fixed z or at the correlation of d_L at different z along a single line of sight

Thank you for your attention!