



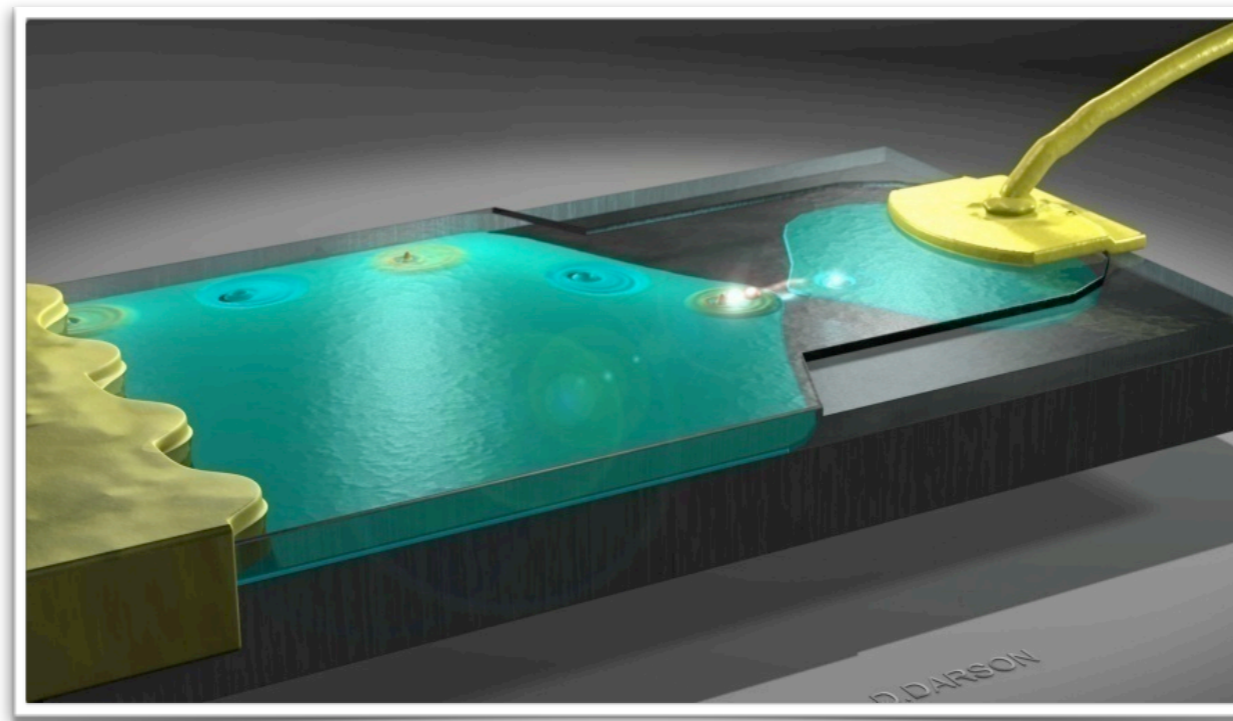
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Erwann Bocquillon

Intensity correlations of a single electron beam

Charge and heat dynamics in nanosystems - Orsay - 10/11/2011



Laboratoire Pierre Aigrain
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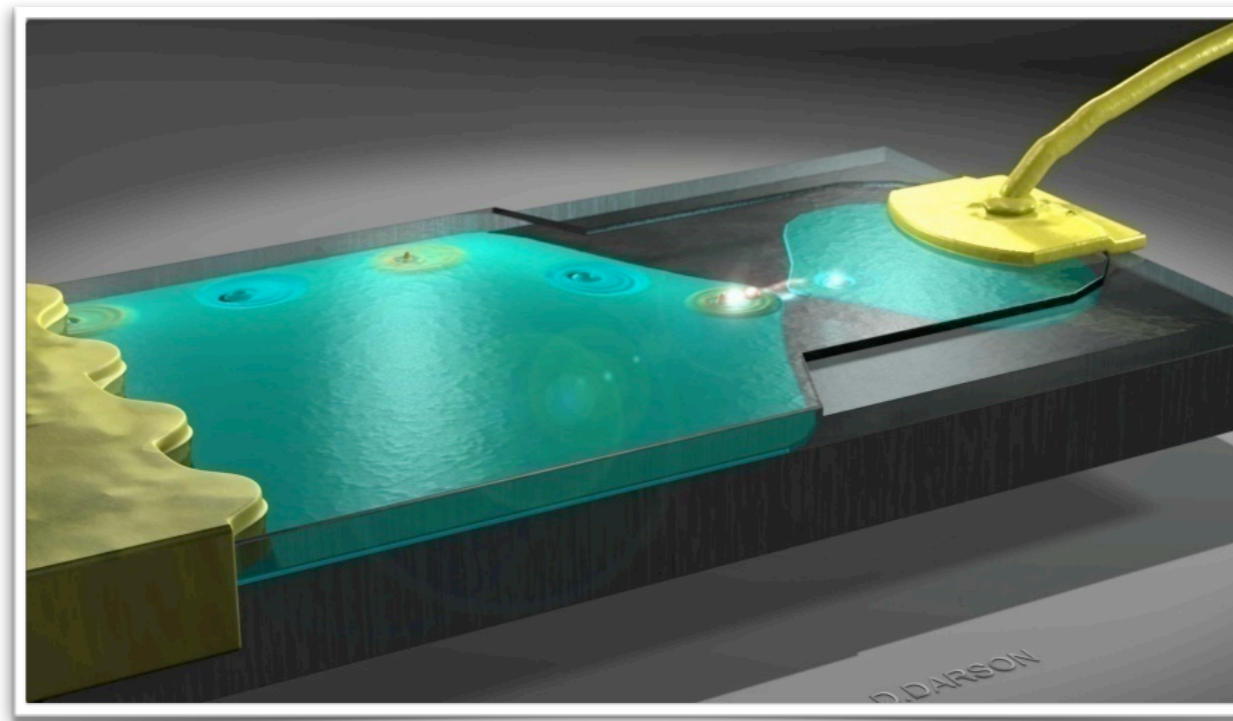
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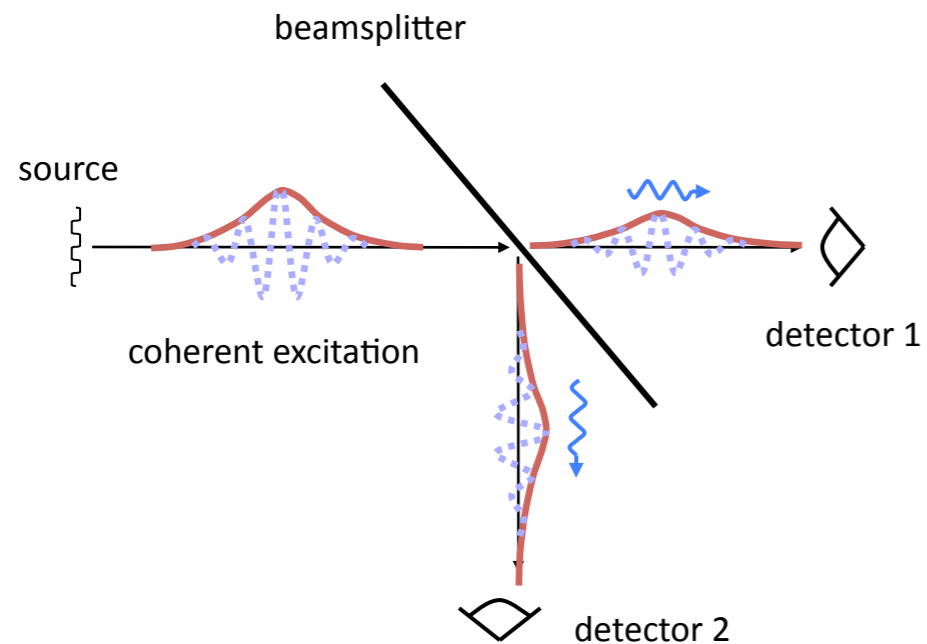
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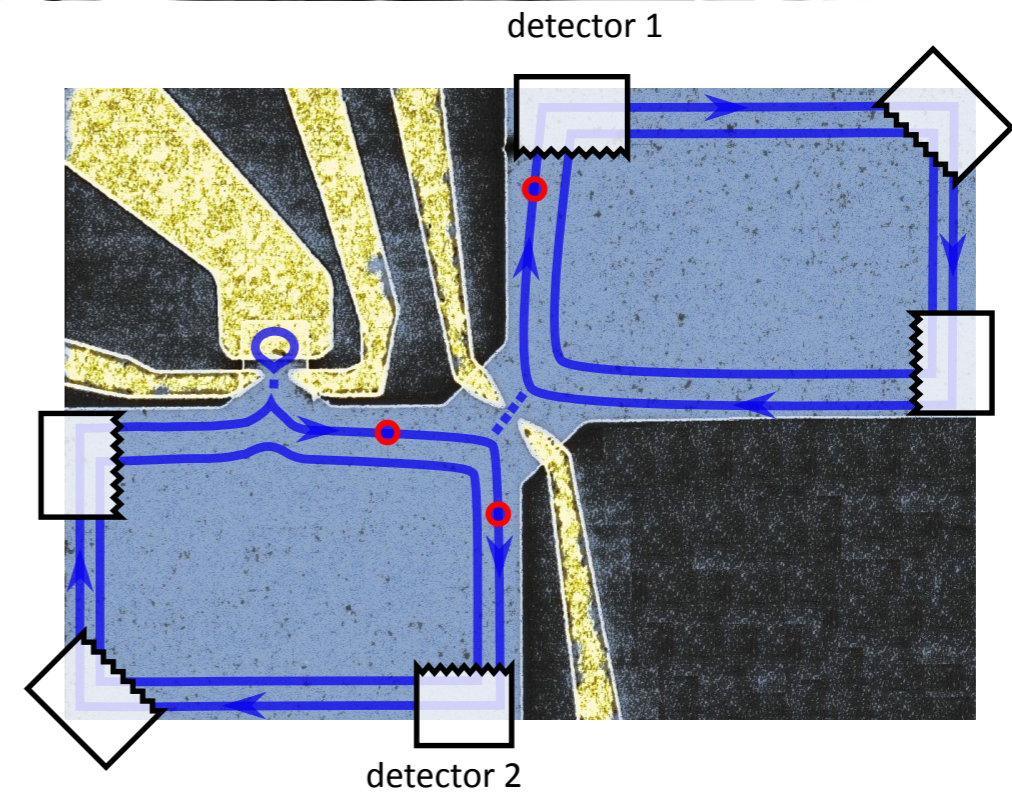
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HBT correlations in optics

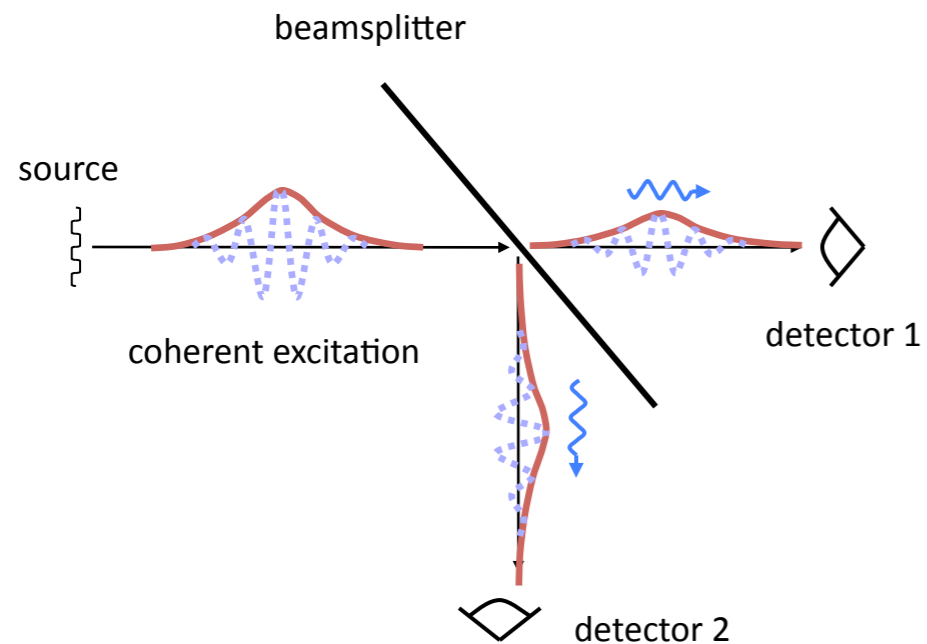


HBT correlation in a ballistic conductor

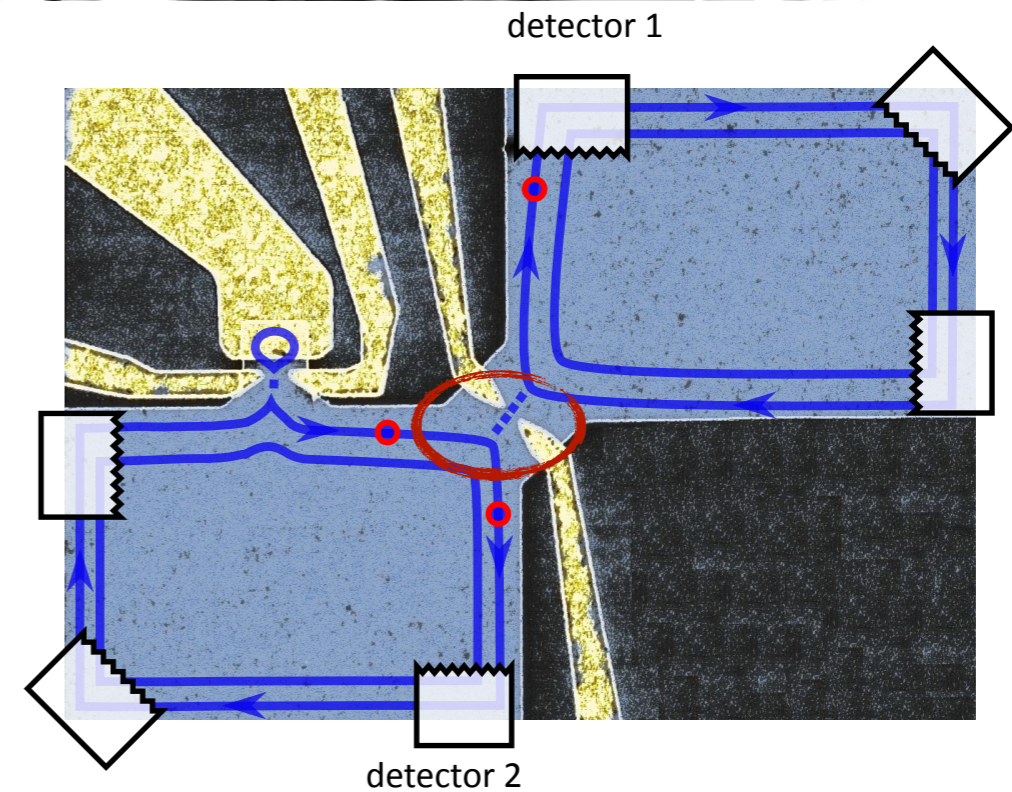
- ballistic propagation in vacuum \Leftrightarrow spin polarised 1-D edge channel of IQHE
- beamsplitter \Leftrightarrow quantum point contact (QPC)

Most experiments : DC sources
 \neq AC single electron source

W. Oliver *et al.*, *Science* **284**, (5412), 299-301, (1999)
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HBT correlations in optics

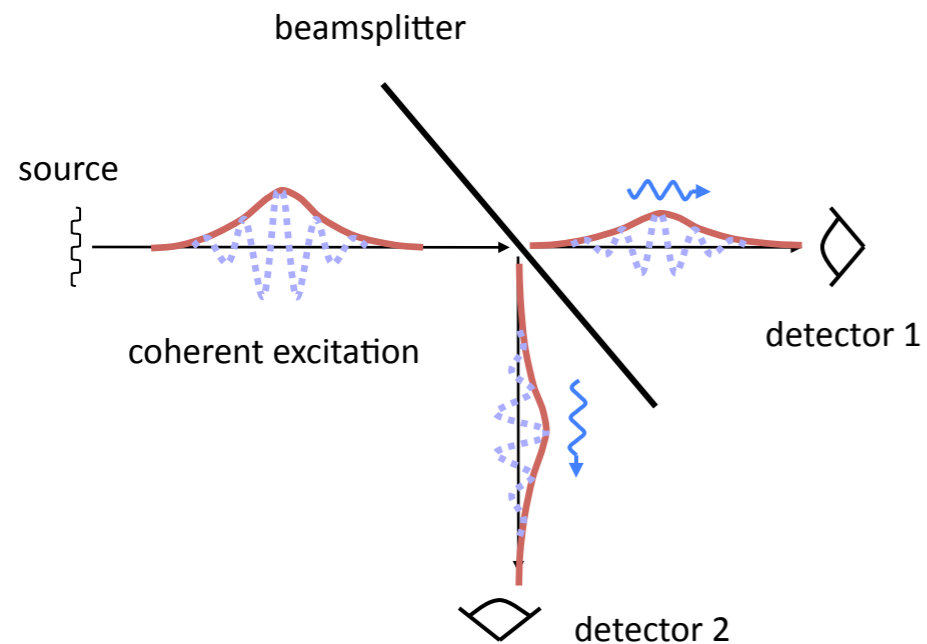


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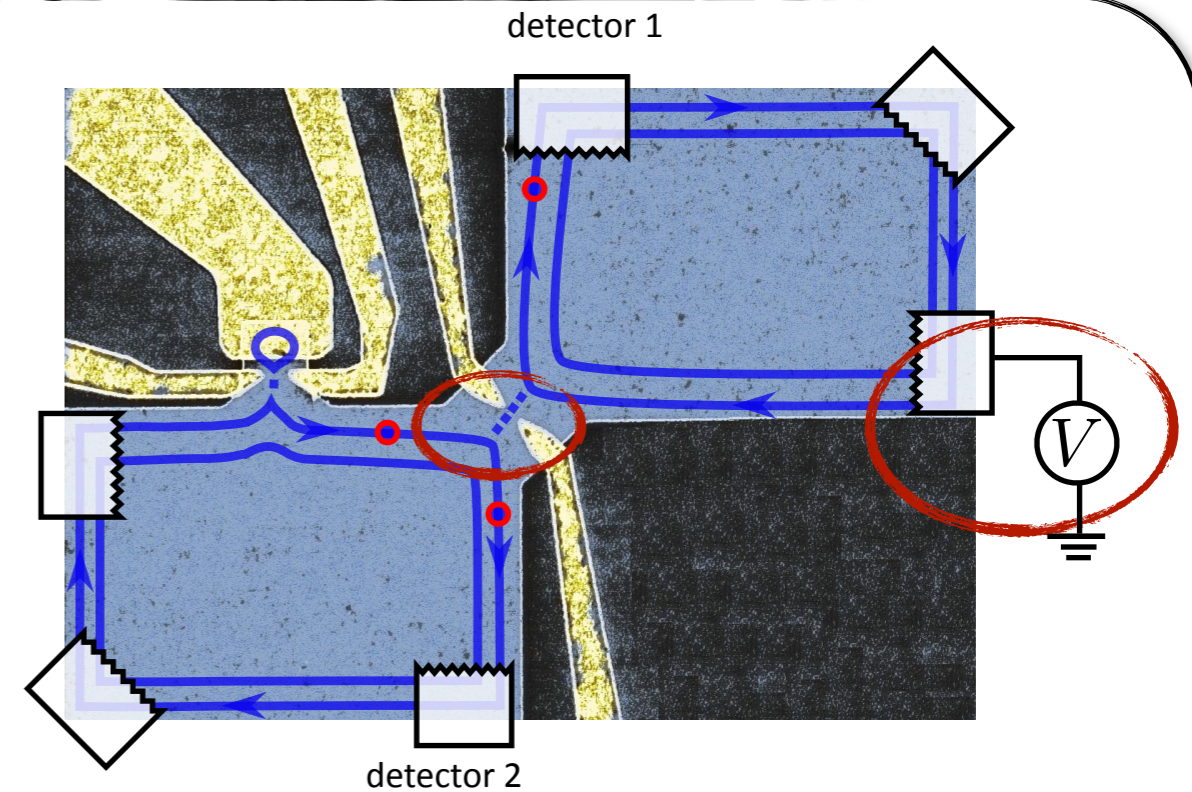
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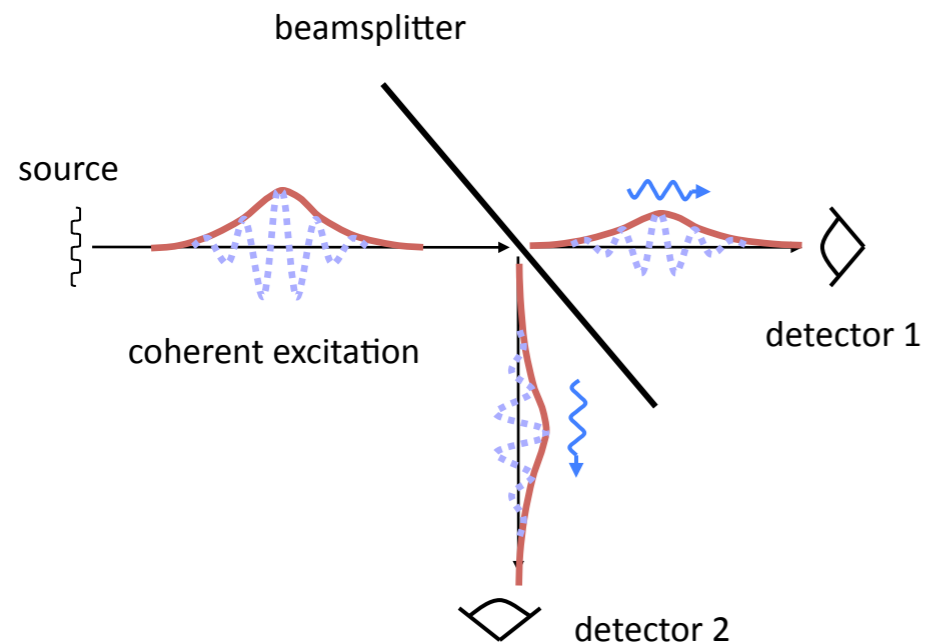


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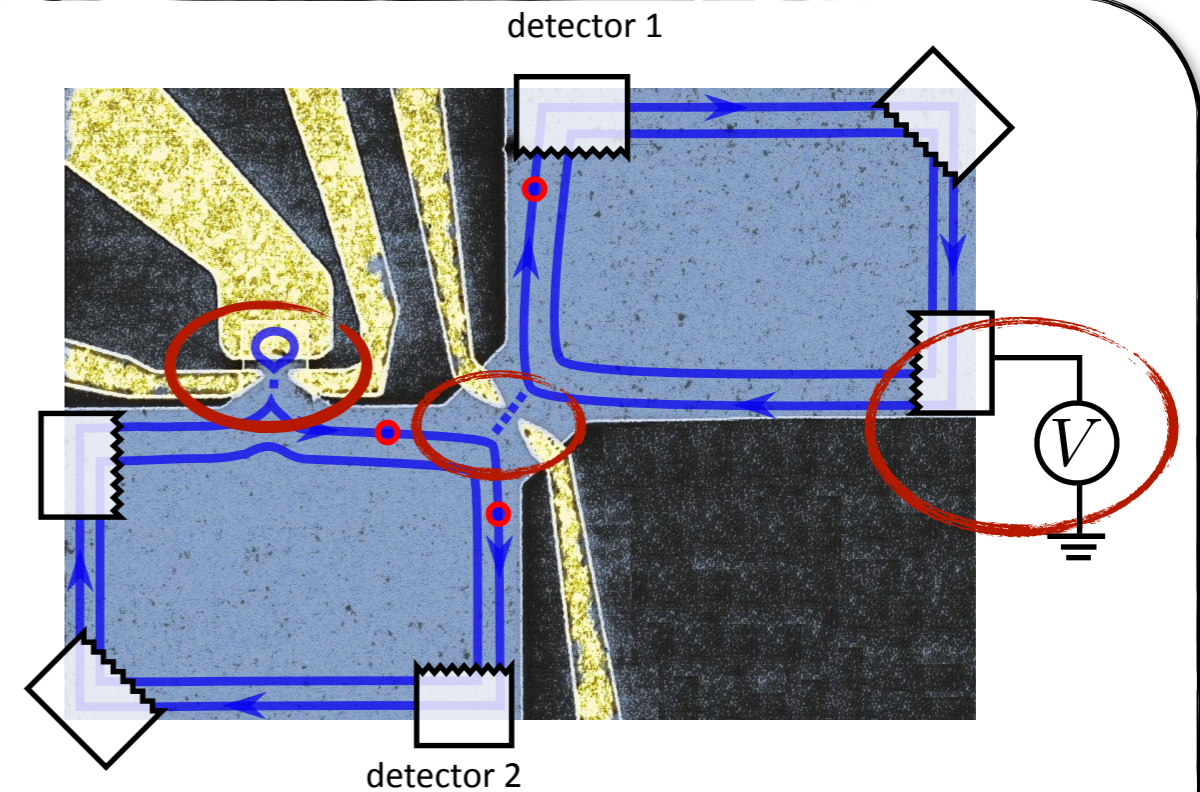
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HBT correlations in optics



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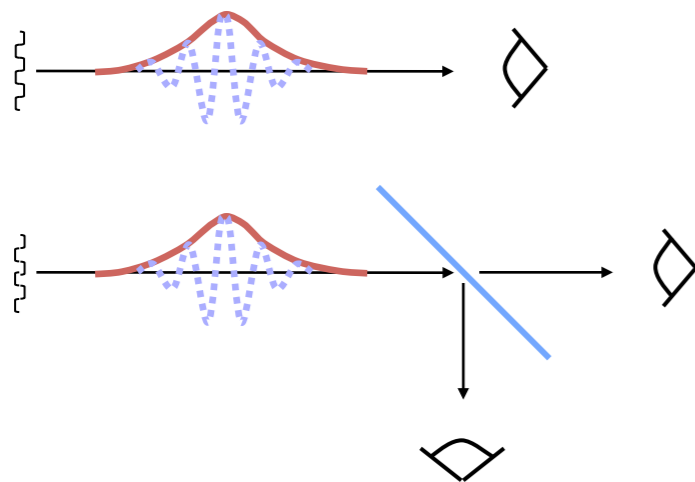
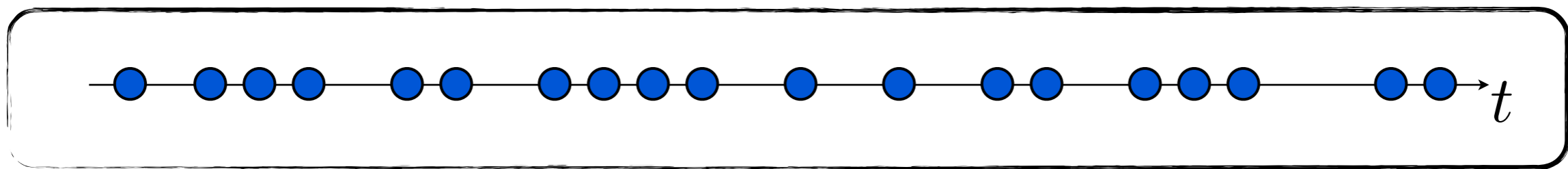
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Source with 1 type of carrier :

- photon source
- biased contact at T=0

$$\text{Average current : } \langle \hat{I} \rangle = e \frac{\langle N_e \rangle}{T_{meas}}$$



Autocorrelations

HBT correlations

$$S(\omega = 0) \text{ gives } \langle \delta N_e^2 \rangle$$

M. Reznikov *et al.*, PRL **75**, 3340 (1995)

A. Kumar *et al.*, PRL **76**, 2778 (1996)

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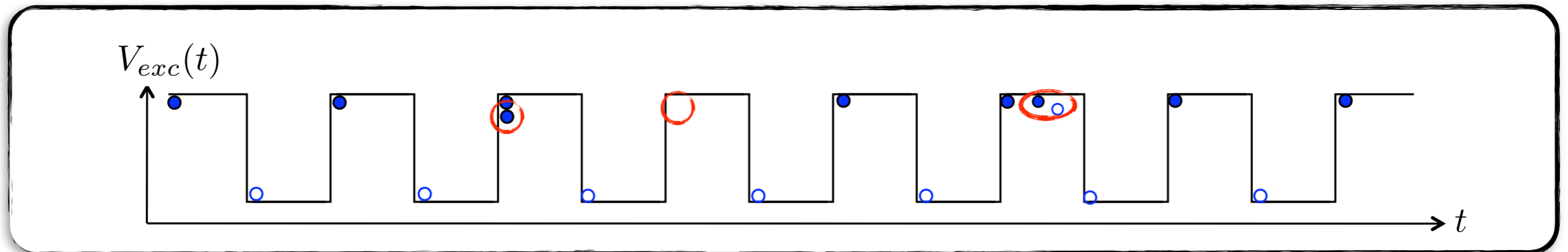
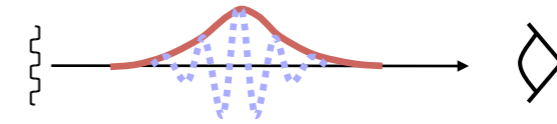
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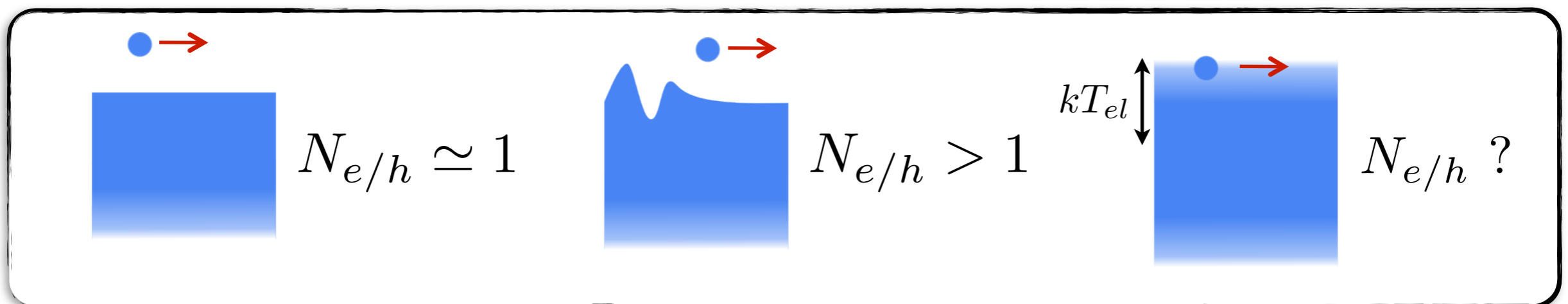
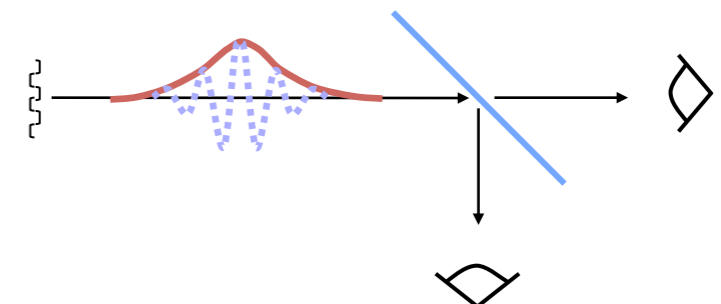
T. Martin *et al.*, PRB **45**(4), 1742-1755 (1992)

Y. Blanter *et al.*, Physics Reports, **336**(1-2), 1-166 (2000)...

Autocorrelations : Statistics of emitted charge
 charge counted in **high-frequency noise**
 demonstrates **single charge emission**



HBT correlations : Counting of emitted quasiparticles
 each particle counted in **low-frequency noise**
 reveal **extra e/h pairs** and effect of **thermal excitations**

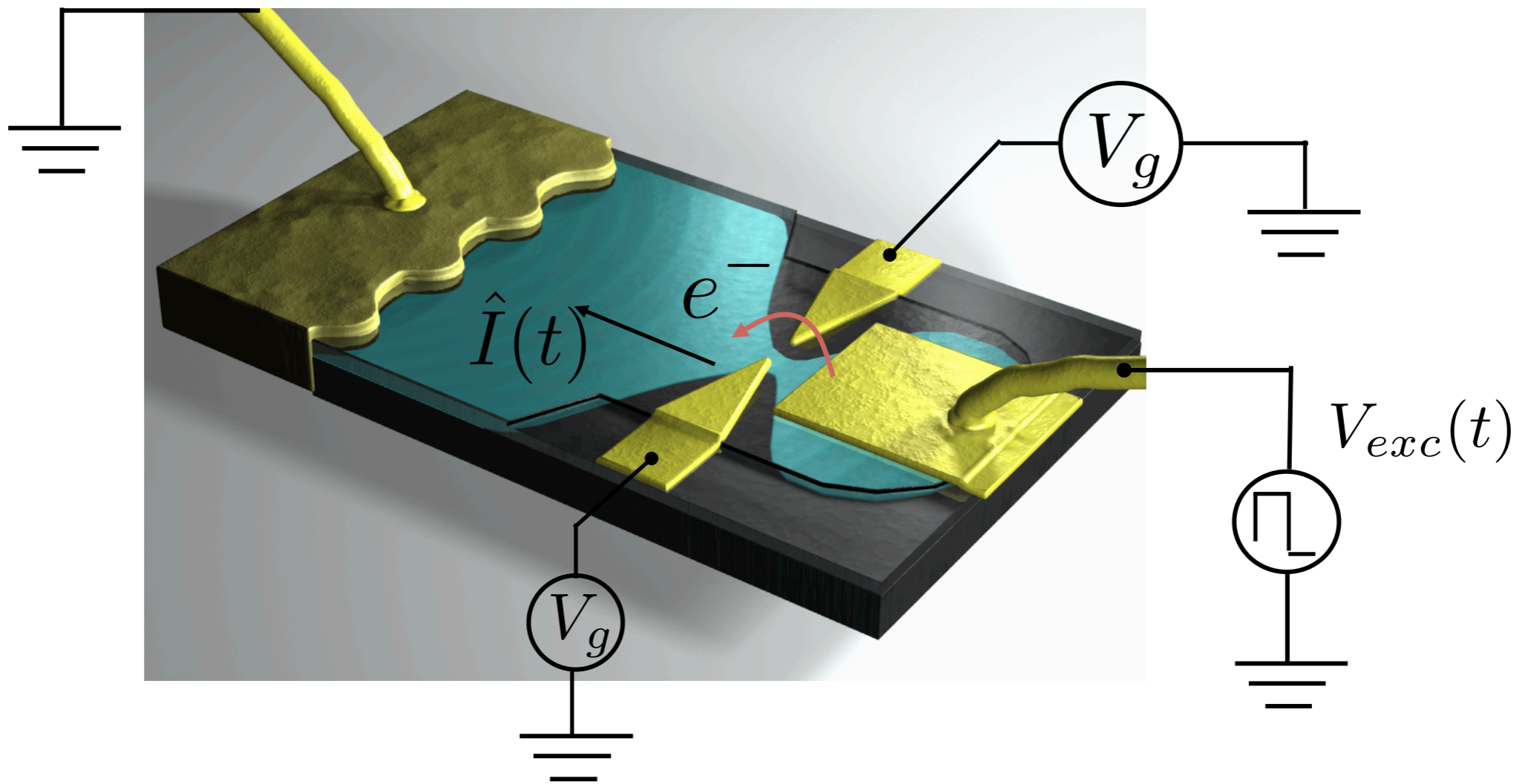


Short time autocorrelations of the current

- single electron source
- high-frequency noise measurement
- phase noise : evidence of single charge emission

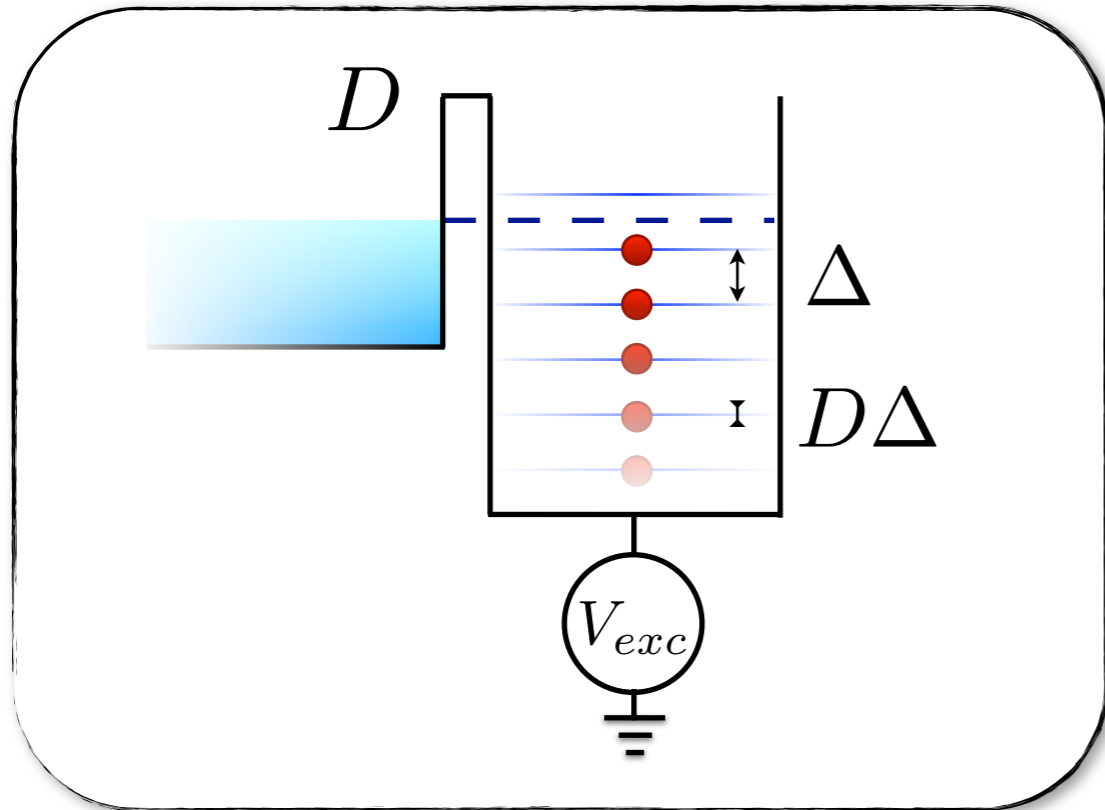
Hanbury-Brown & Twiss experiment with single electrons

- counting quasiparticles with partition noise
- effect of the Fermi sea
- engineering of the wavepacket



M. Büttiker *et al.*, Phys. Lett. A **180**, 364 (1993)

J. Gabelli *et al.*, Science **313**, 499 (2006)

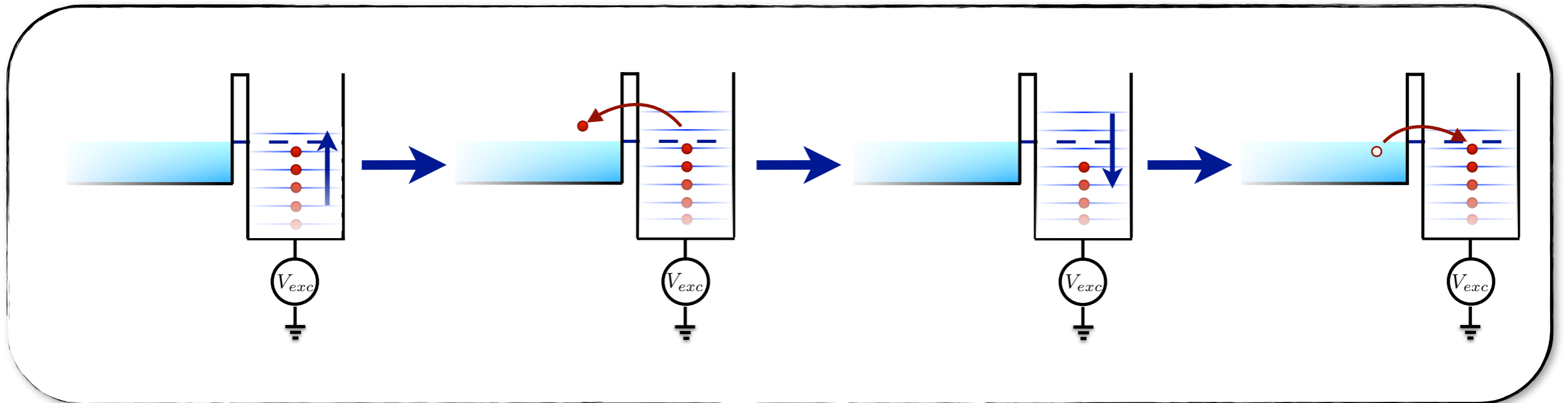


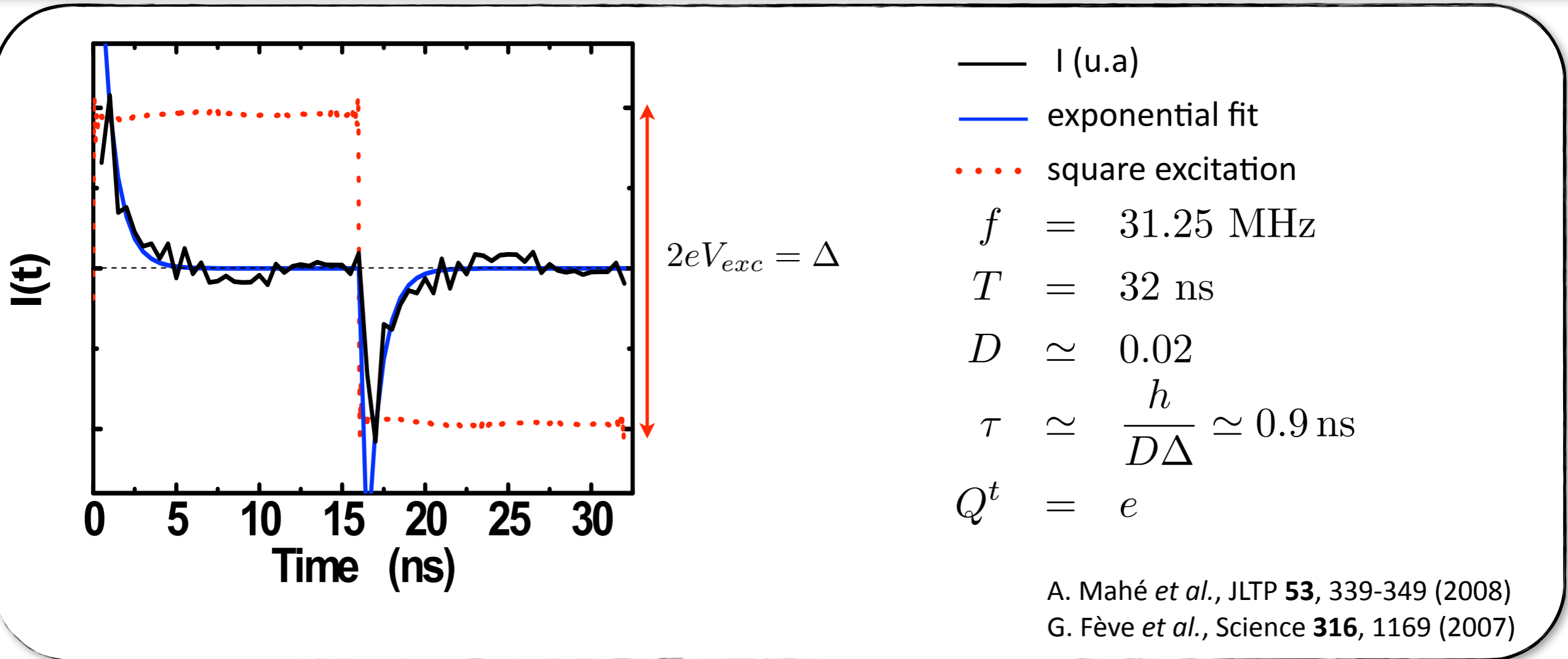
level spacing : $\Delta \simeq 2.1K$

level broadening : $D\Delta$

escape time : $\tau \simeq \frac{h}{D\Delta}$

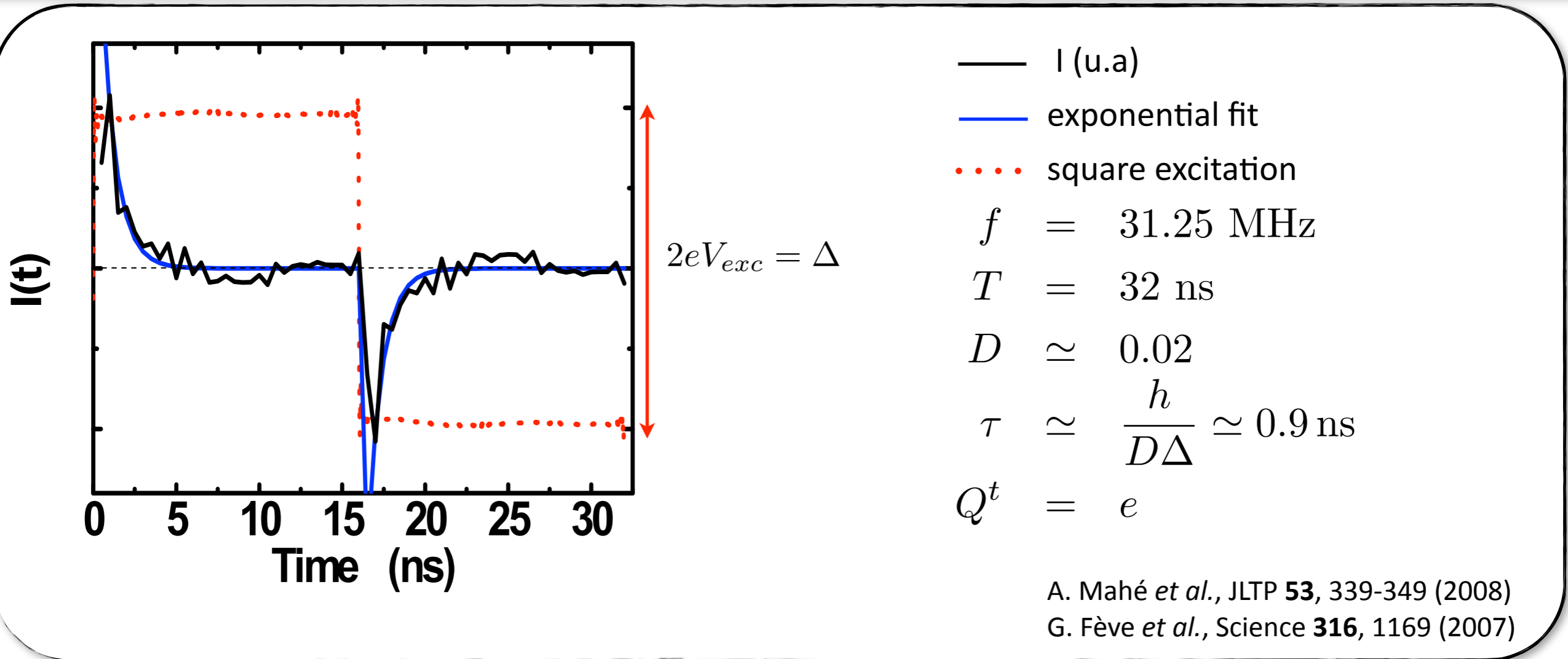
driving frequency : $f = \frac{1}{\tau} \simeq 1.7 \text{ GHz}$





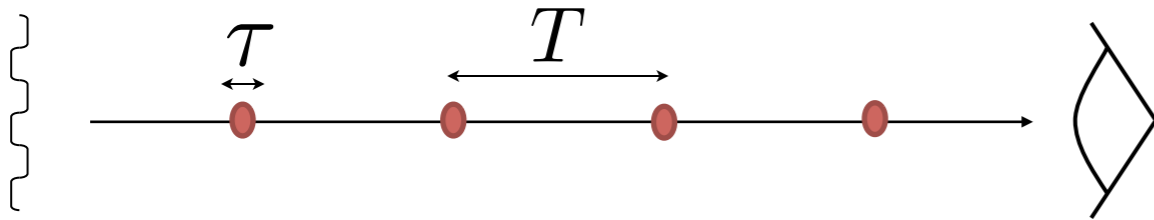
$$\left. \begin{array}{l} \tau \ll T/2 \\ \Rightarrow Q^t = e \end{array} \right\} \Rightarrow \text{single charge emission on average}$$

$$\left. \begin{array}{l} \tau \geq T/2 \\ \Rightarrow Q^t < e \end{array} \right\} \Rightarrow \text{charges not emitted during each cycle}$$

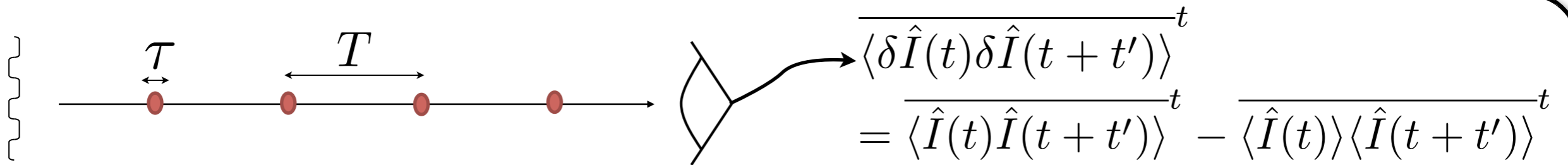


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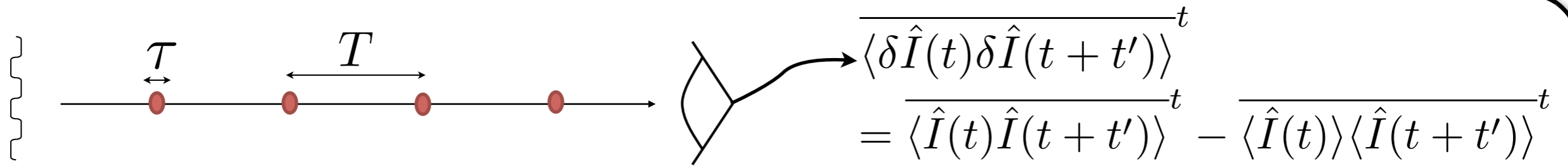
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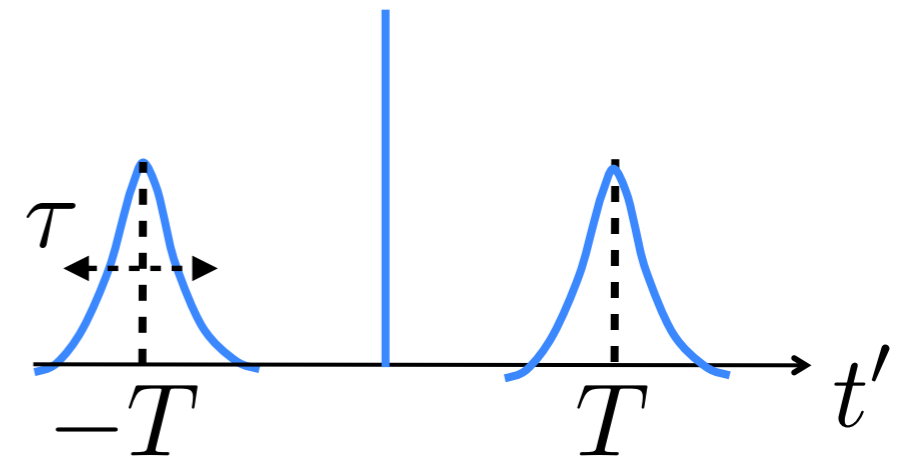
F.D. Parmentier *et al.*, RSI **82**(1), 013904 (2011)
 A. Mahé *et al.*, PRB **82**(20), 201309 (2010)
 M. Albert *et al.*, PRB **82**(4), 41407 (2010)



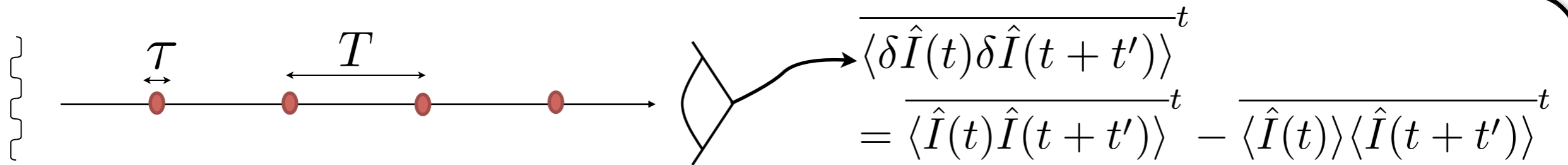
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for $t' < T$ — $\overline{\langle \hat{I}(t) \hat{I}(t + t') \rangle}^t = \frac{e^2}{T} \delta(t')$



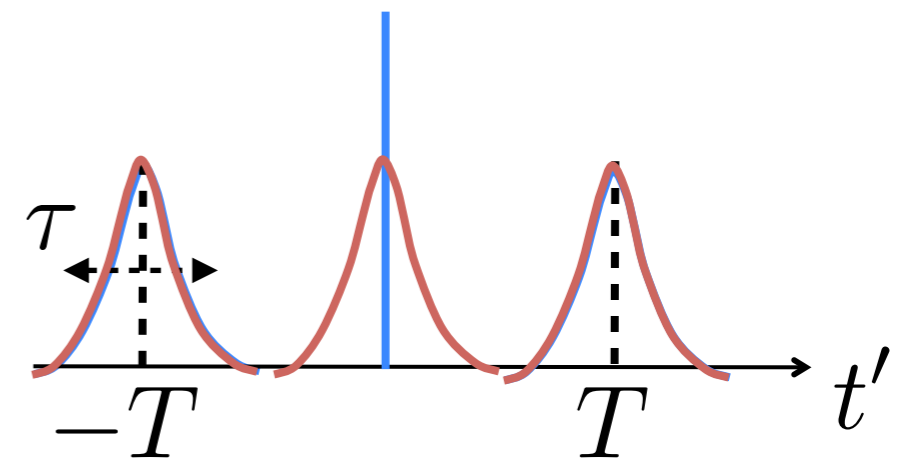
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for $t' < T$

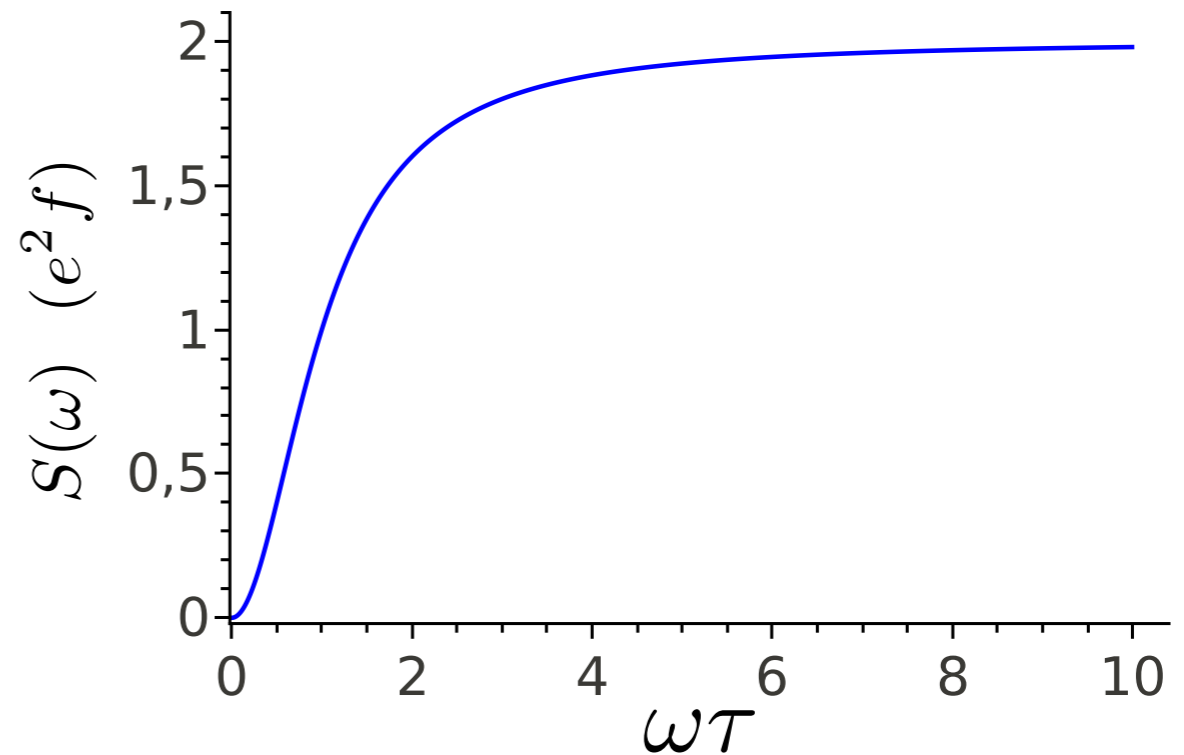
$$- \overline{\langle \hat{I}(t) \hat{I}(t + t') \rangle}^t = \frac{e^2}{T} \delta(t')$$

$$- \overline{\langle \hat{I}(t) \rangle} \overline{\langle \hat{I}(t + t') \rangle}^t = \frac{e^2}{T} \frac{e^{-|t'|/\tau}}{\tau}$$



F.D. Parmentier *et al*, RSI **82**(1), 013904 (2011)
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$$S(\omega) = 2e^2 f \frac{(\omega\tau)^2}{1 + (\omega\tau)^2}$$



Fundamental noise : phase noise

High frequency noise : $\omega\tau \simeq 1$

τ -dependence probed, at fixed $\omega \simeq 2\pi f$

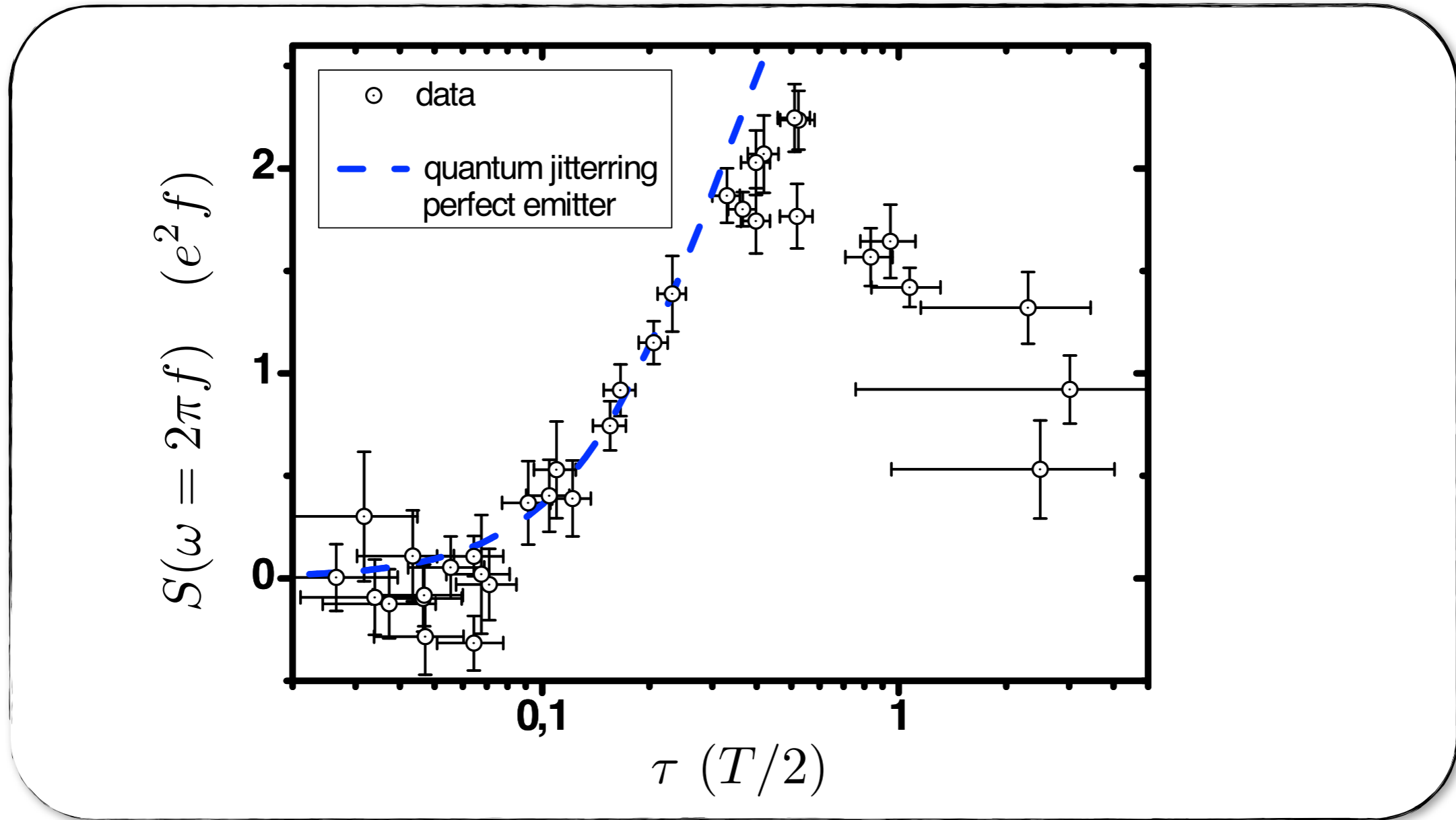
$$S(\omega) \simeq e^2 f \simeq 4 \cdot 10^{-29} \text{ A}^2 \text{ Hz}^{-1}$$

at $f = 1.5 \text{ GHz}$

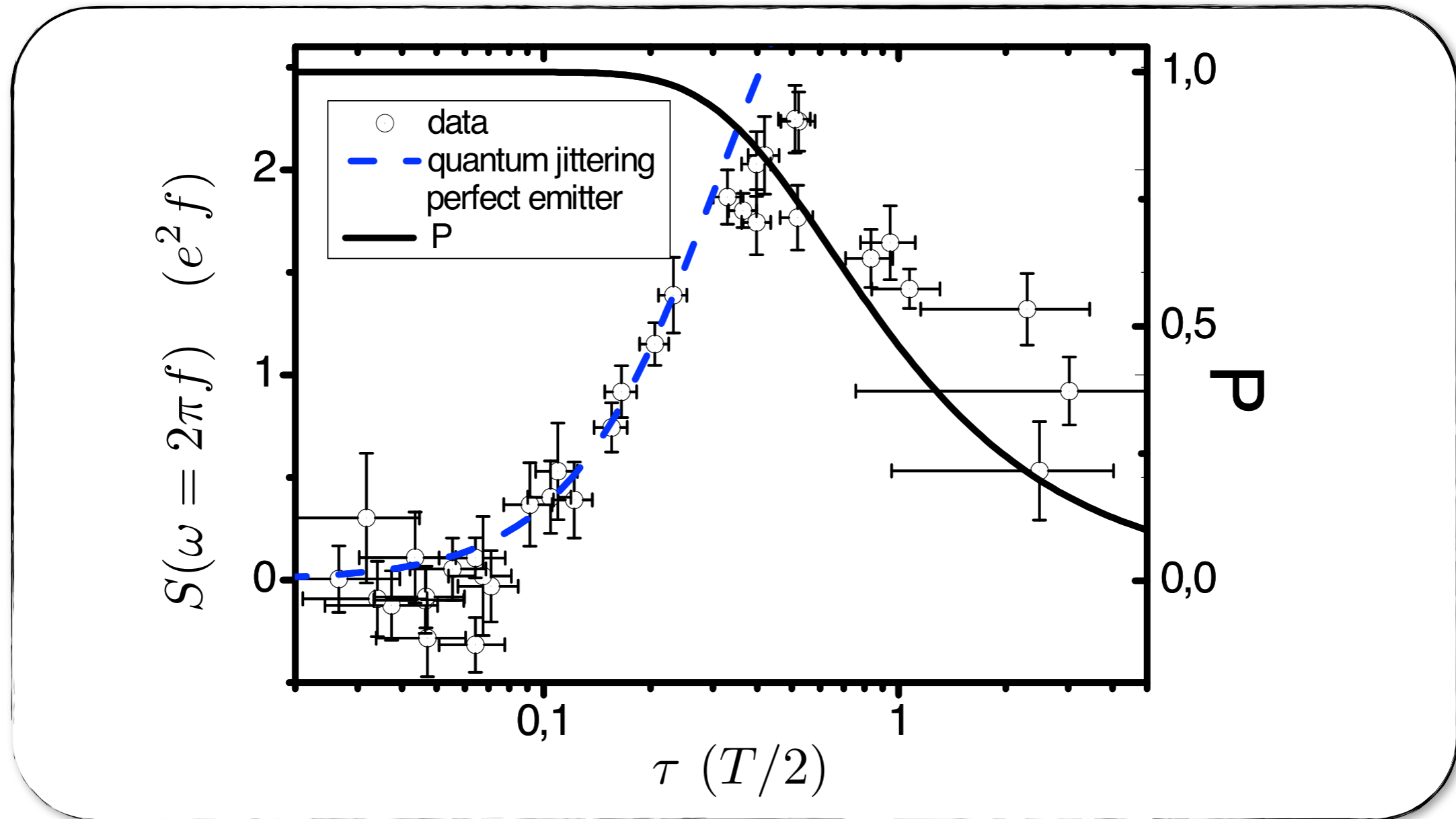
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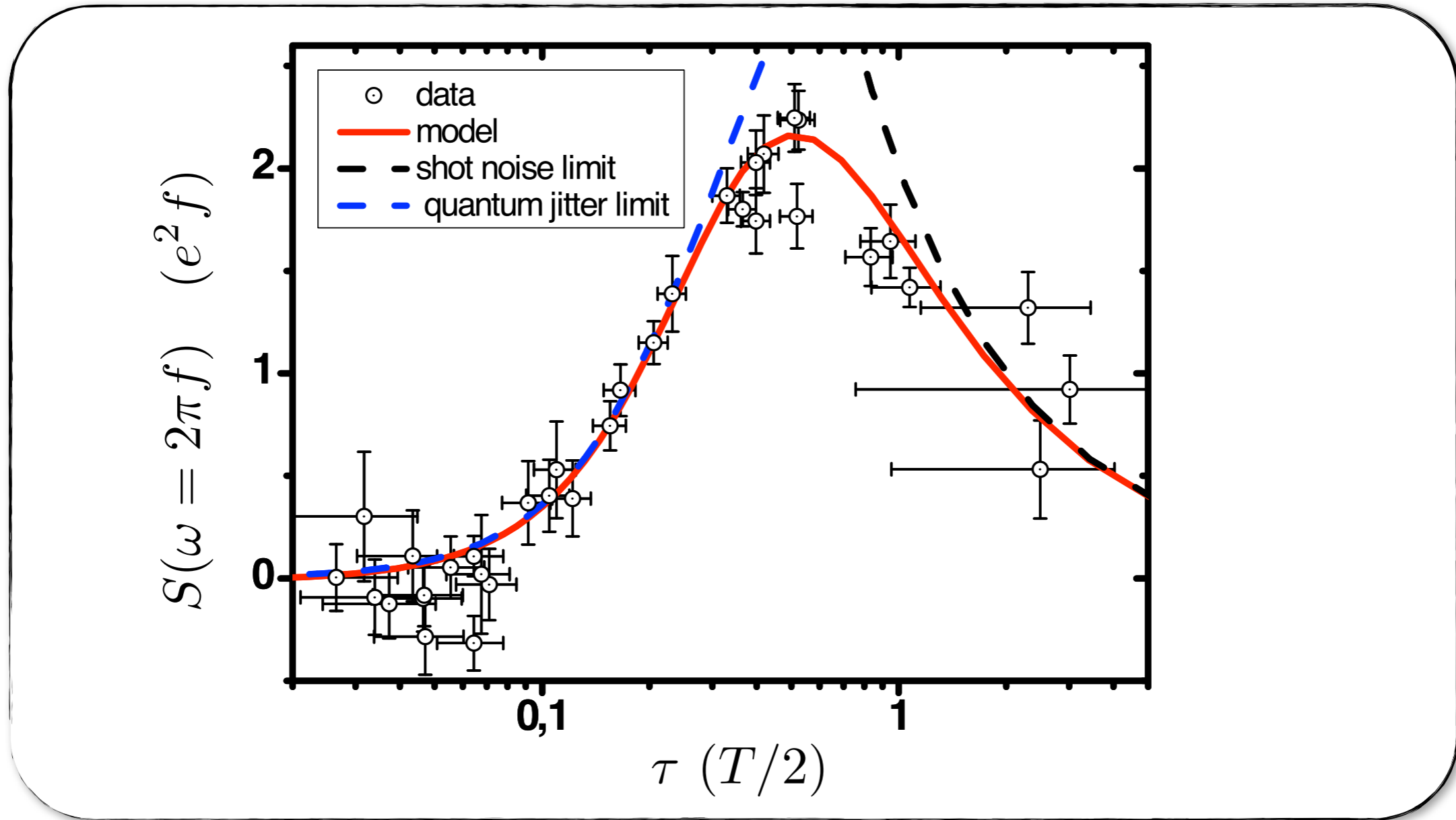
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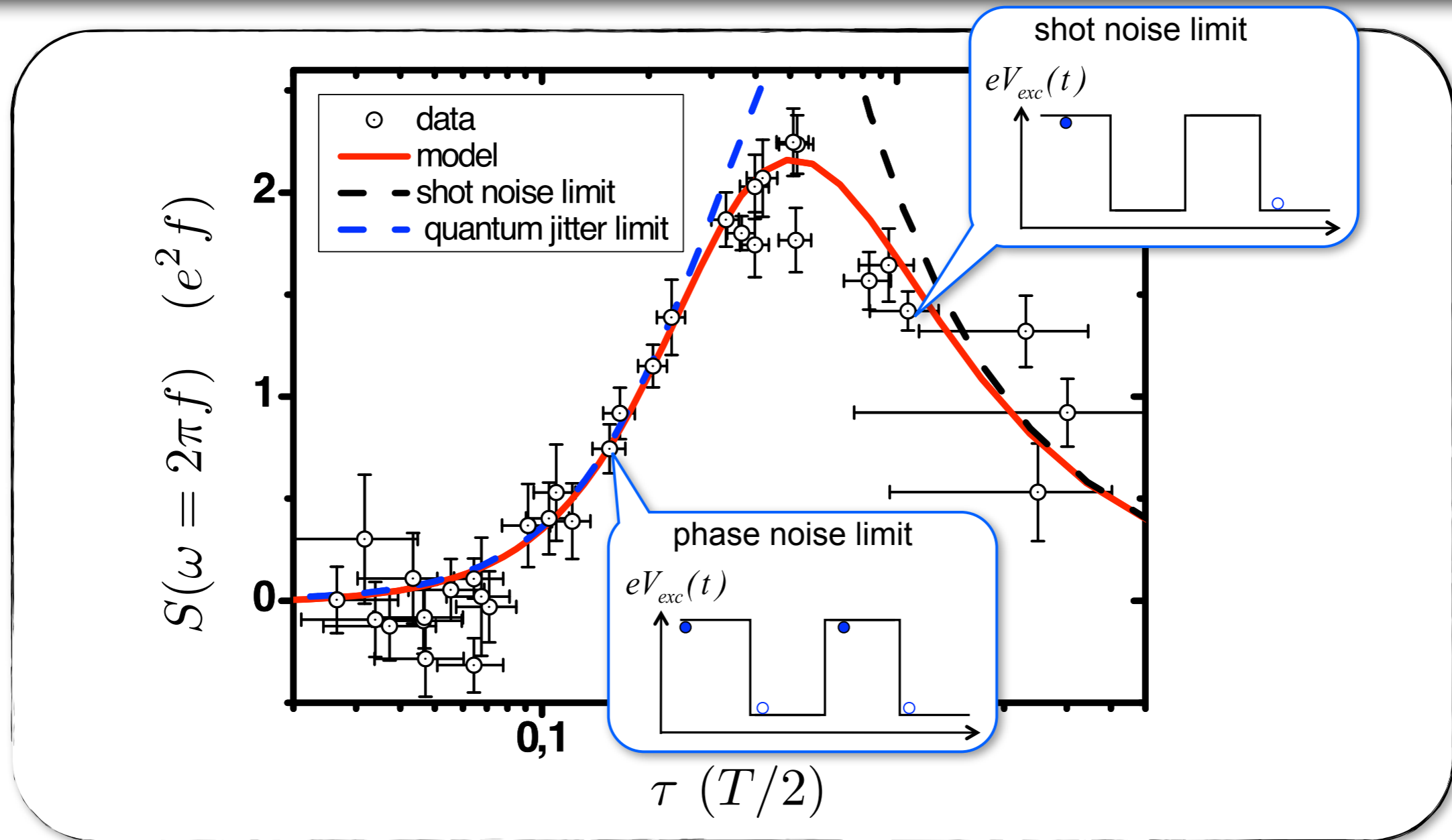
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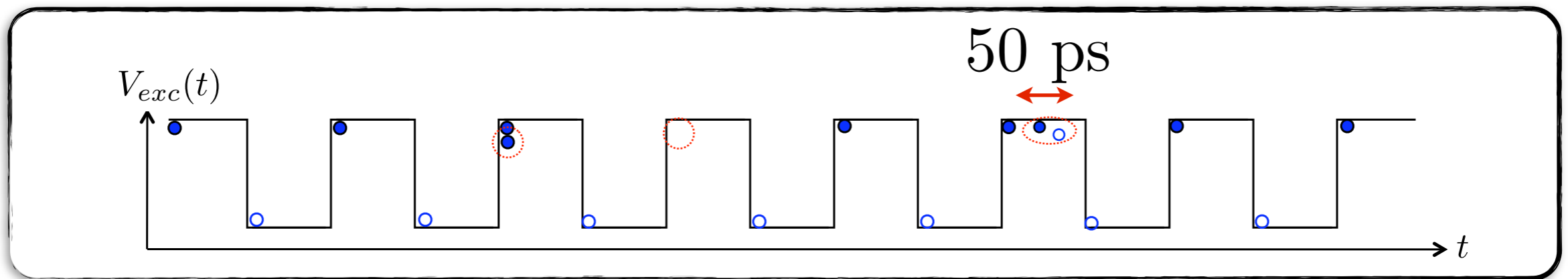
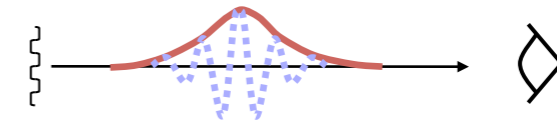
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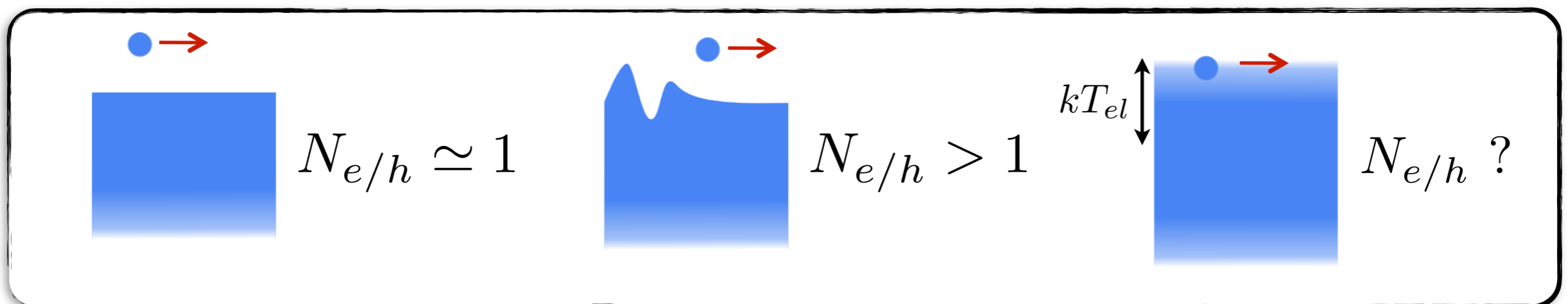
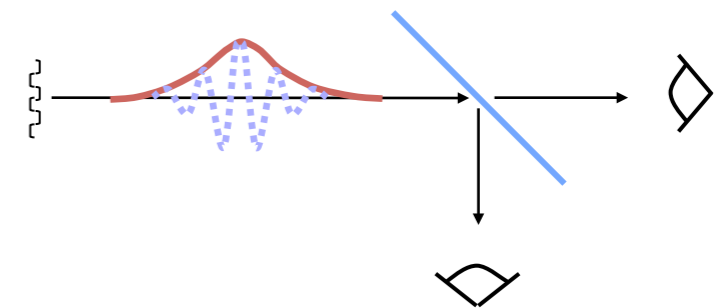
Noise reduced to phase noise \Rightarrow demonstration of **single charge emission** !
 Fully described by 2 models : Floquet scattering theory & heuristic model

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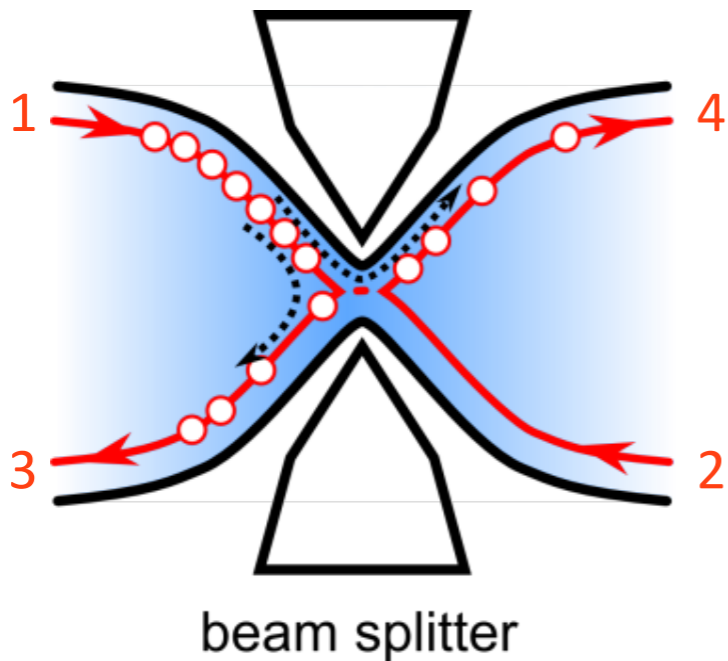
Autocorrelations : Statistics of emitted charge



HBT correlations : Counting of emitted quasiparticles



continuous source

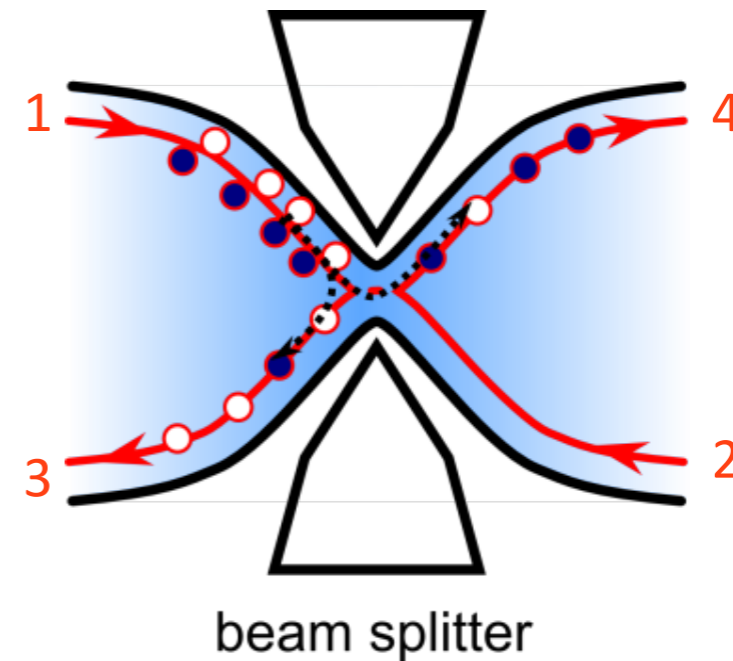


$$S_{3,4}(\omega = 0) = 2e\langle\hat{I}\rangle \times T(1 - T)$$

$$\langle\hat{I}\rangle = e\frac{\langle N_e \rangle}{T_{meas}}$$

- M. Reznikov *et al.*, PRL **75**, 3340 (1995)
- A. Kumar *et al.*, PRL **76**, 2778 (1996)
- G.B. Lesovik, JETP Lett. **70**, 208 (1999)
- T. Martin *et al.*, PRB **45**(4), 1742-1755 (1992)
- Y. Blanter *et al.*, *Physics Reports*, **336**(1-2), 1-166 (2000)

e/h pairs



$$S_{3,4}(\omega = 0) = 2e\langle\hat{I}_{part}\rangle \times T(1 - T)$$

$$\langle\hat{I}_{part}\rangle = e\frac{\langle N_e \rangle + \langle N_h \rangle}{T_{meas}}$$

- L.-H Reydellet *et al.*, PRL **90**, 176803 (2003)
- V. S. Rychkov *et al.*, PRB **72**, 155326 (2005)

Partition noise (in units e^2f) classically yields the **number of excitations** $\langle N_e \rangle + \langle N_h \rangle$

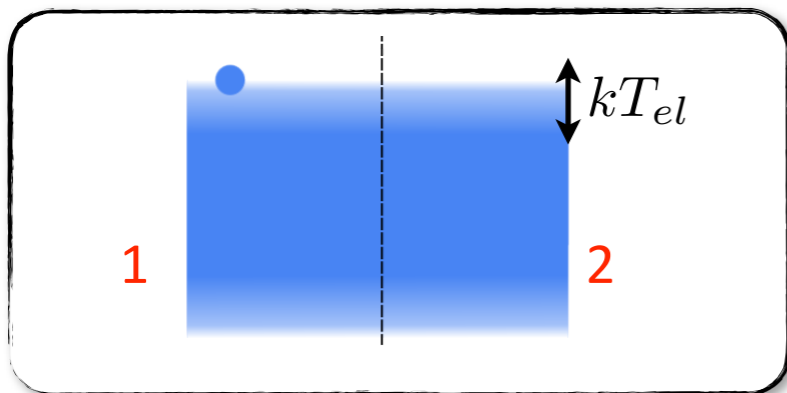
- Autocorrelation on Output 3 :

$$\Delta S_{3,3} = -S_{3,4} = 2e^2 f \times T(1 - T)\Delta Q$$

- Resolution :

$$0.05 e^2 f \simeq 2 \cdot 10^{-30} \text{ A}^2 \text{ Hz}^{-1}$$

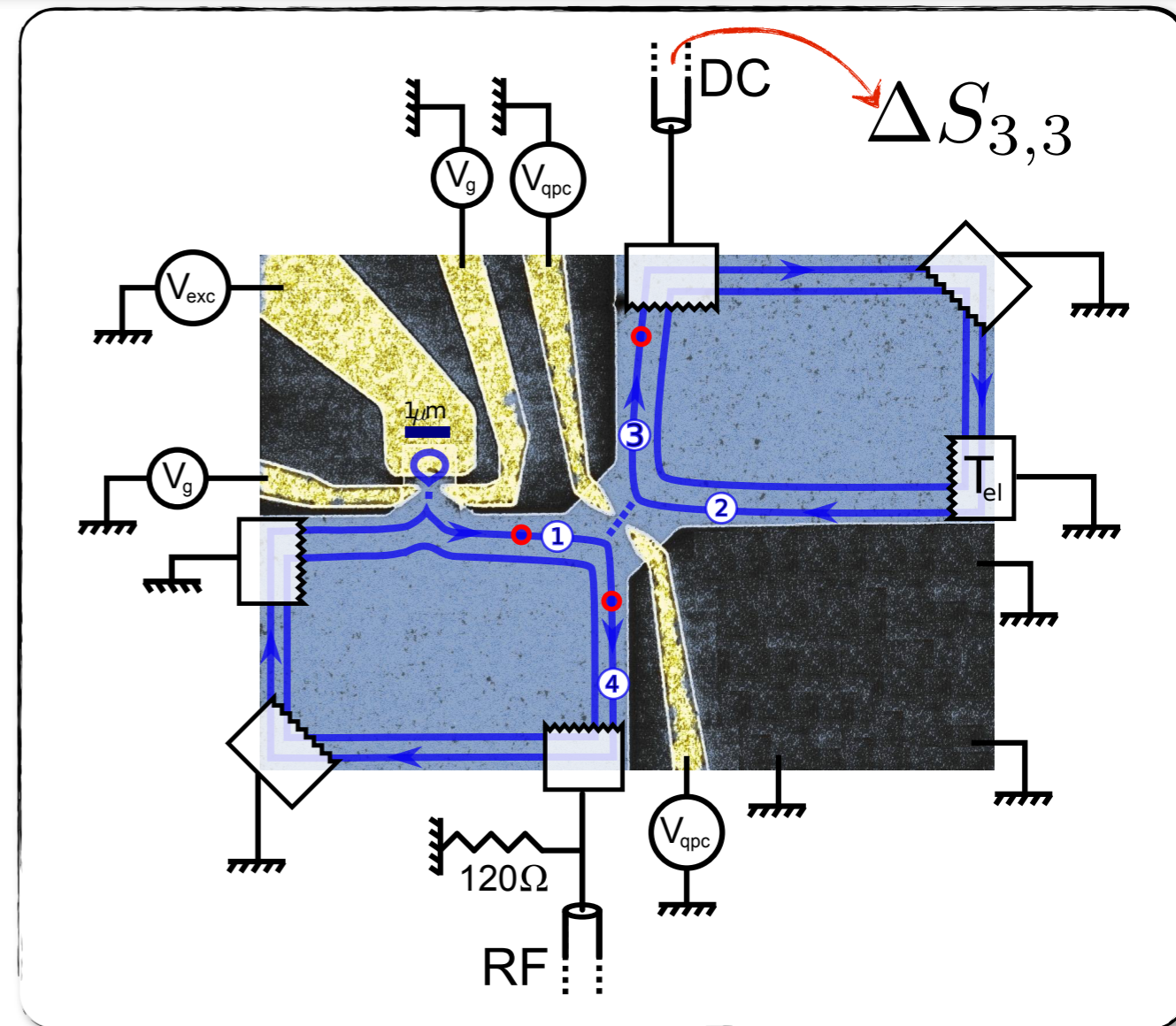
- Input 2 : **Fermi sea at $T_{el} = 150 \text{ mK} \neq \text{vacuum}$!**



- HBT signal :

$$\Delta Q = \langle N_e \rangle + \langle N_h \rangle - 2 \int_0^\infty d\epsilon (n_e(\epsilon) + n_h(\epsilon)) f(\epsilon)$$

$$\langle N_{e/h} \rangle = \int_0^\infty d\epsilon n_{e/h}(\epsilon)$$



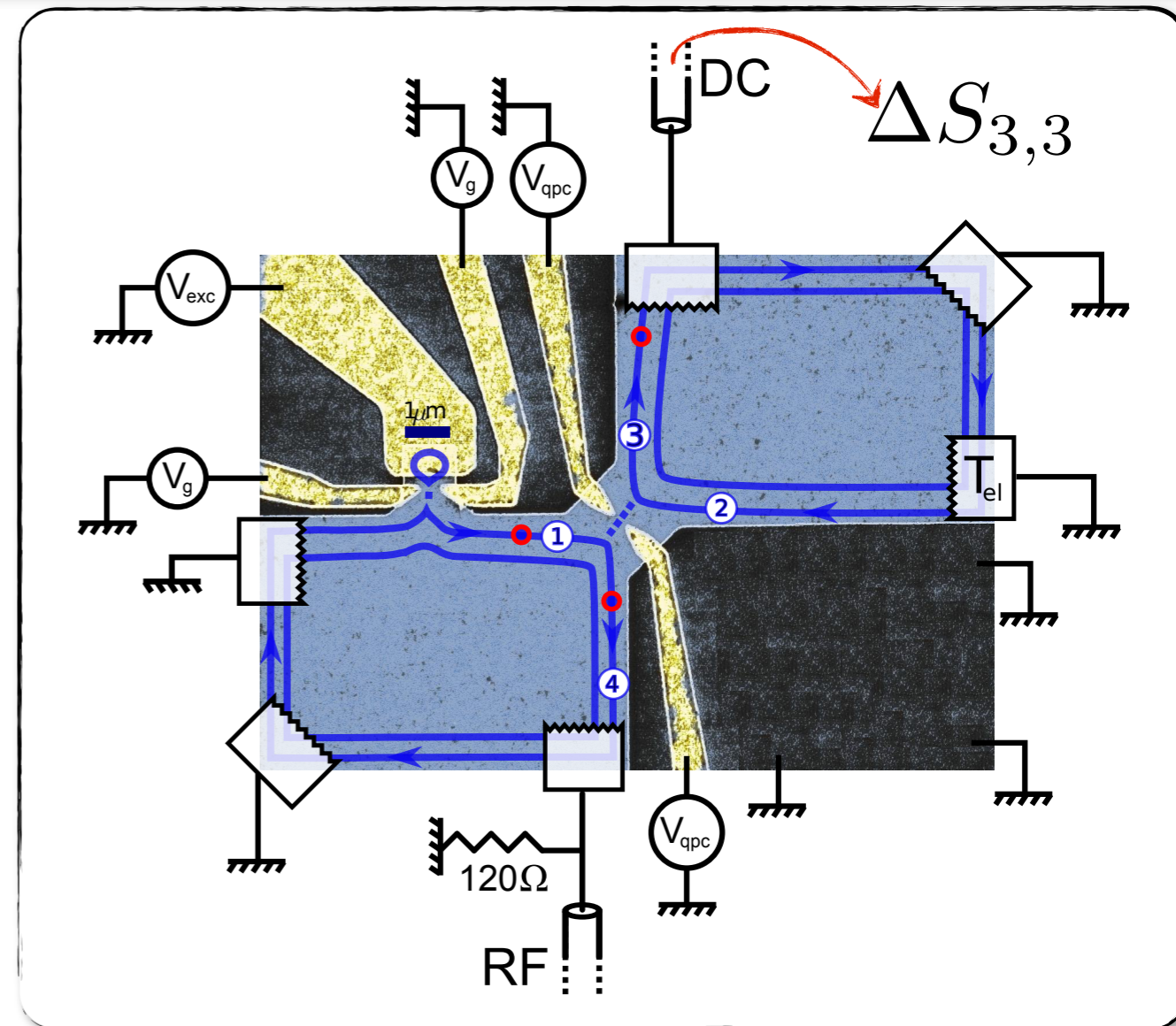
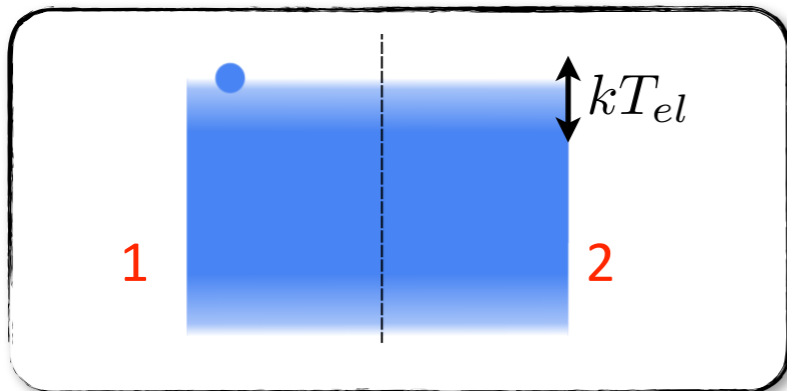
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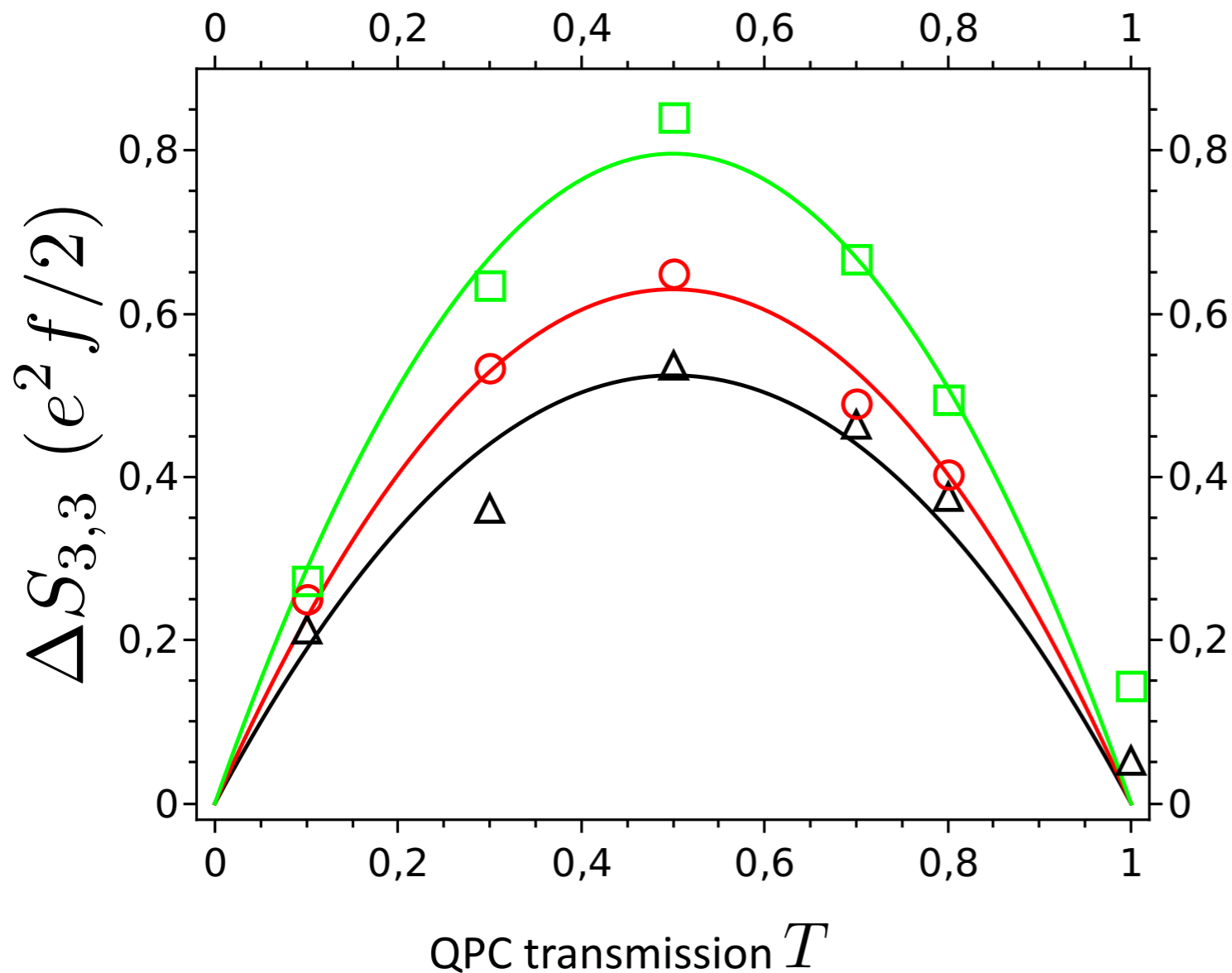
$$\Delta Q = \langle N_e \rangle + \langle N_h \rangle - 2 \int_0^\infty d\epsilon (n_e(\epsilon) + n_h(\epsilon)) f(\epsilon)$$

$$\langle N_{e/h} \rangle = \int_0^\infty d\epsilon n_{e/h}(\epsilon)$$

overlap between Fermi sea and single quasiparticles

$$\Delta S_{3,3} = 2e^2 f T(1 - T)\Delta Q$$

Noise vs QPC transmission



— Sine, $D = 1$

$Q^t = 1.27$ $\Delta Q = 0.51$

— Sine, $D = 0.3$

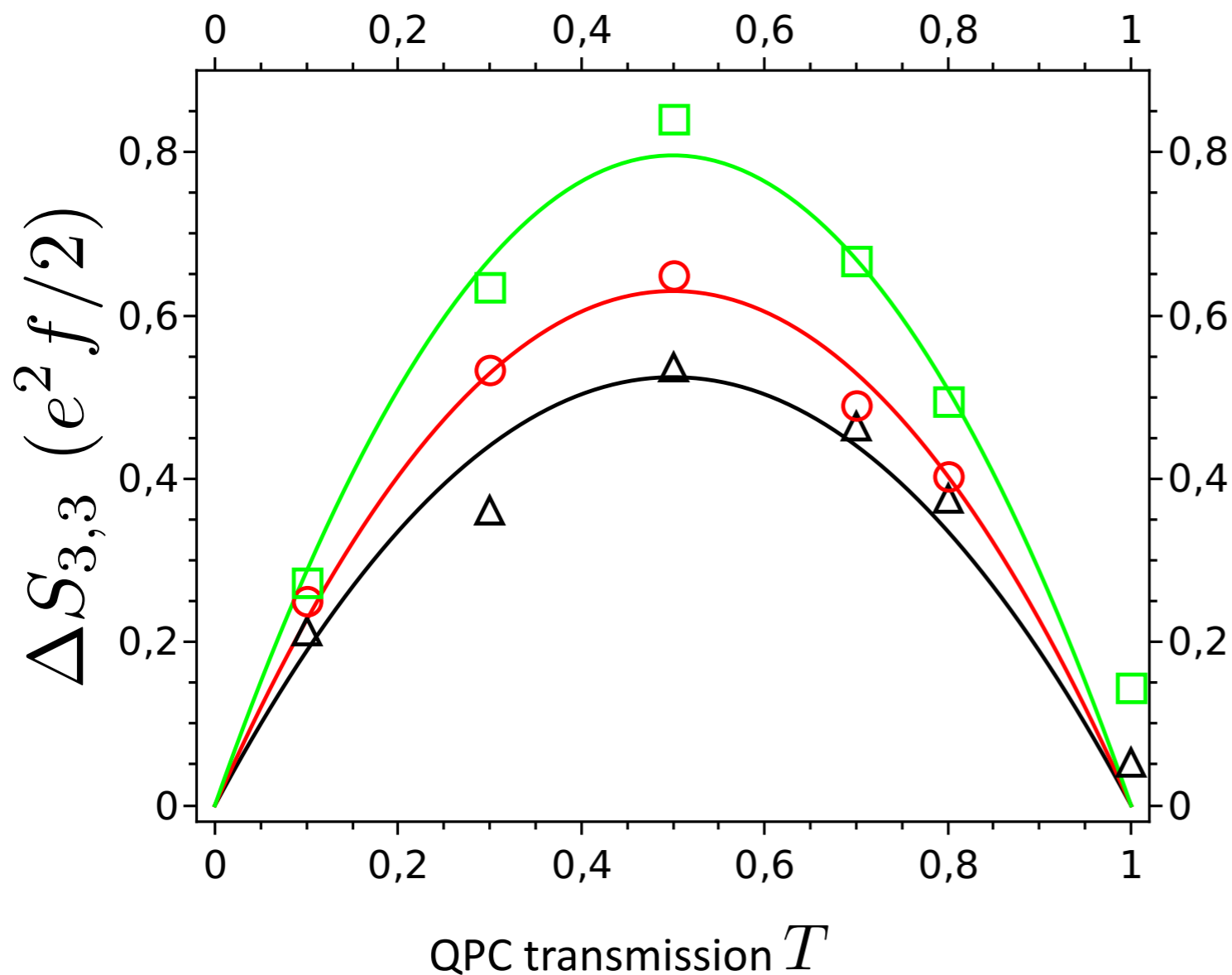
$Q^t = 0.93$ $\Delta Q = 0.63$

— Square, $D = 0.4$

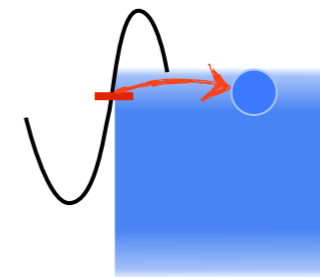
$Q^t = 1$ $\Delta Q = 0.80$

$$\Delta S_{3,3} = 2e^2 f T(1 - T)\Delta Q$$

Noise vs QPC transmission



— Sine, $D = 1$



— Sine, $D = 0.3$

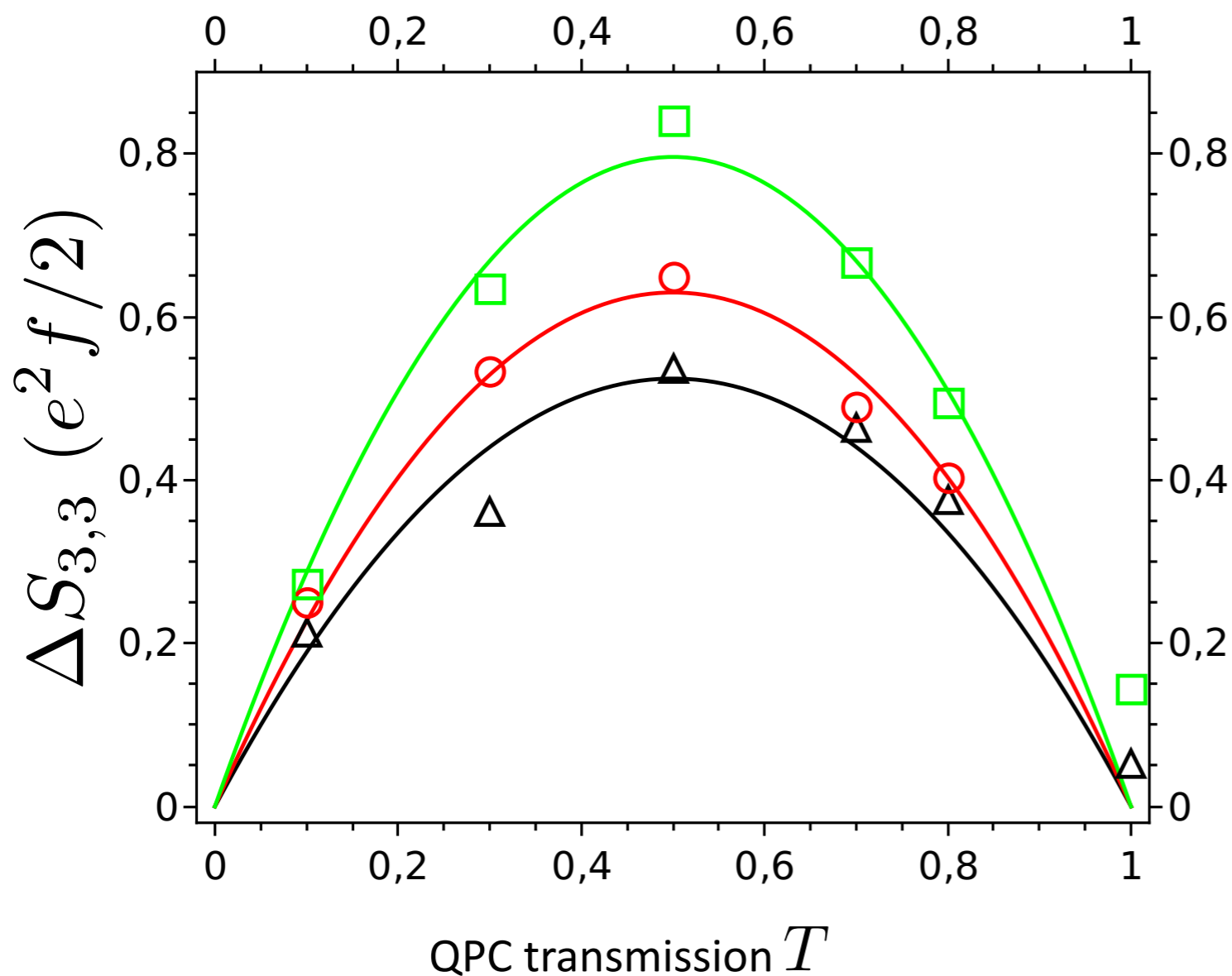
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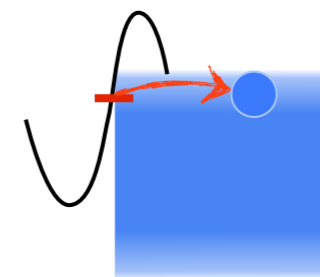
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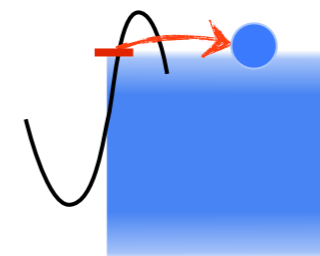
Noise vs QPC transmission



— Sine, $D = 1$



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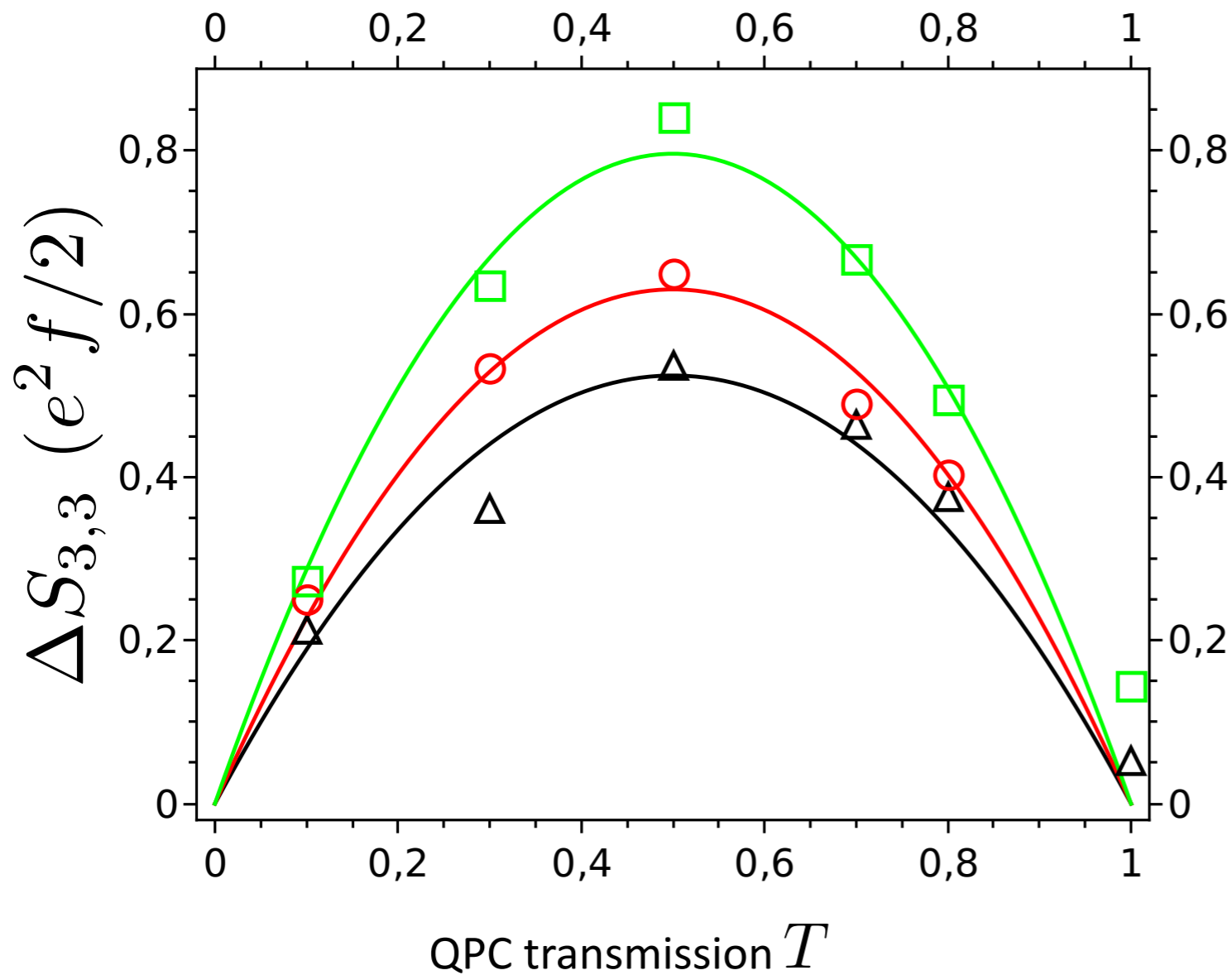


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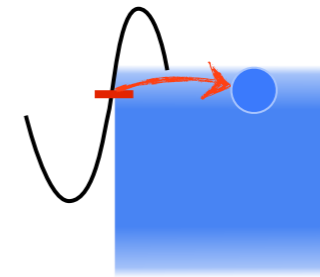
$Q^t = 1$ $\Delta Q = 0.80$

$$\Delta S_{3,3} = 2e^2 f T(1 - T)\Delta Q$$

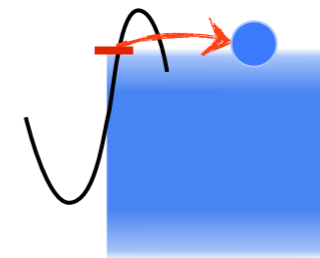
Noise vs QPC transmission



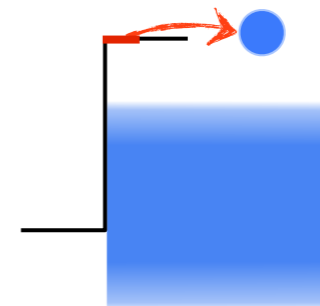
— Sine, $D = 1$

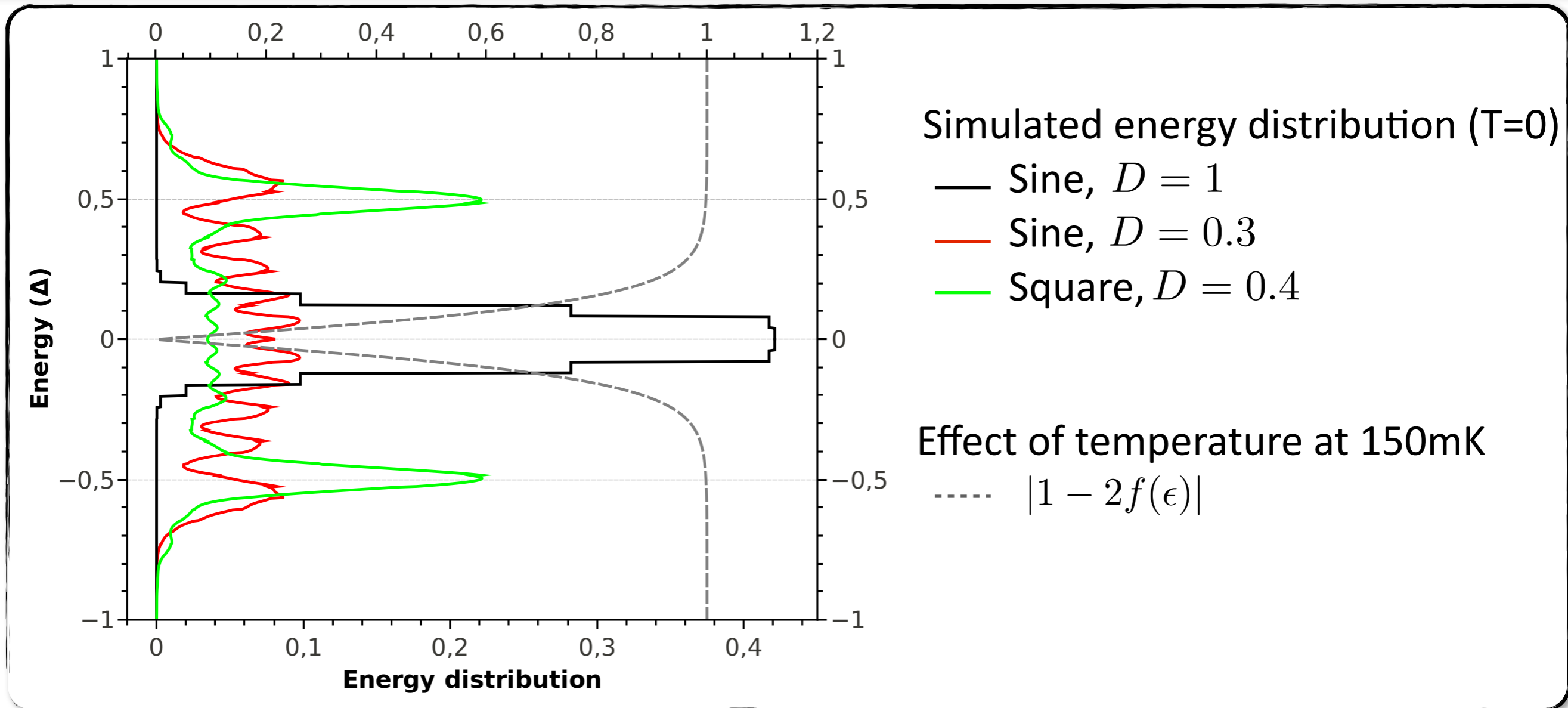


— Sine, $D = 0.3$



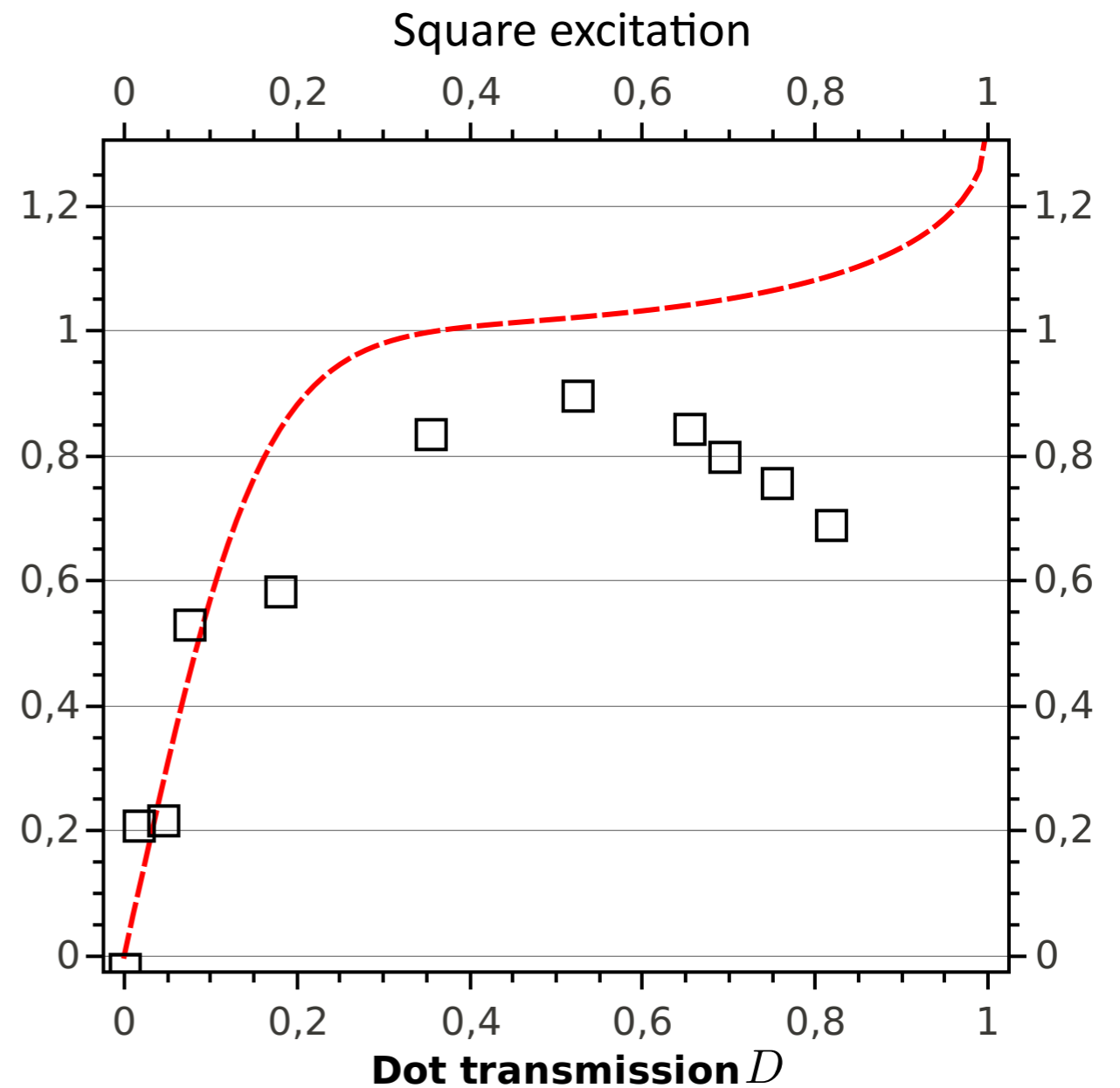
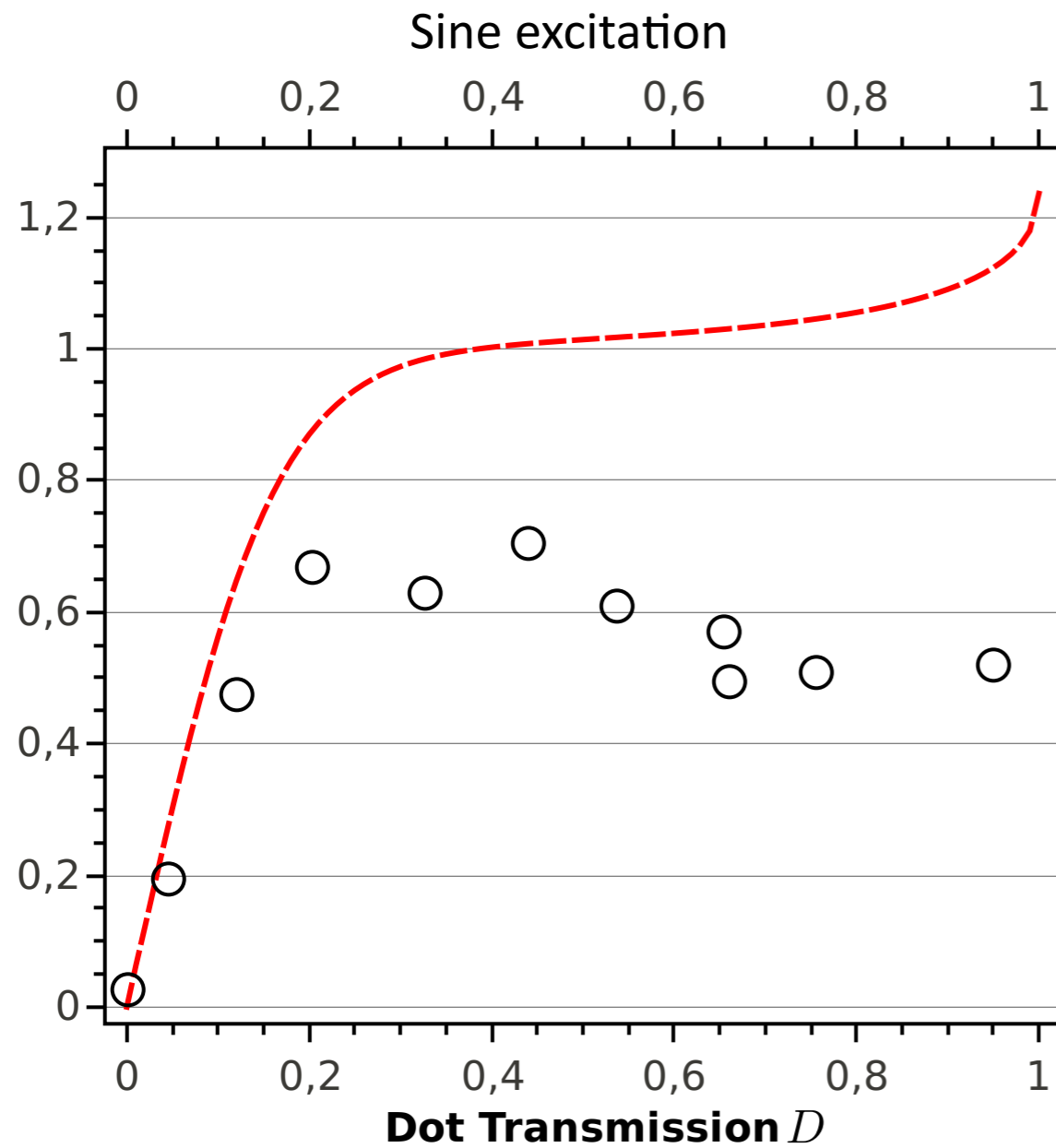
— Square, $D = 0.4$



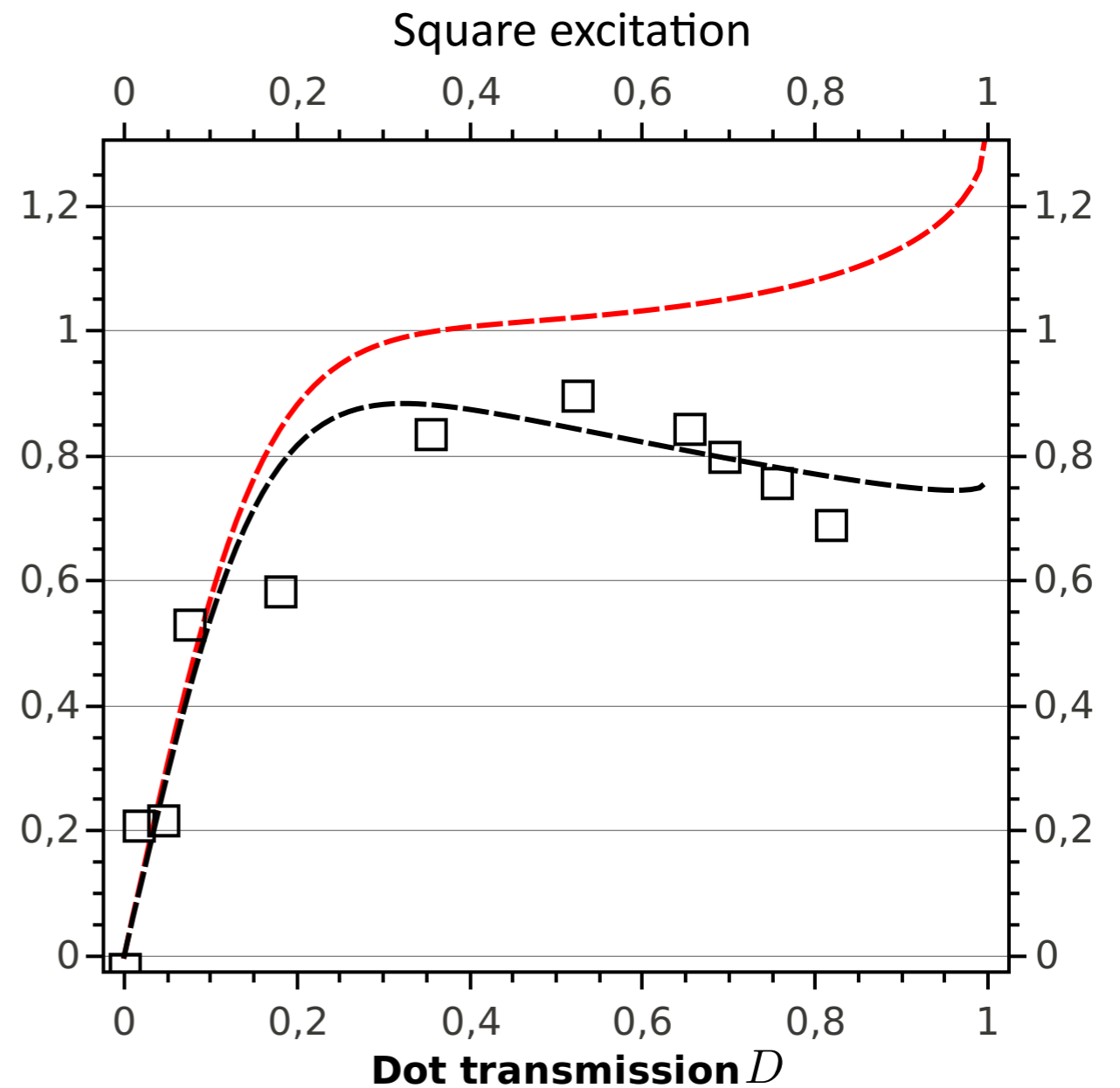
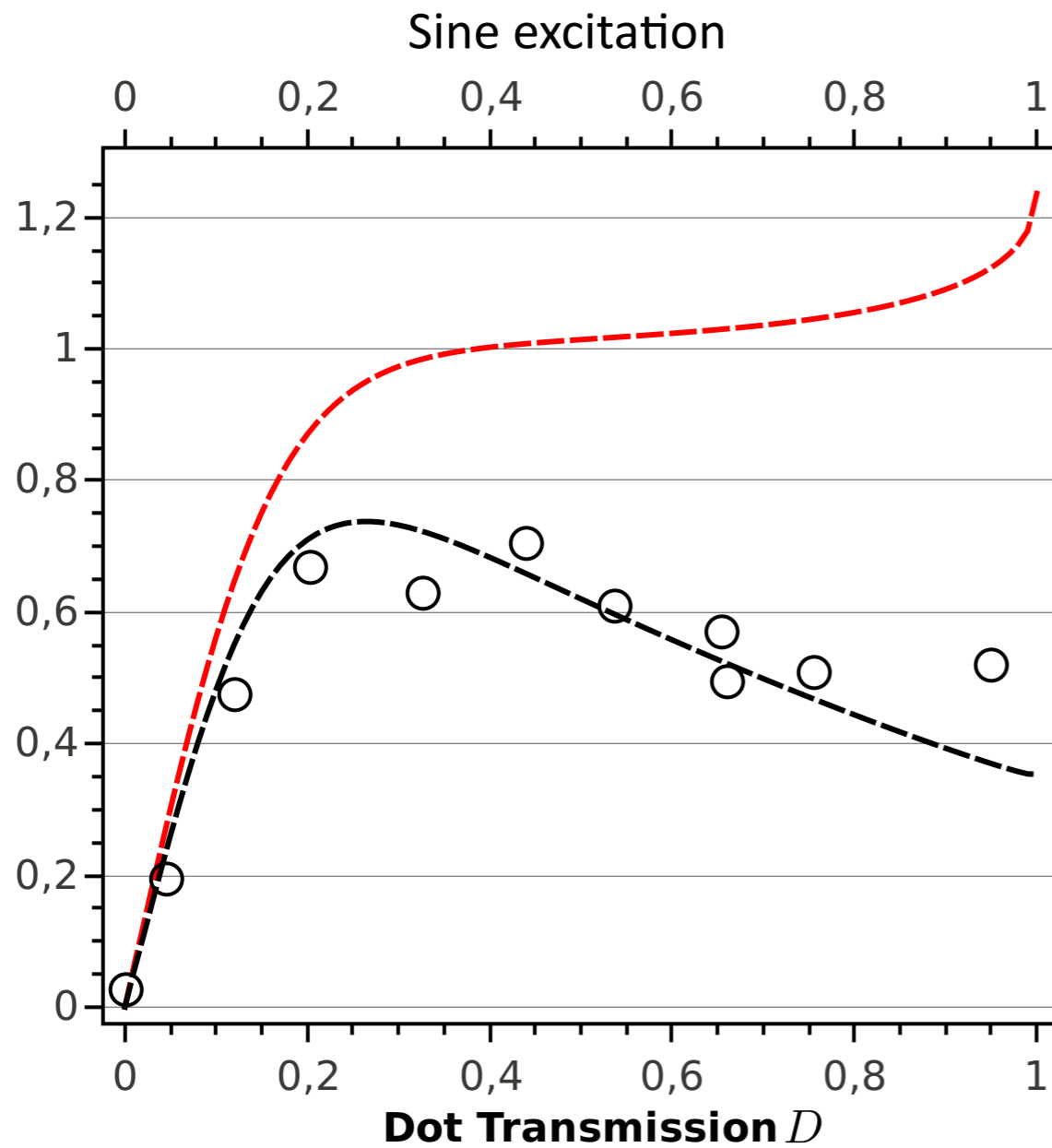


Engineering of the wavepacket through :

- excitation : sine, square
- dot transmission



○ □ Experimental data
 - - - Simulation of ΔQ , $T_{el} = 0$ mK

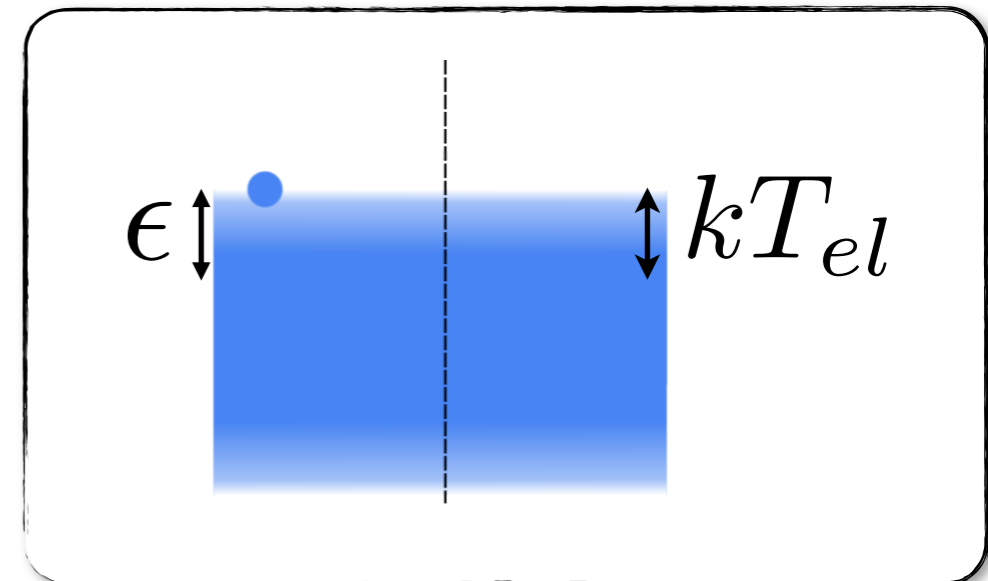


- □ Experimental data
- Simulation of ΔQ , $T_{el} = 0$ mK
- Simulation of ΔQ , $T_{el} = 150$ mK

Autocorrelations proves **single charge emission**

HBT correlations :

- reveal the effect of **thermal excitations**
- evidence of **engineering of the wavepacket**



Perspectives :

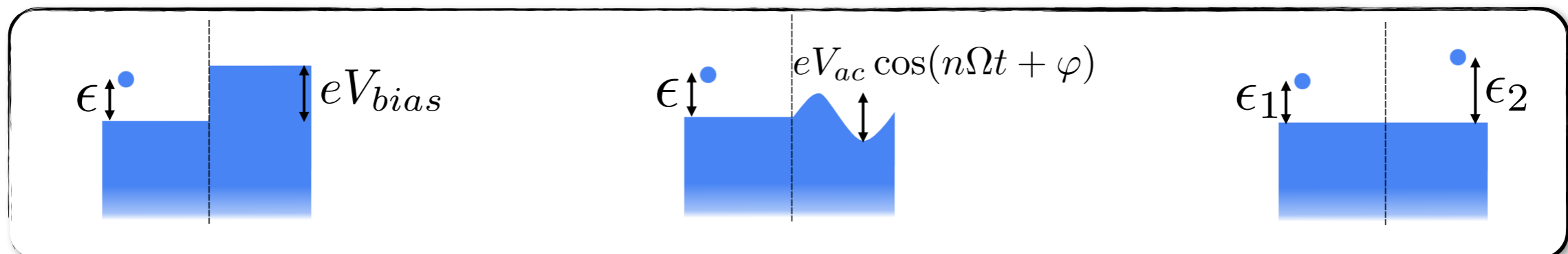
Energy distribution of emitted quasiparticles

Complete tomography protocol (cf. P. Degiovanni's talk)

Hong-Ou-Mandel experiment

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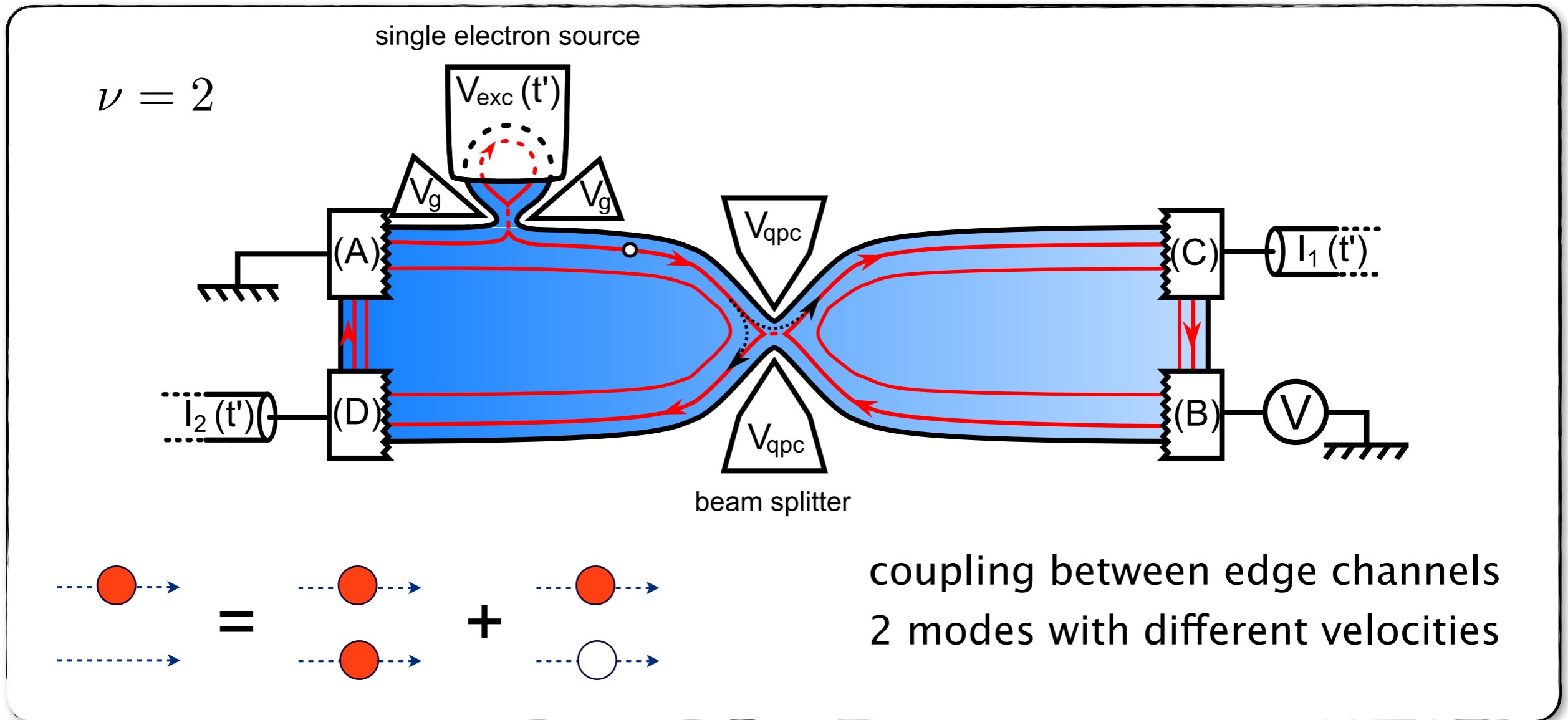
P. Degiovanni
C. Grenier
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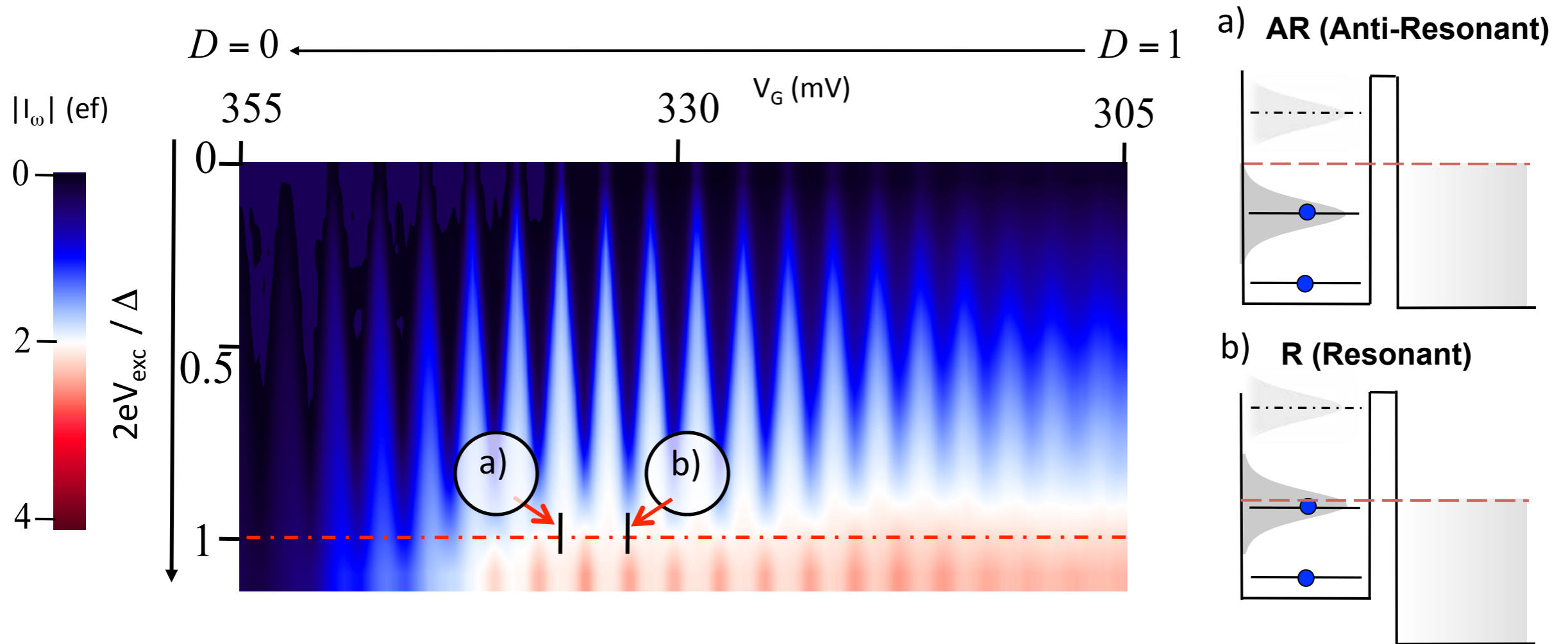
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