

Galactic perturbations of the cosmic flow

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credit: Raphael Errani

Rencountres du Vietnam, Quy Nhon
8 July 2016

The timing argument

I - Relative motion between mass-less particles
“A” and “G” in Universe with no radiation

$$\frac{\ddot{r}}{r} = -\frac{4\pi G}{3}\rho + H_0^2\Omega_\Lambda$$

Friedmann eqs.

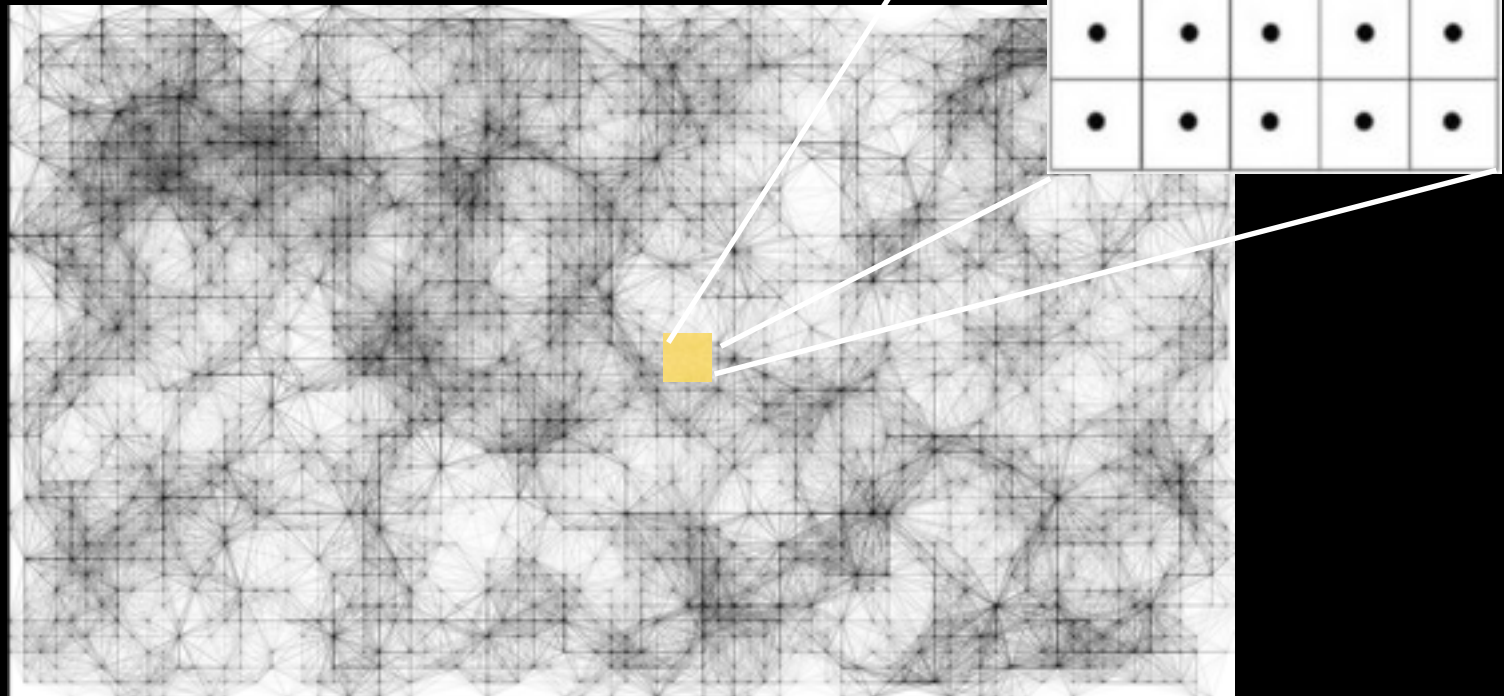
2- (Local) perturbations of Hubble flow by mass $M = M_A + M_G$

$$\ddot{r} = -\frac{GM}{r^2} + H_0^2\Omega_\Lambda r$$

‘Timing argument’

Kahn & Woltjer (1957)

$r(t = t_{\text{now}}), v(t = t_{\text{now}})$



Orbits

$$\ddot{r} = -\frac{GM}{r^2} + H_0^2 \Omega_\Lambda r.$$

relative motion between 2 massive particles in an expanding Universe

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The general solution is a Keplerian **radial** orbit

$$r = a(1 - \cos 2\eta); \quad (1)$$

where

$$a = GM/(-2E)$$

is the semi-major axis of the orbit, **E** is the orbital energy, and **η** is an angle typically referred to as the eccentric anomaly, which can be calculated numerically from the following equation

$$2\eta - \sin 2\eta = (GM/a^3)^{1/2} t. \quad (2)$$

THE DYNAMICAL AGE OF THE LOCAL GROUP OF GALAXIES

By D. Lynden-Bell

Institute of Astronomy, The Observatories, Cambridge

The distance to those Local Group members whose expansion has just been stopped by the gravity of the Local Group yields Mt^2 where M is the mass of the Group and t the time since expansion began. The distance and radial velocity of M31 yield a relationship between Mt^2 and t . Thus M and t may be deduced. Sandage–Tammann distances to Local Group members yield $t = 1.6 \times 10^{10}$ years and $M = 3.6 \times 10^{12} M_{\odot}$. Accurate distances to outlying members of the Local Group could refine this method and make a lasting contribution to cosmogony and cosmology.

Perturbed Hubble flow

Define 2 orbital frequencies (Lynden-Bell 1981):

$$\Omega = \left(\frac{GM}{r^3} \right)^{1/2}$$

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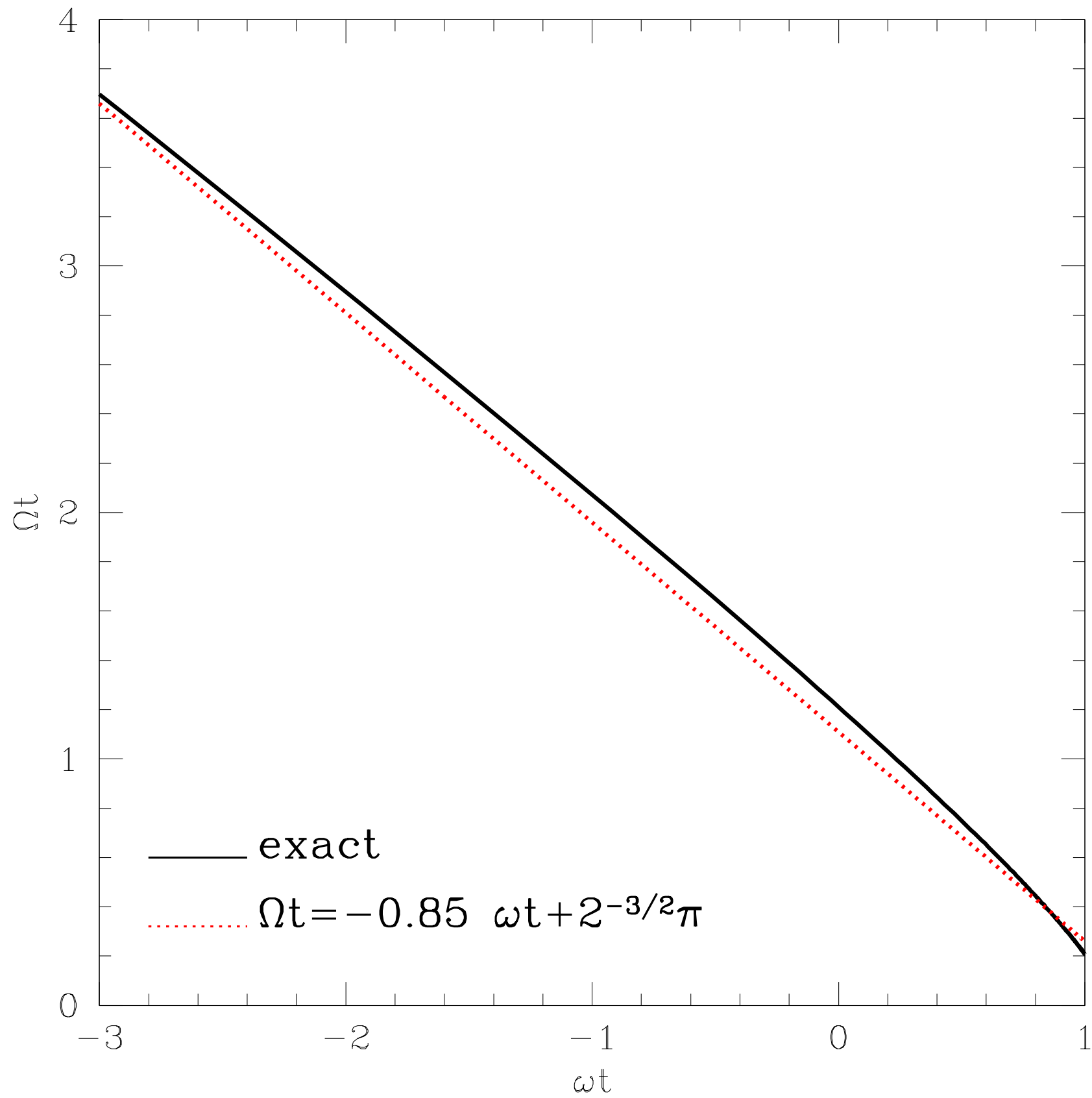
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Multiply **ω** by the time **t** and insert **Ωt**

$$\omega t = \frac{\eta - \sin \eta \cos \eta}{\sin^3 \eta} \cos \eta.$$



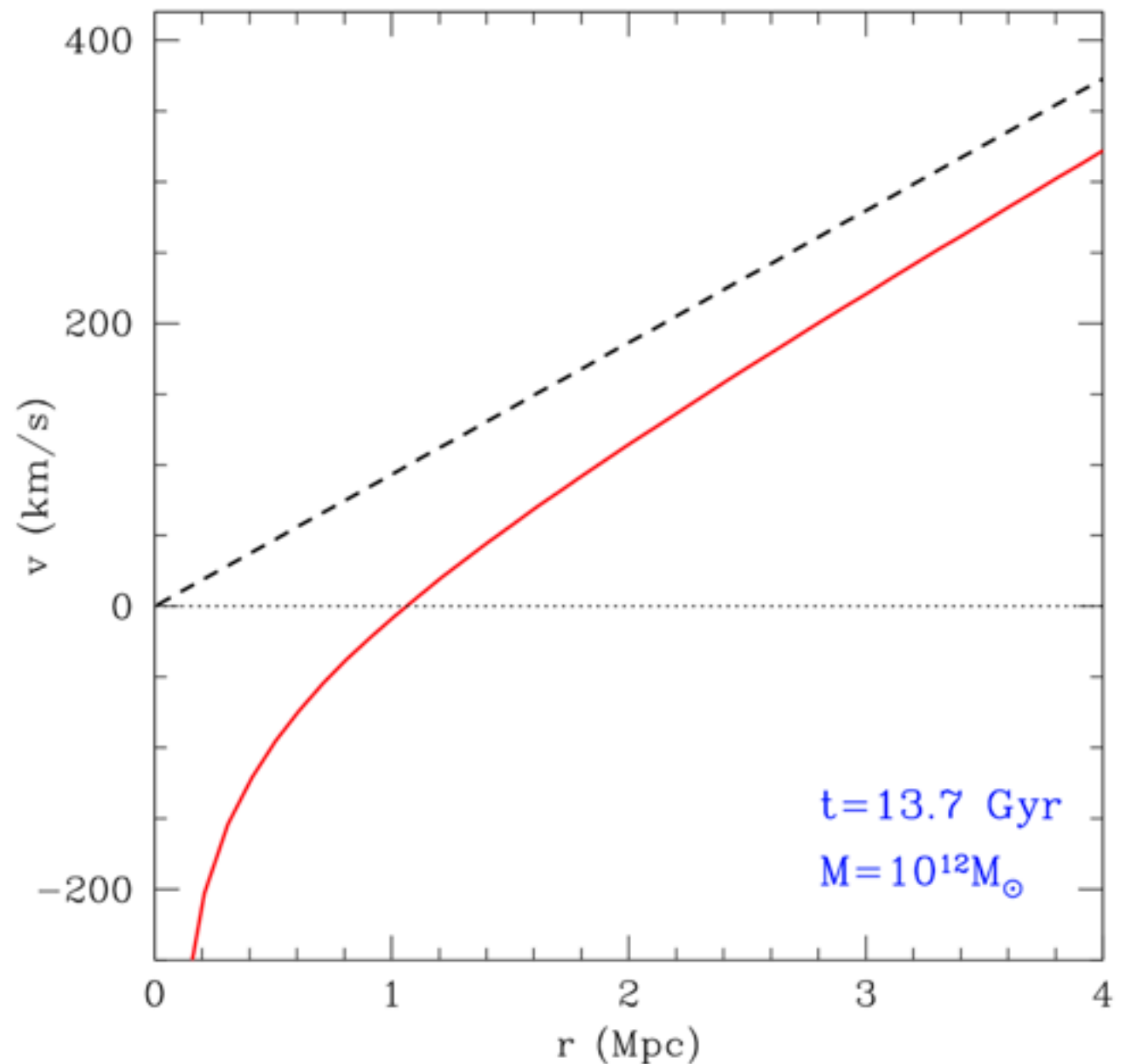
Perturbed Hubble Flow

$$\Omega t = -0.85\omega t + 2^{-3/2}\pi$$

$$\left(\frac{GM}{r^3}\right)^{1/2} t = -0.85\frac{v}{r}t + 2^{-3/2}\pi$$

solving for $v=v(r)$

$$v \simeq 1.2\frac{r}{t} - 1.1\left(\frac{GM}{r}\right)^{1/2}. \quad (\text{red curve})$$



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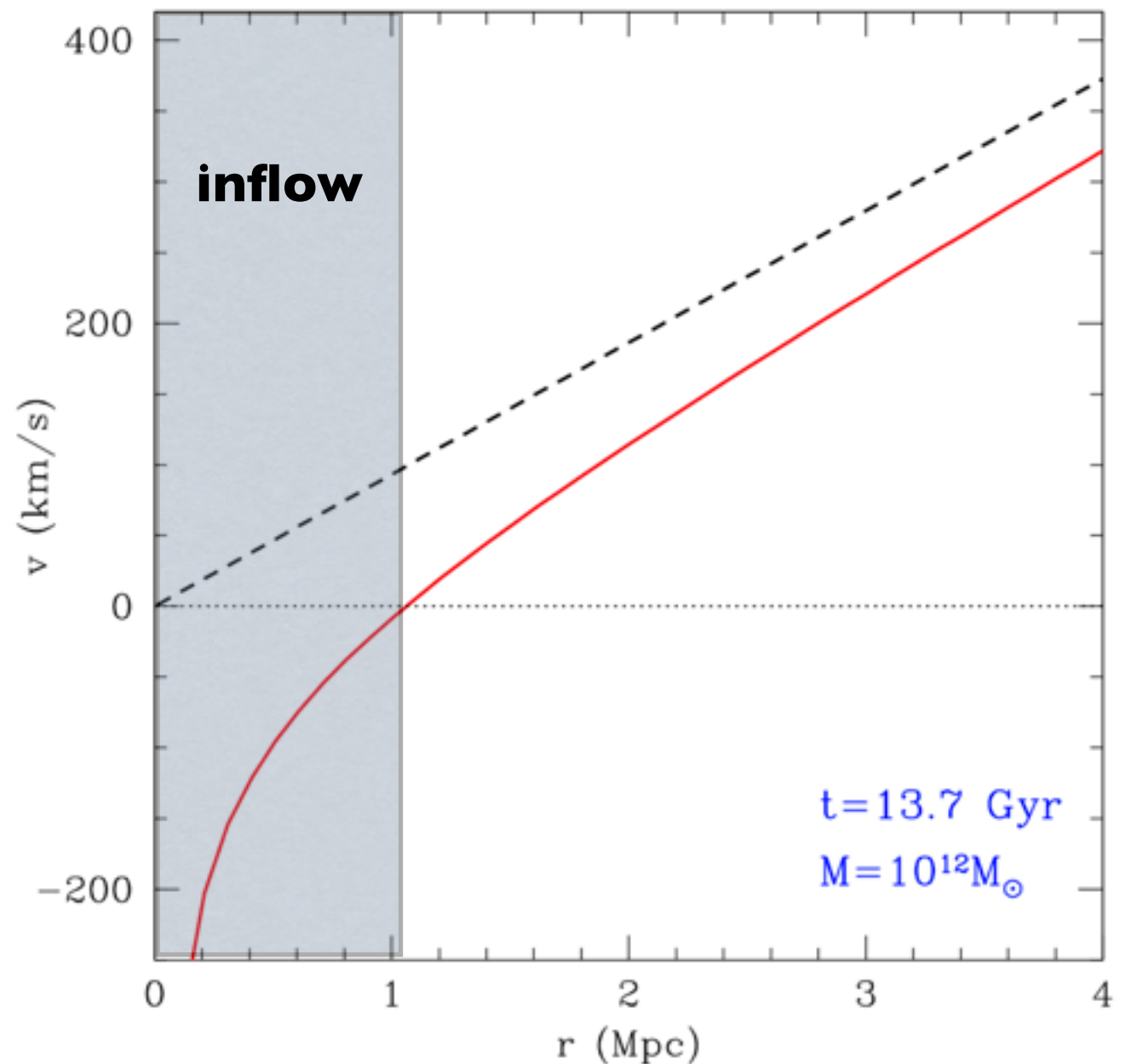
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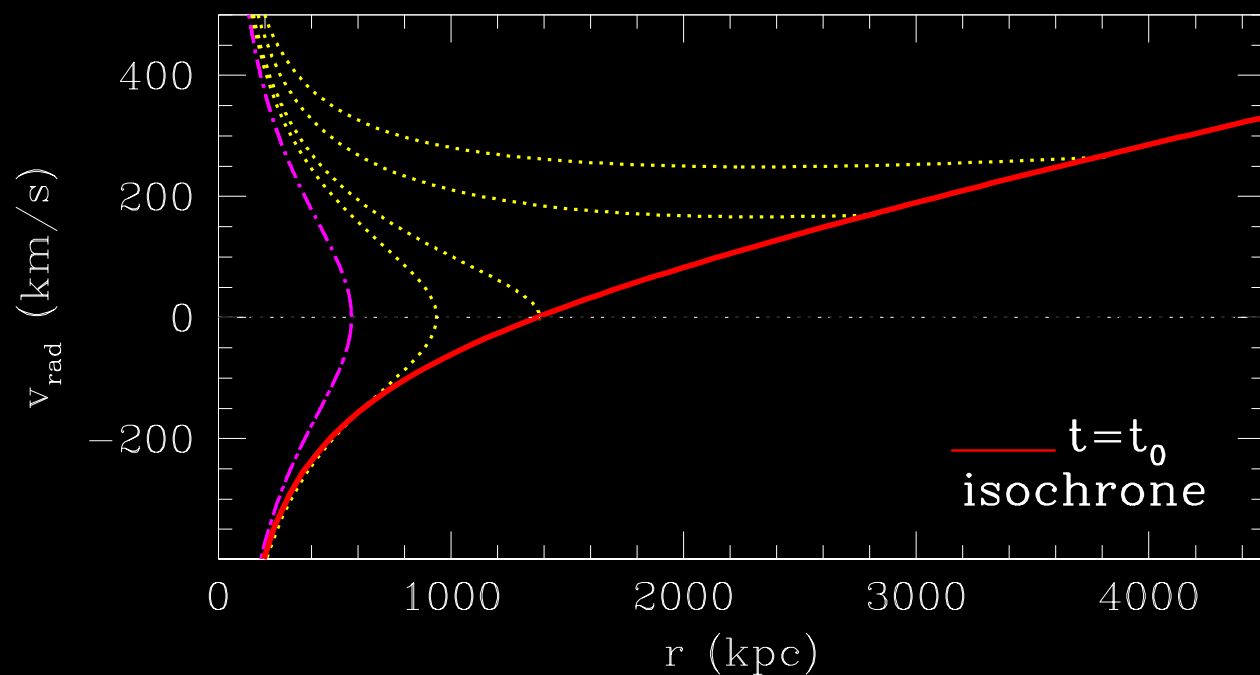
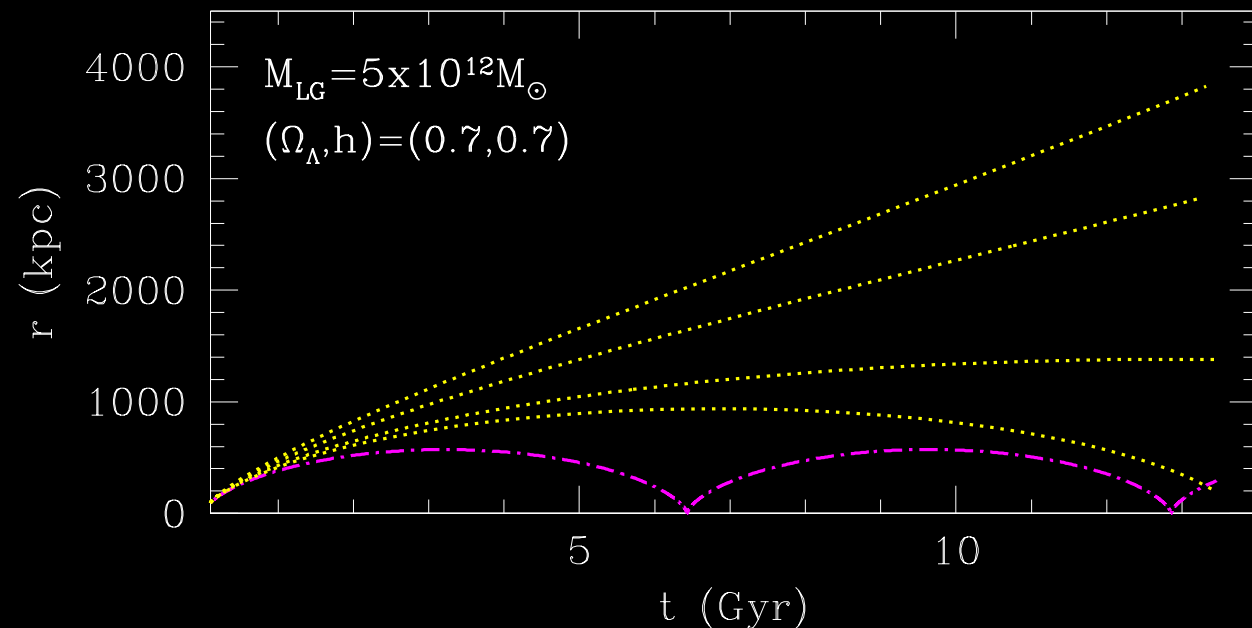
The expansion of the Universe is momentarily halted ($\mathbf{v=0}$) at $\mathbf{r_0=(GM\ t^2)^{1/3}}$ (turn-around radius).



NOTE: $\mathbf{r_0}$ grows with time even when $\mathbf{M=const.}$

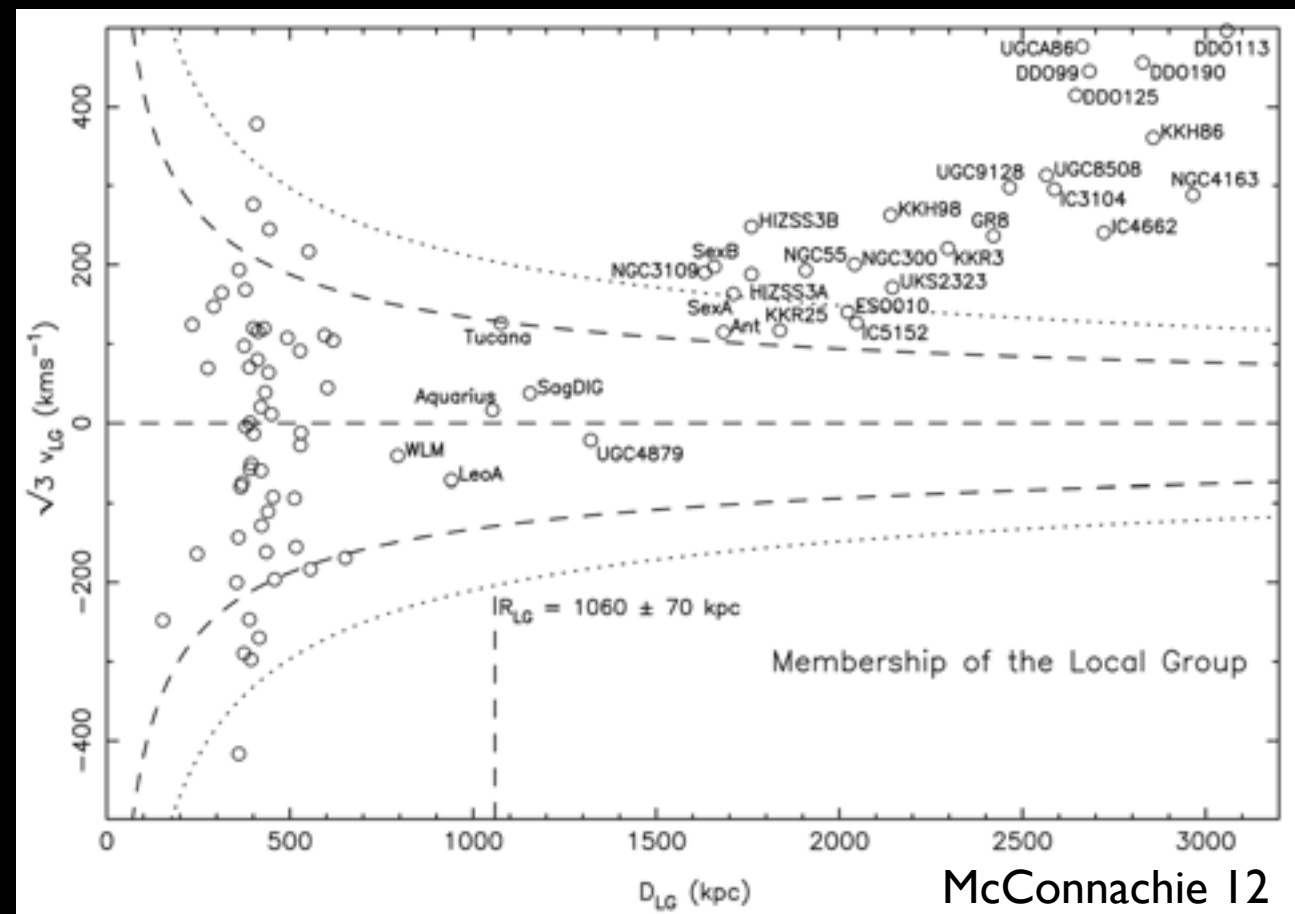
$$\ddot{r} = -\frac{GM}{r^2}$$

Galaxies around the Local Group



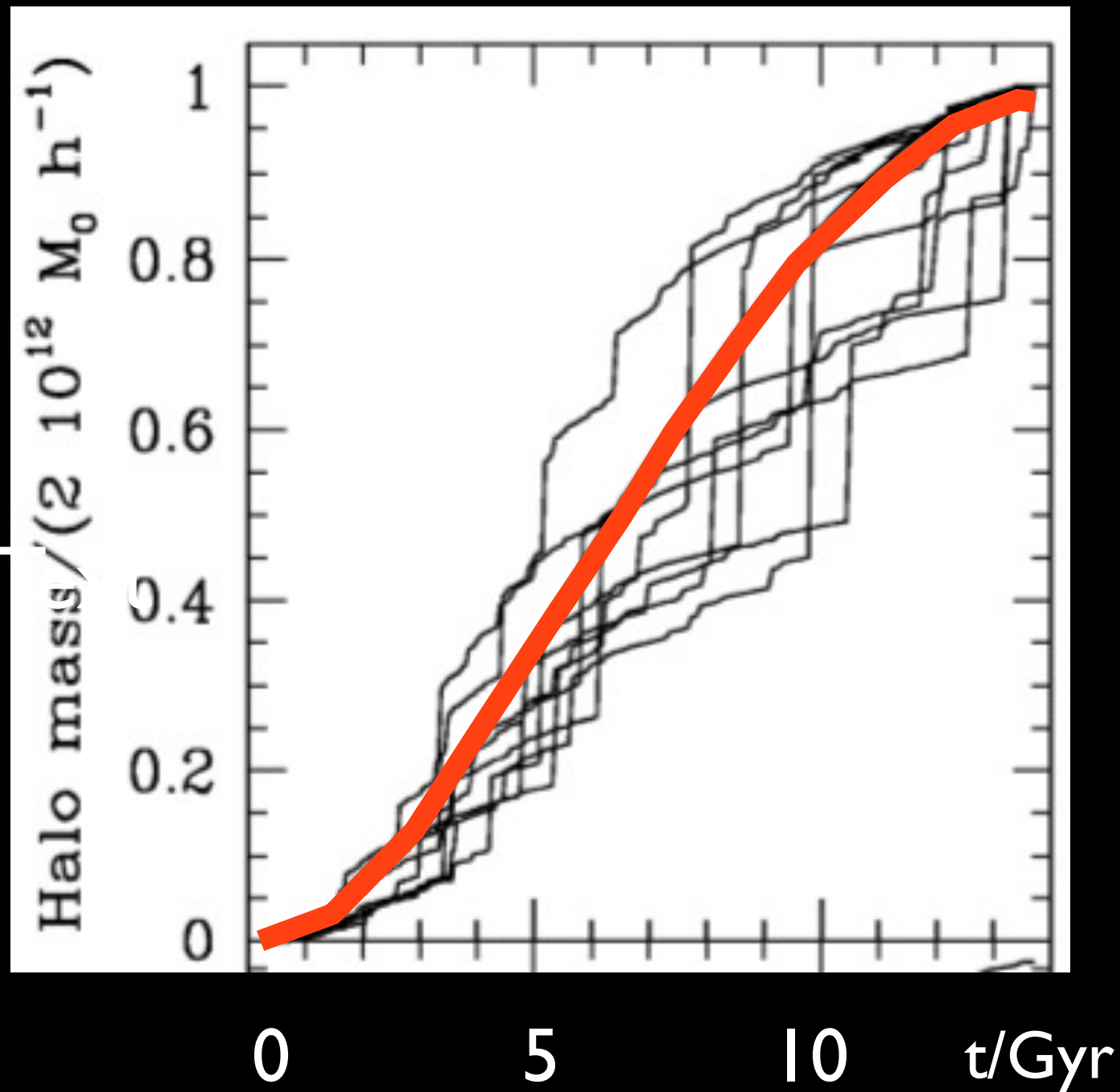
Isochrones can be used to model the local Hubble flow

Lynden-Bell (1981); Sandage (1986);
Peñarrubia et al (2014, 2016)



Universe has a finite age!

But ... how accurate is to assume $M=\text{const}$?



hierarchical growth via mergers

$$\langle M(z) \rangle \approx M_0 \exp[-z/z_c]$$

Wechsler et al. (2002)

What is the effect of $M(t)$ on the mass derived from timing argument?

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$$F(r, t) = -GM(t)r^n$$

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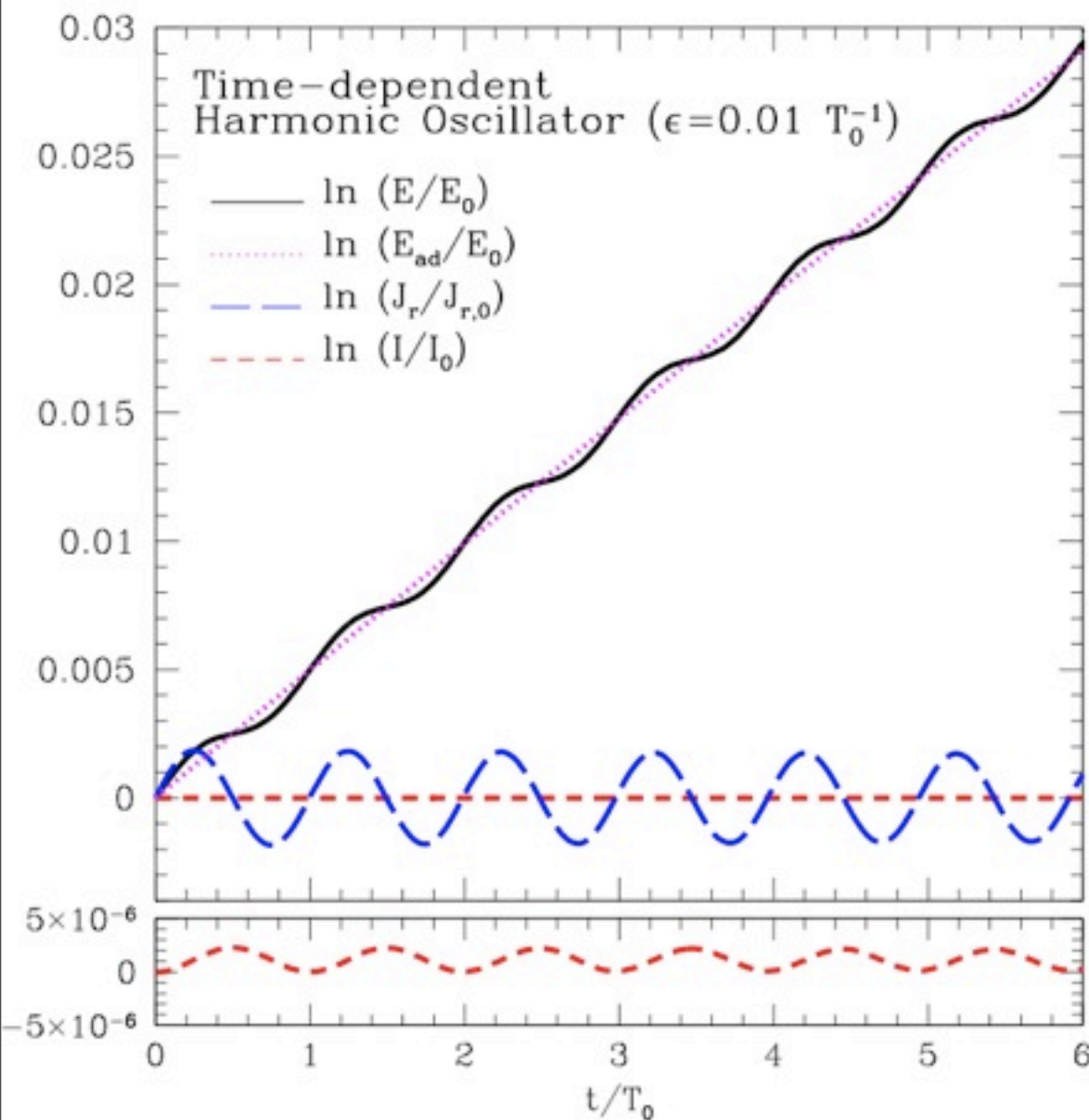
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For power-law fields

$$\ddot{R} + r'^{n-1} GM(t) R^n = -r'^{n-1} GM_0 R^{-3}. \quad \longrightarrow \quad R(t) \approx \left[\frac{M_0}{M(t)} \right]^{1/(3+n)}$$

Invariant space

Integrals of motion in (\mathbf{r}', τ) become **dynamical invariants** in (\mathbf{r}, t) Peñarrubia 2014, 2016



$$I = \left(\frac{d\mathbf{r}'}{d\tau} \right)^2 + \Phi'(\mathbf{r}') = \frac{1}{2}(R\mathbf{v} - \dot{R}\mathbf{r})^2 + \frac{1}{2}\ddot{R}Rr^2 + R^2\Phi(\mathbf{r}, t)$$

$$E = \underbrace{\frac{I}{R^2}}_{\text{Adiabatic term } \mathbf{E}_a} + \underbrace{(\mathbf{r} \cdot \dot{\mathbf{r}})\frac{\dot{R}}{R} - \frac{1}{2}r^2\left(\frac{\ddot{R}}{R} + \frac{\dot{R}^2}{R^2}\right)}_{\text{Deviation } \Delta};$$

$\mathbf{r} \cdot \mathbf{v}$ independent of angular momentum

Δ varies in phase with the particle orbit

Isochrones in time-dependent potentials

Expressing the frequencies in the invariant coordinates (\mathbf{r}', τ)

$$\Omega(t) \mapsto \frac{1}{R^2} \left[\frac{GM_0}{r'^3} \right]^{1/2} = \frac{\Omega'}{R^2(t)}$$

$$\omega(t) \mapsto \frac{1}{r'} \frac{dr'}{d\tau} = \frac{1}{R(t)r'} \left[2 \left(\frac{E_0}{R^2(t)} + \frac{GM_0}{r'R^2(t)} \right) \right]^{1/2} = \frac{\omega'}{R^2(t)} \quad \text{for} \quad \frac{dM}{dt} \ll \frac{M_0}{t_0} \quad (\text{adiabatic})$$

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In a Keplerian potential **n=-2**

$$R(t) \approx \left[\frac{M_0}{M(t)} \right]^{1/(3+n)} = \frac{M_0}{M(t)}$$

The isochrones evolve as

$$\Omega' t + m \omega' t = n R^2(t) = n \left[\frac{M_0}{M(t)} \right]^2 ;$$

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At present **M(t₀)=M₀** and thus **R(t₀)=1**

$$\Omega' t_0 + m \omega' t_0 = n$$

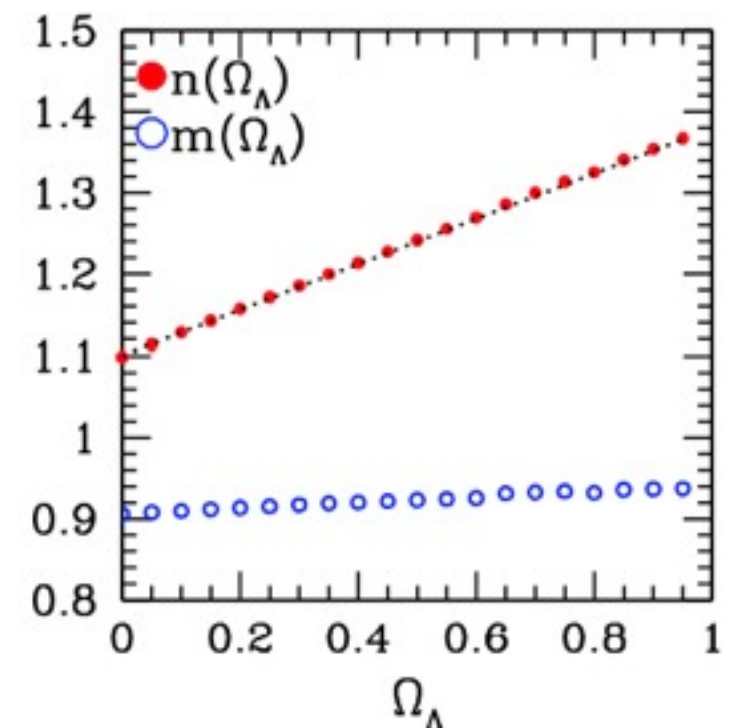
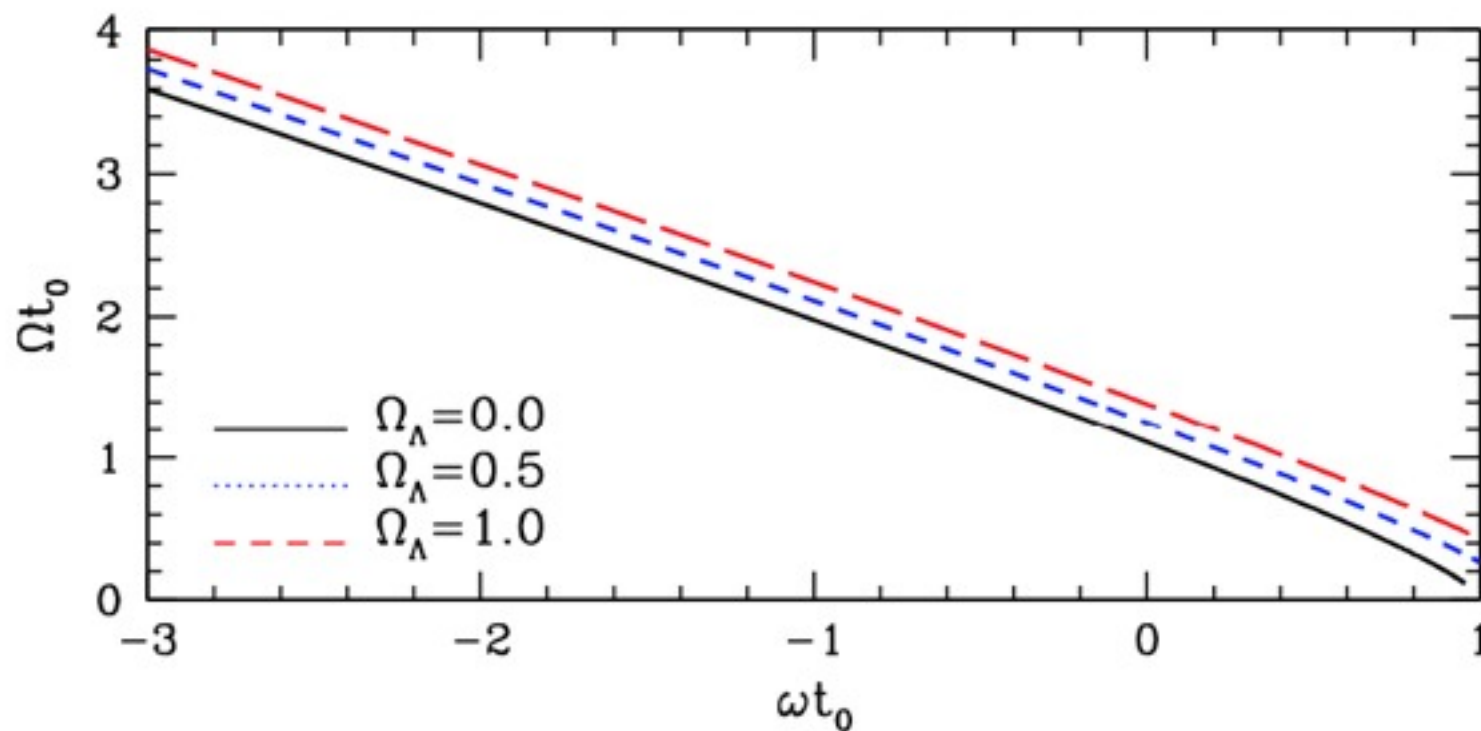
*Perturbed Hubble flow contains **NO** information on past **M(t)***

Isochrones with Dark Energy

The above relation can be generalized to $\Omega_\Lambda \neq 0$

$$\ddot{r} = -\frac{GM}{r^2} + H_0^2 \Omega_\Lambda r.$$

$$v \approx (n - \Omega t_0) \frac{r}{mt_0}$$

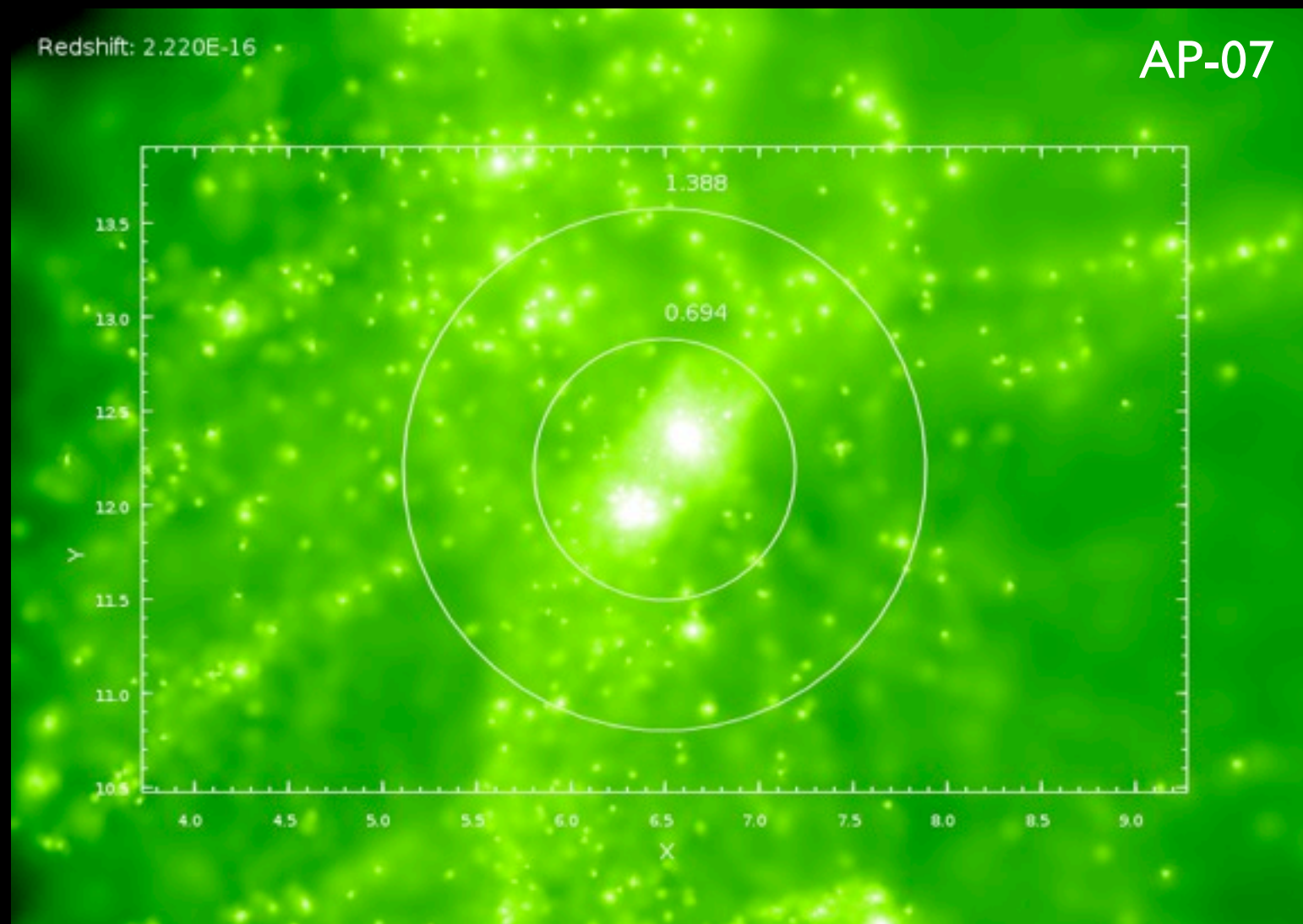


$$v \approx 1.2 \frac{r}{t_0} - 1.1 \left(\frac{GM}{r} \right)^{1/2} + 0.16 \Omega_\Lambda \frac{r}{t_0}$$

LCDM isochrones

Dark energy steepens the Hubble flow by $\delta_v = v - v(\Omega_\Lambda = 0) \sim 0.16 \Omega_\Lambda r/t_0$ independently of M . For galaxies within 3Mpc and $t_0 = 13.46$ Gyr this implies $\delta_v < 24$ kms

Tests with cosmological N-body sims.

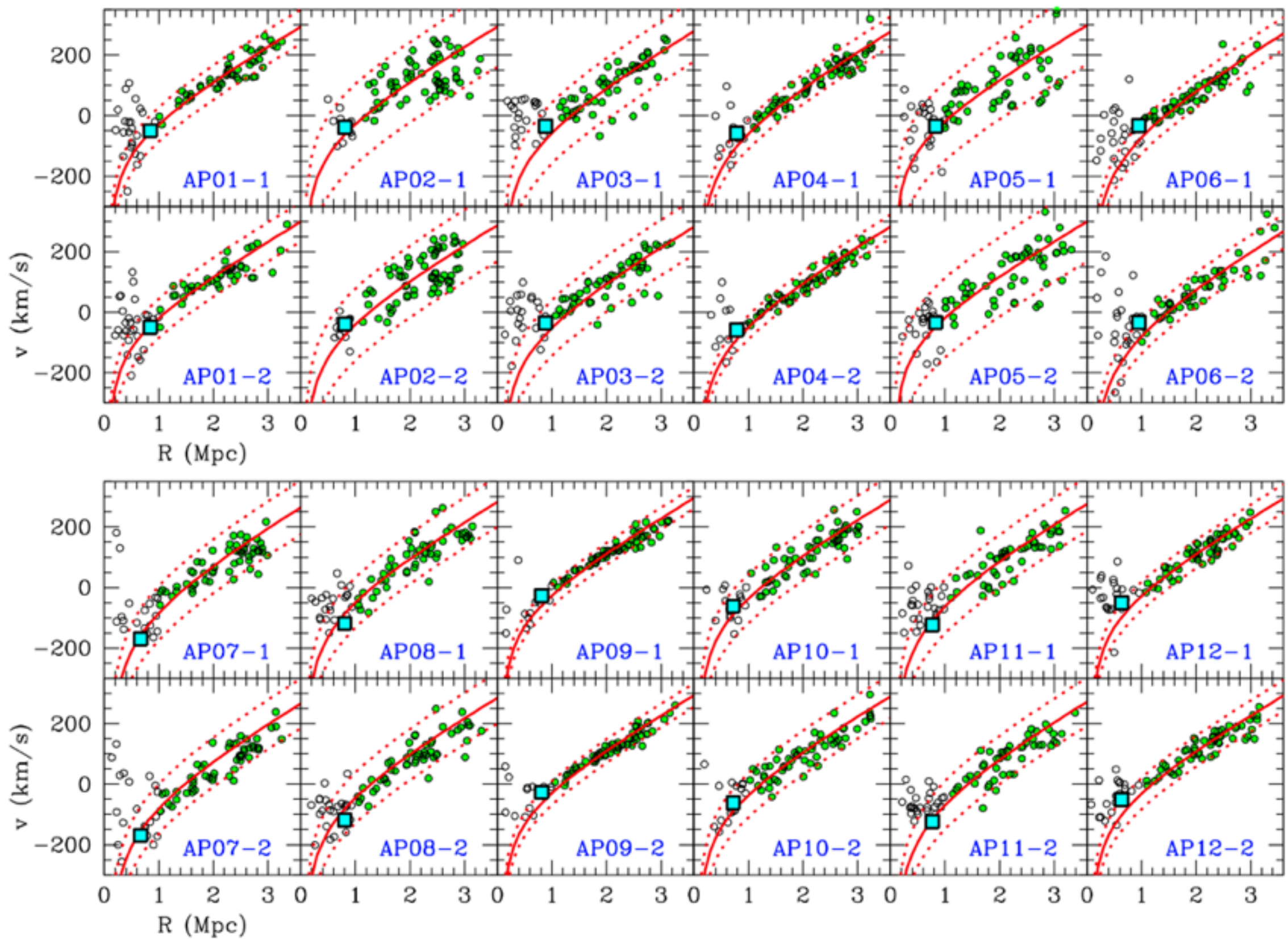


12 Local Group realizations
(Sawala+16; Fattahi+16)

50 random substructures between
1-3 Mpc

Fit total mass (M), mass ratio (q) and
hyperparameter (σ)

$$v \approx 1.2 \frac{r}{t_0} - 1.1 \left(\frac{GM}{r} \right)^{1/2} + 0.16 \Omega_\Lambda \frac{r}{t_0}$$



Summary

- **The timing argument describes the local expansion of the Universe in a quasi non-linear regime (no shell crossing)**
- **For a point-mass the perturbed Hubble flow in LCDM can be derived analytically**
- **The analytical isochrones provides a simple tool for modelling Local Group dynamics within a Bayesian framework (Peñarrubia+14,16)**
- **The iso-chrone $v=v(r,t_0)$ is an adiabatic invariant, i.e. independent of $M(t)$ insofar as $dM/dt \ll M/t_0$**
- **Future work: test analytical approach with cosmological N-body models of the Local Group (e.g. Apostle, CLUES?) as well as external galaxies (issue: projection effects)**

