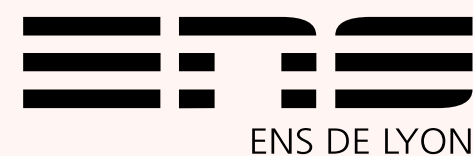




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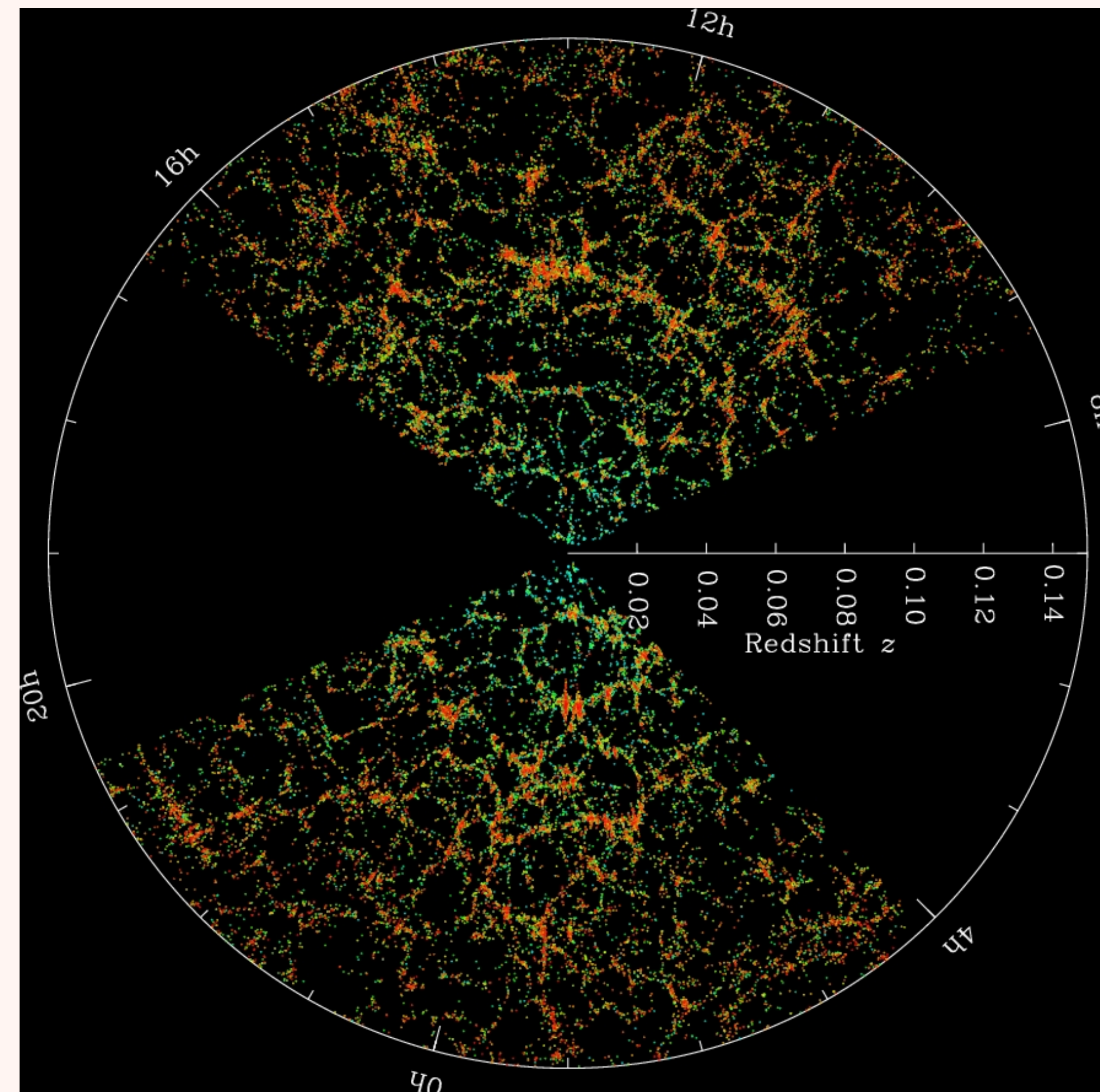


COSMOLOGICAL BACKREACTION FROM INHOMOGENEITIES/ CORRELATIONS

PhD student : Pascal WANG

Supervisors : Gilles CHABRIER (CRAL, Lyon, France),
Pascal TREMBLIN (MDLS, Saclay, France)

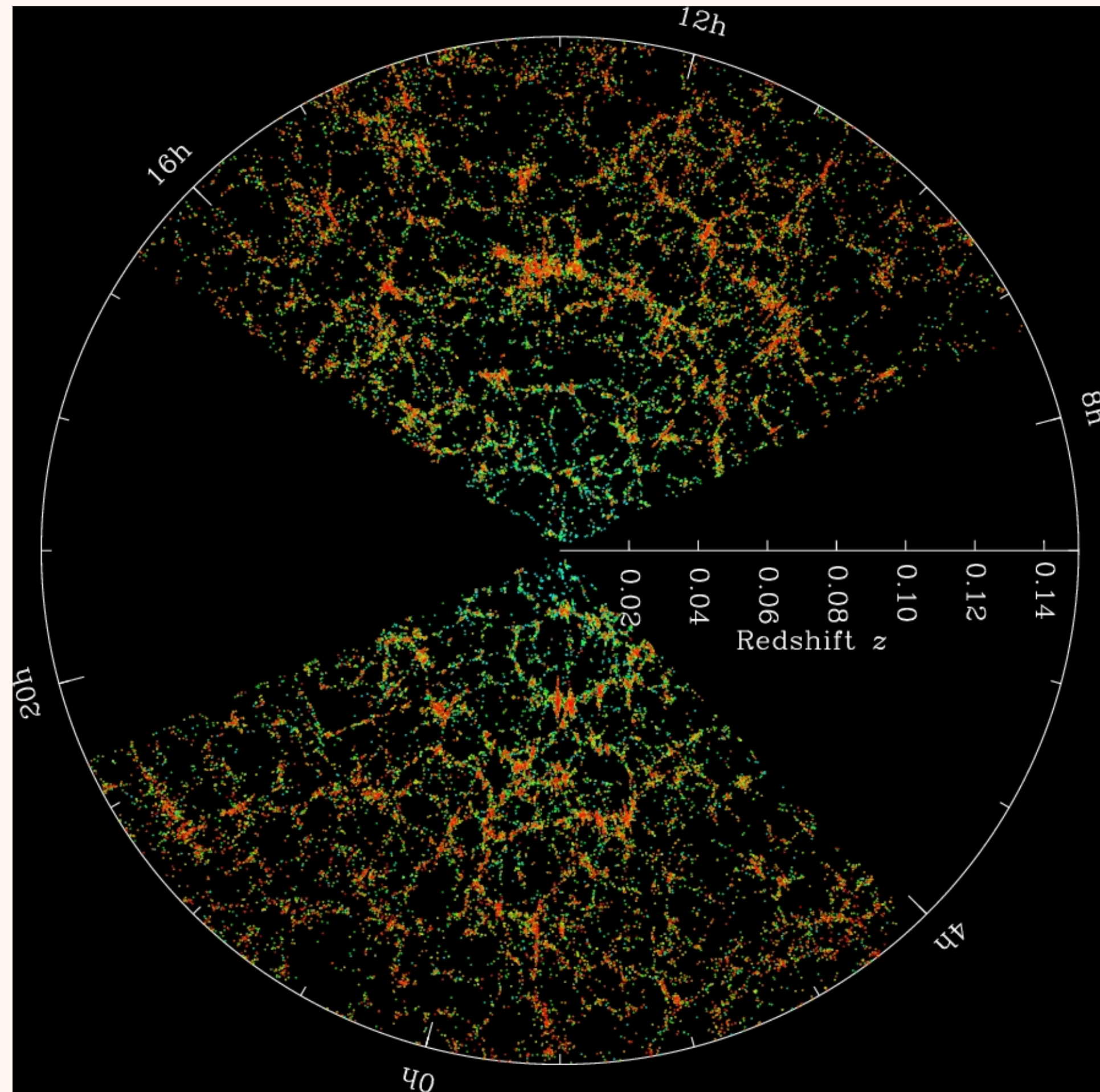
Future cosmology school - 28.04.2023 - Cargèse



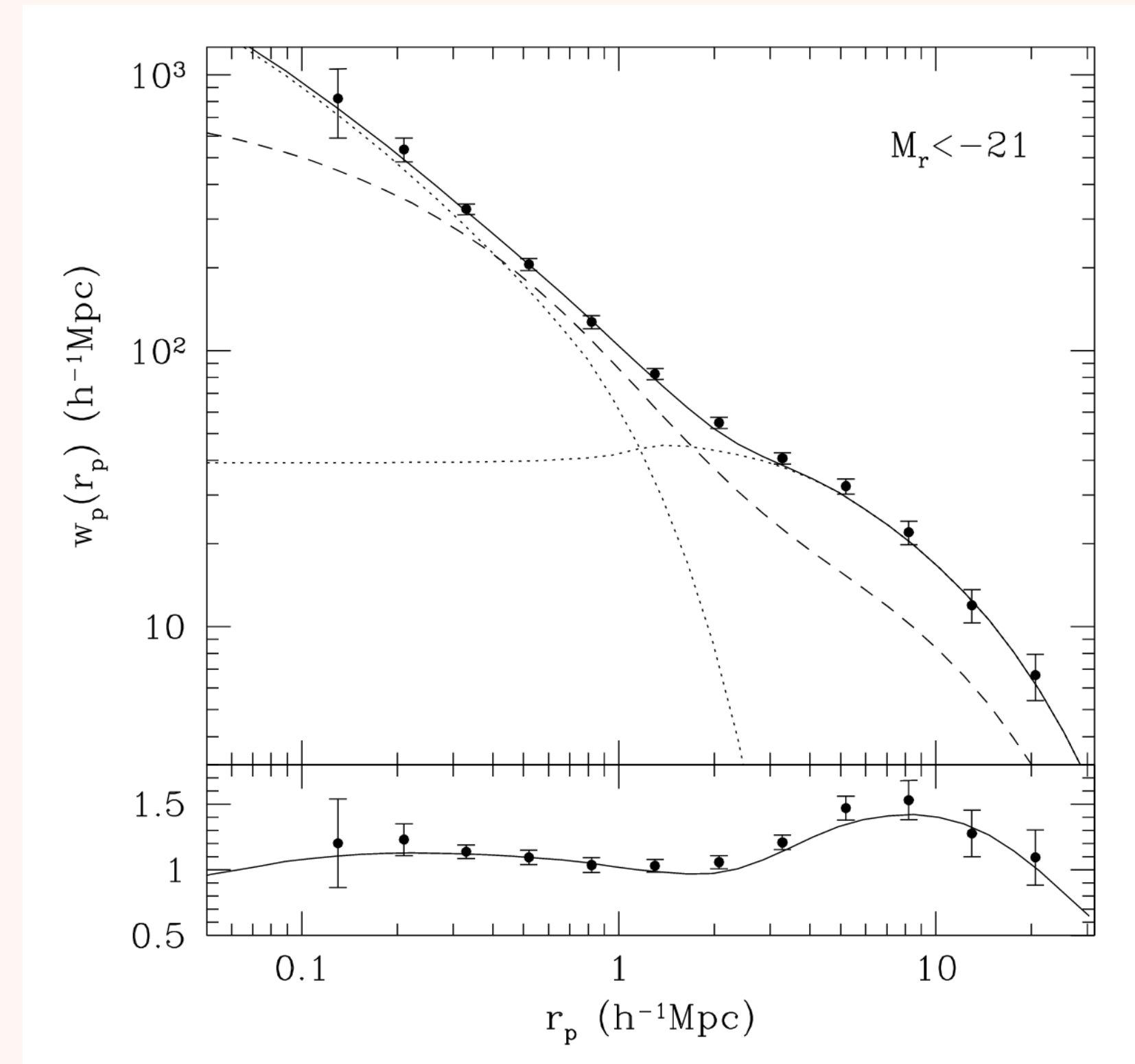
A zoomed-out view of galaxies identified by the Sloan Digital Sky Survey. Filaments and voids can be identified.

Credit : M. Blanton, SDSS

CORRELATIONS IN LARGE SCALE STRUCTURES



A zoomed-out view of galaxies identified by the Sloan Digital Sky Survey. Credit : M. Blanton, SDSS



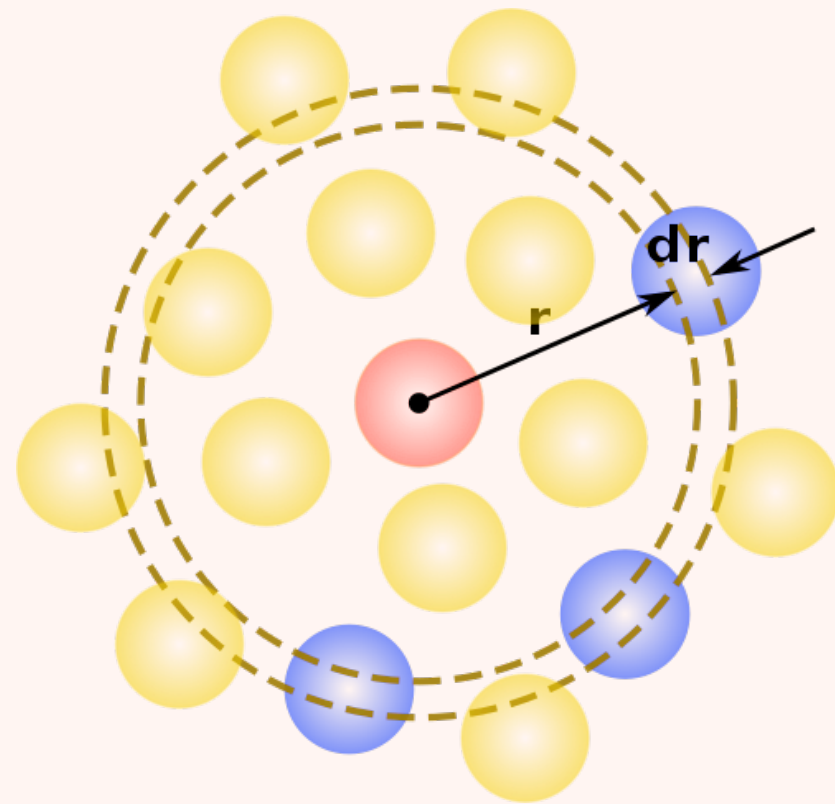
The projected correlation function for SDSS galaxies (data points with error bars). Taken from Zehavi et al. (2004)

The large-scale galaxy matter distribution of the Universe is highly correlated at small scales.

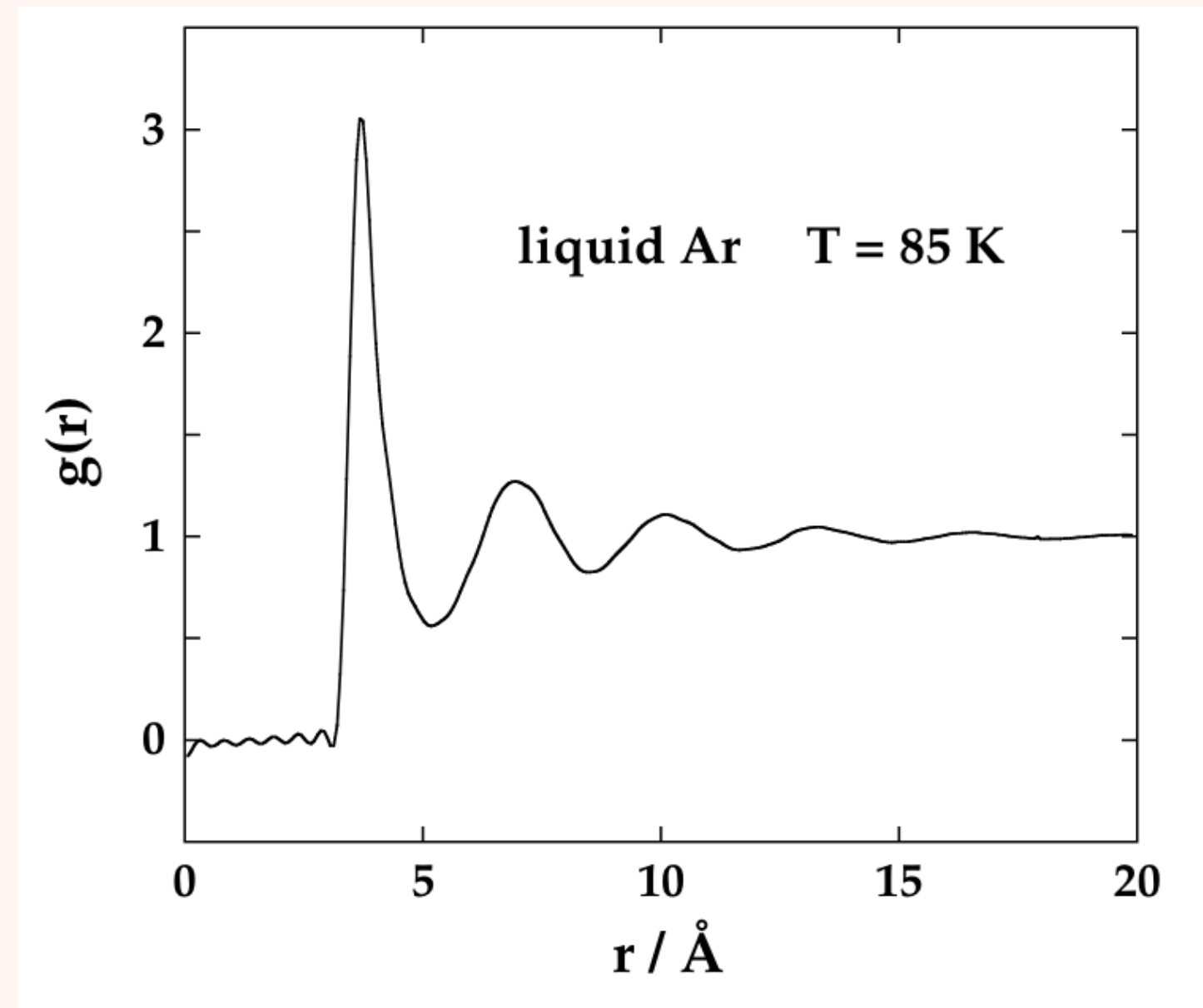
CORRELATIONS IN LIQUIDS

Radial distribution function

$$g(r) = \frac{1}{\rho} \left\langle \frac{1}{N} \sum_{1 \leq i \neq j \leq N} \delta(\mathbf{r} - \mathbf{r}_j + \mathbf{r}_i) \right\rangle$$



$4\pi r^2 \rho g(r) dr$ is the mean number of particles within the range r to $r + dr$, given that there is a particle at $r = 0$



Radial distribution function of liquid argon near the triple point, after Yarnell et al.

Thermodynamic properties

Energy

$$E = \frac{Nk_B T}{\gamma - 1} + \frac{2\pi N^2}{V} \int_0^\infty g(r) \phi(r) r^2 dr$$

Equation of state

$$PV = Nk_B T - \frac{2\pi N^2}{3V} \int_0^\infty g(r) \frac{d\phi(r)}{dr} r^3 dr$$

Transport coefficients

- shear viscosity η
- bulk viscosity χ

Liquids are correlated due to short-range interactions.

Correlations determine their physical properties.

BACKREACTION AND AVERAGING PROBLEMS

Question : Do correlations have an **impact on the dynamics of the Universe (e.g. its expansion) ?**

In connection with

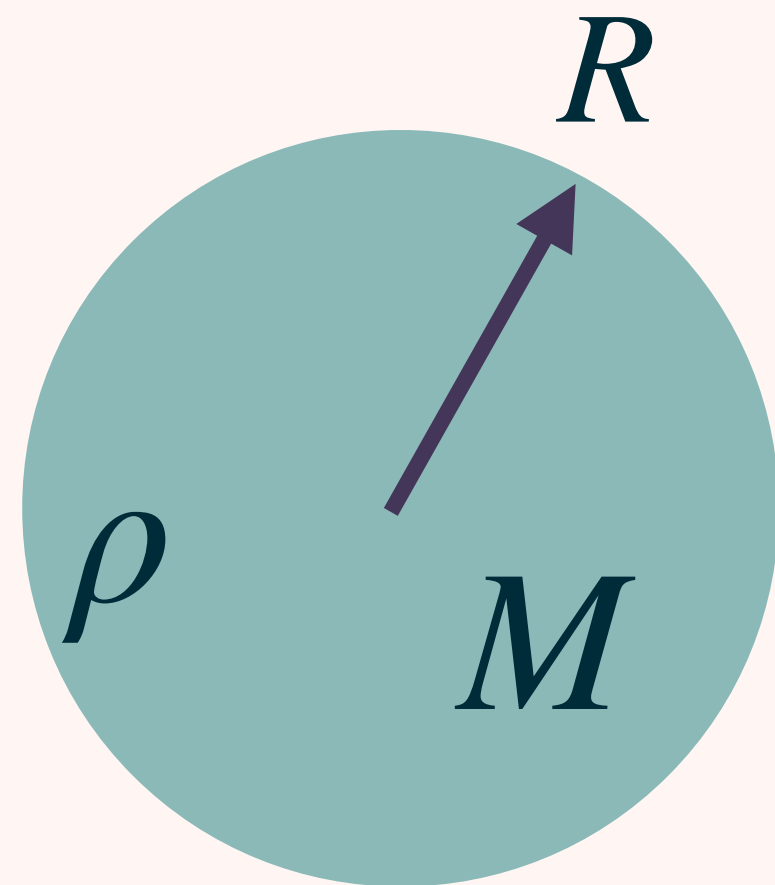
- **backreaction problem:** to what extent do inhomogeneities impact the cosmological expansion ?
- **averaging problem:** how to derive a smooth spacetime from an inhomogeneous spacetime, in a covariant way?
e.g. Macroscopic gravity (Zalaletdinov 1993), volume average of scalars (Buchert 2000)

However, many results pointing to backreaction being a small effect of order $(v/c)^2 \sim 10^{-5}$

- Effective field theory (Baumann 2012)
- Relativistic simulations (Adamek+ 2019)
- And others (see references in Peebles 2010)

NEWTONIAN DERIVATION OF FRIEDMANN EQUATIONS

Homogeneous (uncorrelated)
sphere of fluid



Conservation of mechanical energy (assumed zero here)

$$E_m = E_{kin} + E_{int} = 0$$

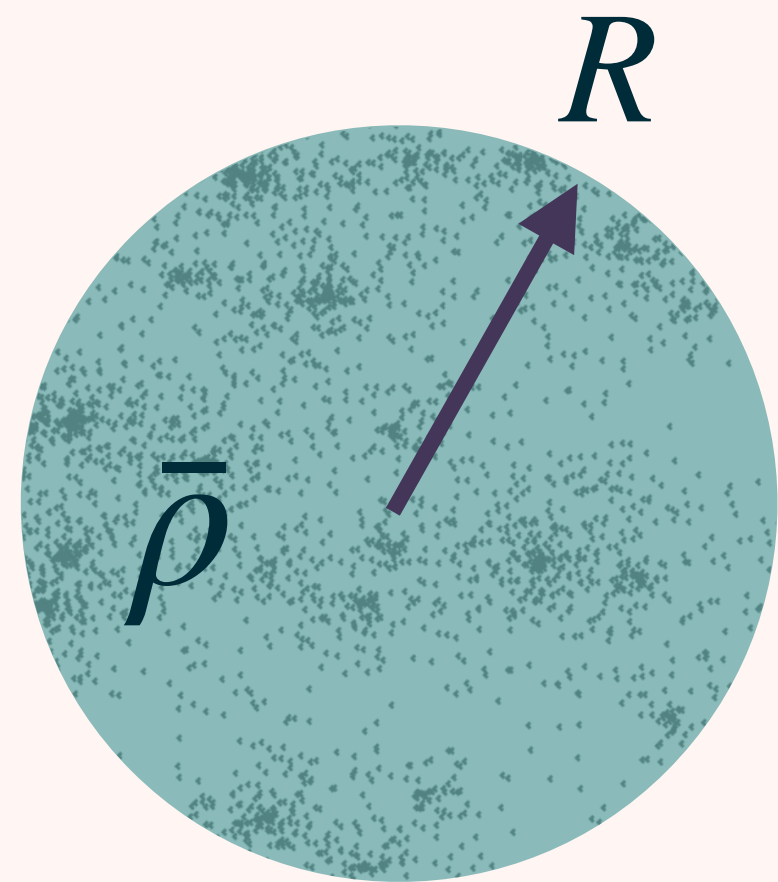
$$E_m = \frac{1}{2}m\dot{R}^2 - \frac{GmM}{R} = 0$$

Leads to the Friedmann equation

$$H^2 = \left(\frac{\dot{R}}{R} \right)^2 = \frac{8\pi G}{3} \rho$$

HEURISTIC NEWTONIAN DERIVATION IN TREMBLIN ET AL. (2022)

Sphere composed of a correlated fluid $P_2(r, r')$



Energy conservation : $E_{kin} + E_{int} = 0$

$$\frac{1}{2}M\dot{R}^2 \approx \frac{1}{2}M\langle v^2 \rangle = \alpha_{ni} \frac{3GM^2}{5R}$$

with $\alpha_{ni} = \frac{\iint_{V,V'} P_2(\mathbf{r}, \mathbf{r}') / |\mathbf{r} - \mathbf{r}'| dV dV'}{\iint_{V,V'} P_1(\mathbf{r}) P_1(\mathbf{r}') / |\mathbf{r} - \mathbf{r}'| dV dV'}$

Modified Friedmann equation

$$H^2 = \frac{\dot{R}^2}{R^2} = \alpha_{ni} \frac{8\pi G}{3} \bar{\rho}$$

Result

The Hubble rate is modified by an **amplification coefficient** α_{ni} related to substructures \rightarrow backreaction effect

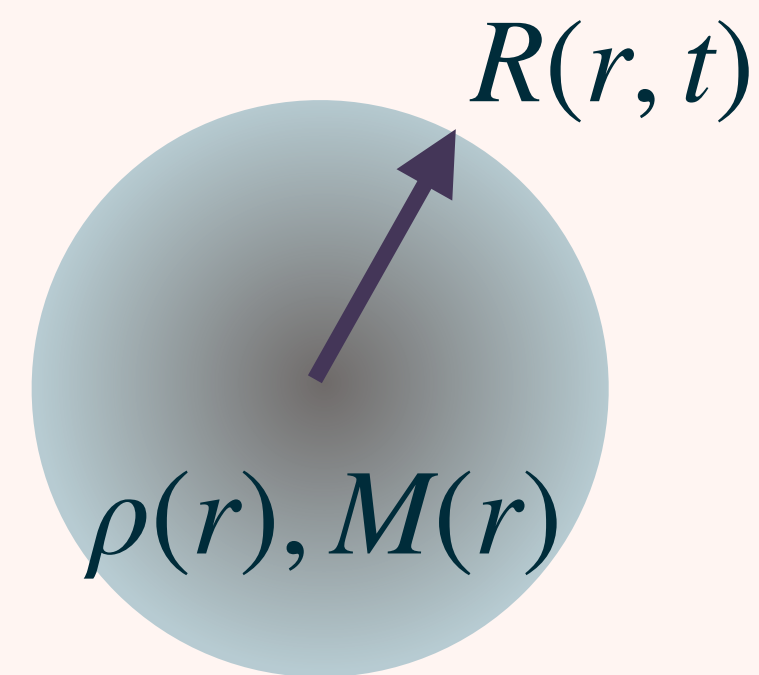
How does this Newtonian heuristic hold in a general relativistic setting?

AVERAGING INHOMOGENEITIES IN THE LTB MODEL

Idea : derive equations similar to the modified Friedmann equation of Tremblin et al. 2022, in a rigorous and general relativistic setting [1]

Lemaître-Tolman-Bondi (LTB) model

Spherically symmetric dust
(pressureless)
inhomogeneous model



Local expansion equation

$$\left(\frac{\partial R}{\partial t}\right)^2 = \frac{2GM(r)}{R} + 2E(r)$$

With line element

$$ds^2 = - dt^2 + \frac{(\partial R/\partial r)^2}{1 + 2E} dr^2 + R^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

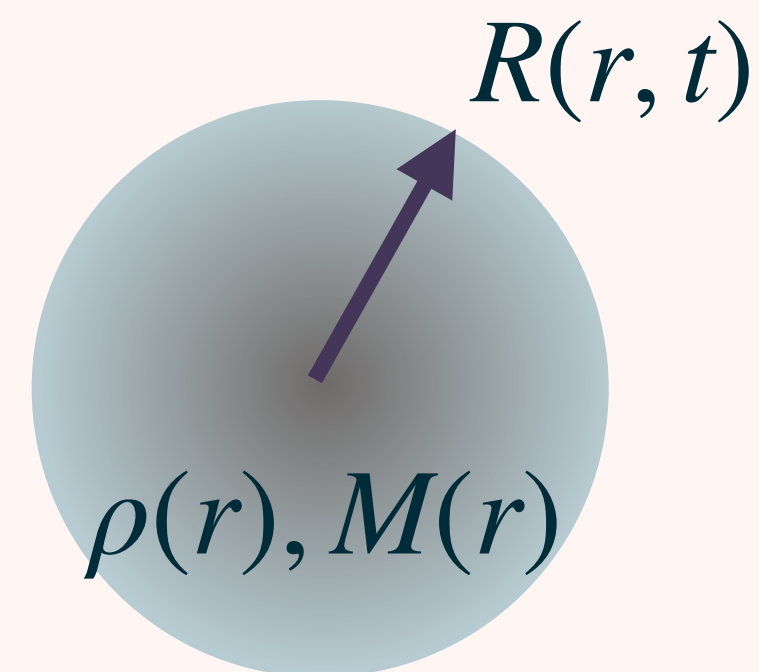
[1] PW, Tremblin, Chabrier, submitted

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Mass-weighted
averaging operator

$$\langle A \rangle_M \equiv \frac{1}{M(r_0)} \int_{V_{r_0}} dV \rho A$$

$$\dot{R}_M \equiv \left\langle \left(\frac{\partial R}{\partial t}\right)^2 \right\rangle_M^{1/2}$$

Mass-weighted effective
scale factor

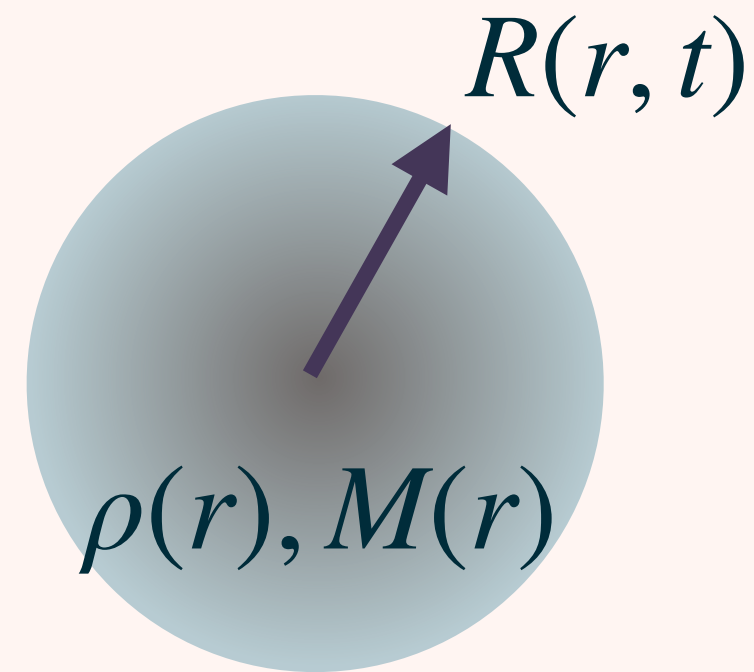
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Mass-weighted effective
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Effective, averaged, expansion equation

$$H_M^2 = \alpha \frac{8\pi G}{3} \bar{\rho}(R_0) + 2 \frac{\langle E \rangle_M}{R_M^2}$$

$$\alpha = \frac{\int_{V_{R_0}} \int_{V_R} d\tilde{V}_R dV_{R'} P_2(R, R') / R}{\int_{V_{R_0}} \int_{V_R} d\tilde{V}_R dV_{R'} \bar{\rho}(R_0)^2 / R} \cdot \frac{\langle R^2 \rangle_V}{R_M^2}$$

Results

- **Two-point function** $P_2(R, R') = \rho(R)\rho(R')$ contribution to the amplification coefficient α , similar to Newtonian heuristics
- **Amplification of matter contribution** compared to FLRW spacetime

[1] PW, Tremblin, Chabrier, submitted

DISCUSSION

We showed how the Tremblin et al. 22 Newtonian heuristic can be derived from a **GR model using a mass average**.

We illustrated and compared the **influence of the averaging procedure** (e.g. with Buchert's formalism) on the form and magnitude of backreaction terms.



Limitations

- The LTB model contains only 1 spherical overdensity → extension to Swiss-cheese models?
- Lack of justification for the averaging formalism (why pick this one over others?) and of connection with observables related to light propagation