

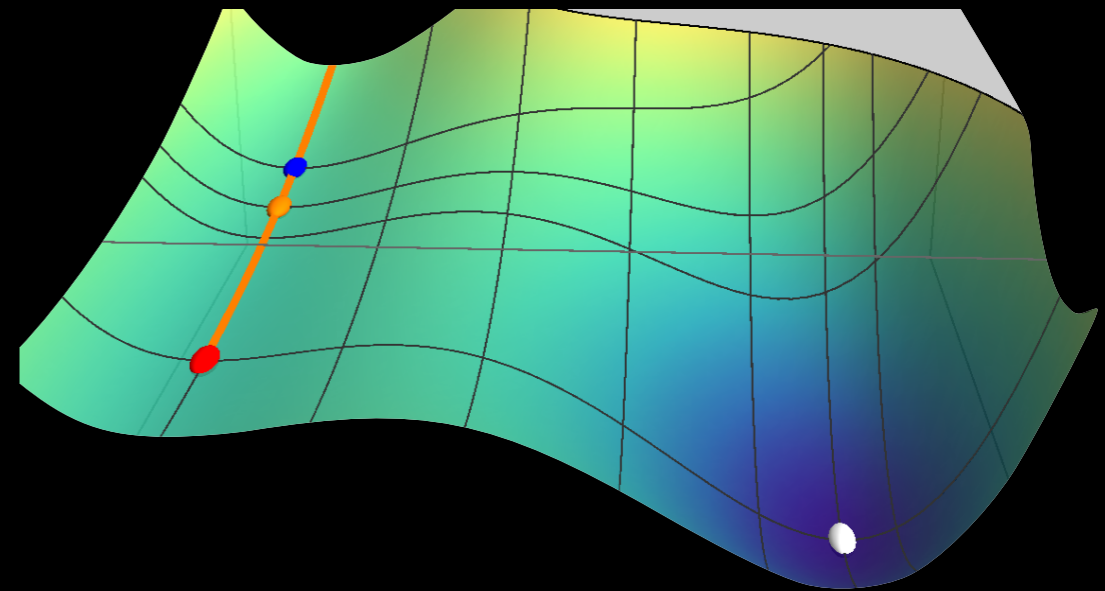


# New Early Dark Energy as a resolution to the Hubble tension

Florian Niedermann  
Nordita

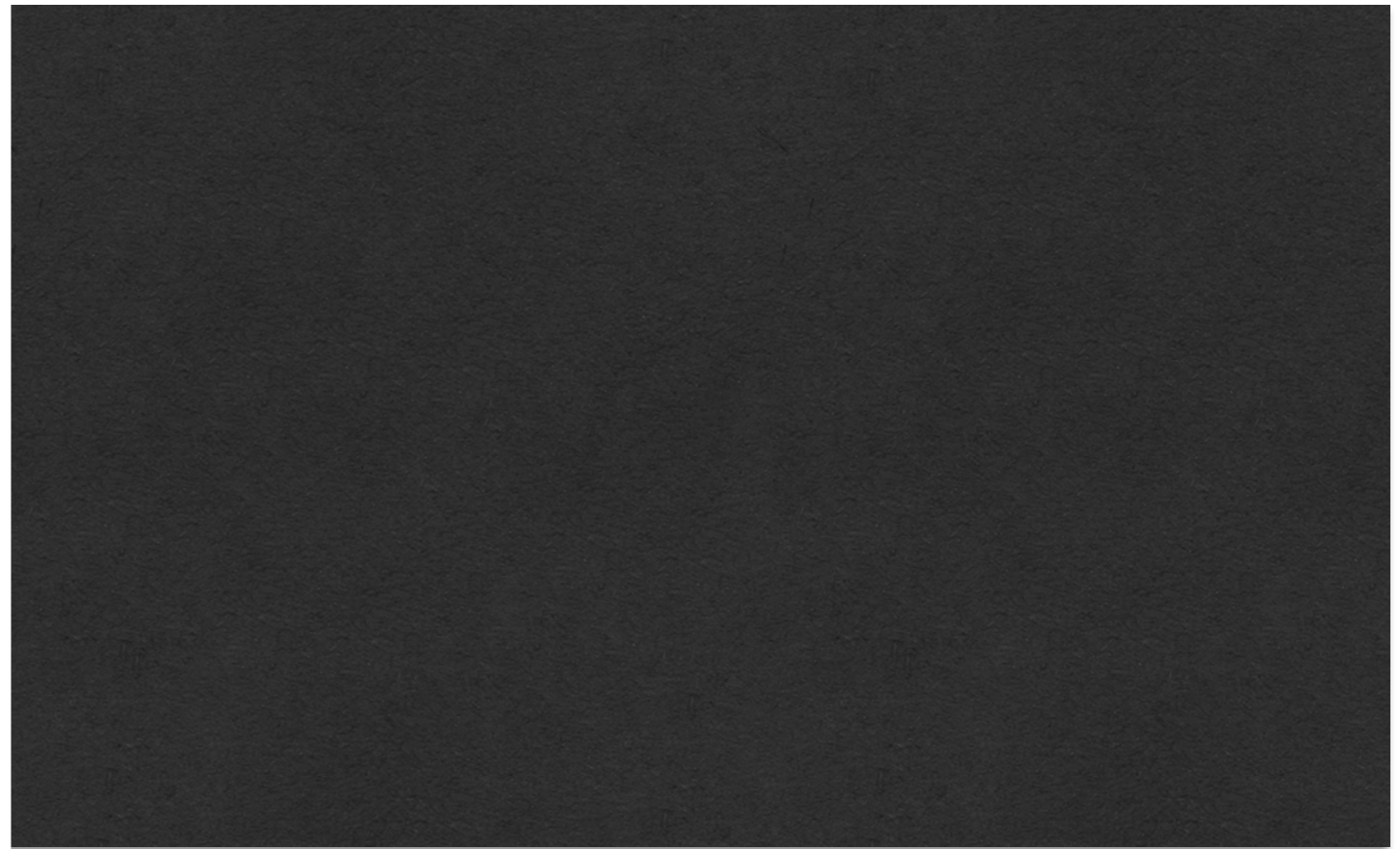
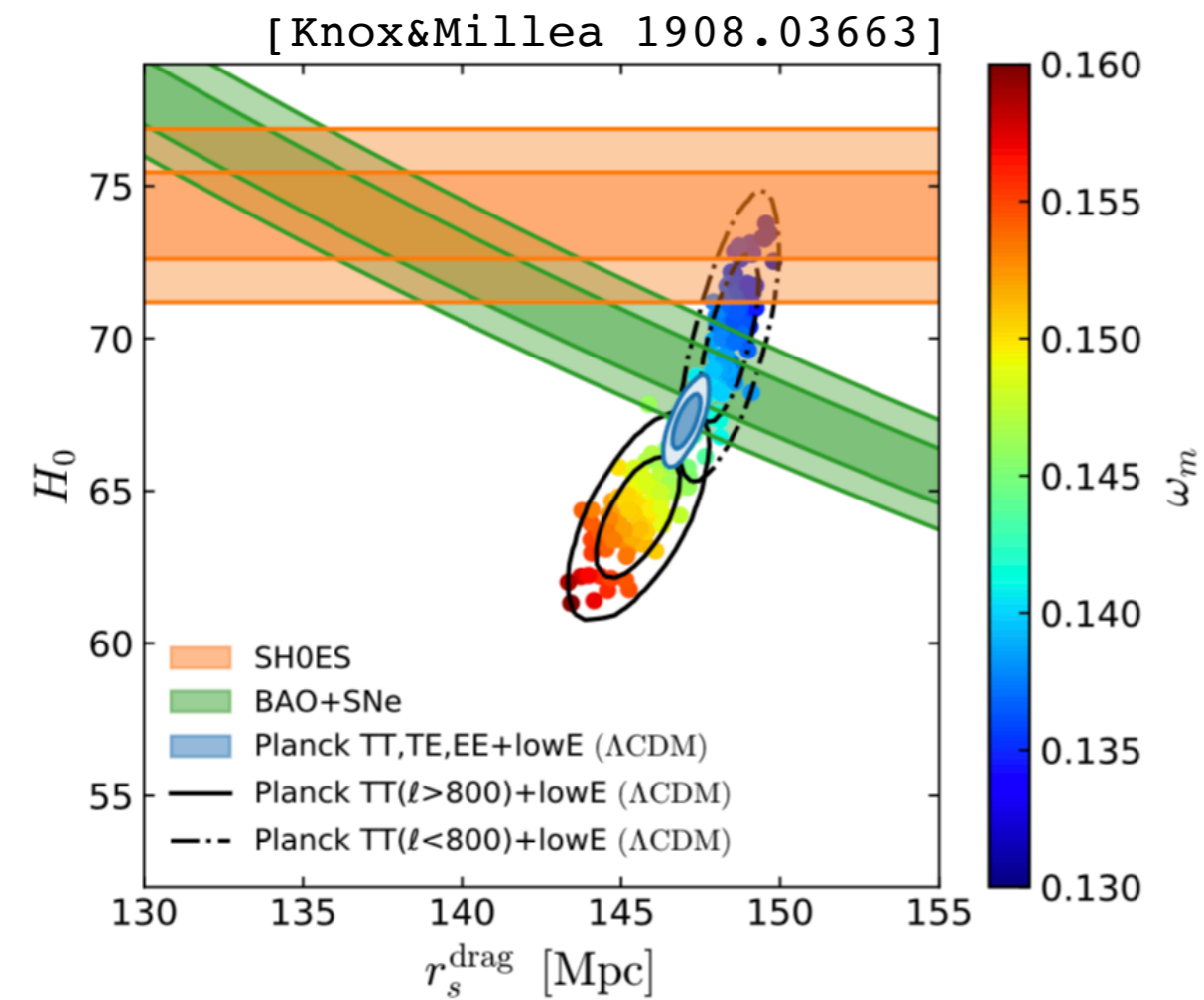
in collaboration with:  
Martin S. Sloth (CP3–Origins, DK)

arXiv:1910.10739 (PRD letter)  
arXiv:2006.06686 (PRD)  
arXiv:2009.00006 (PRD)  
arXiv:2112.00770 (PRD)  
arXiv:2112.00759

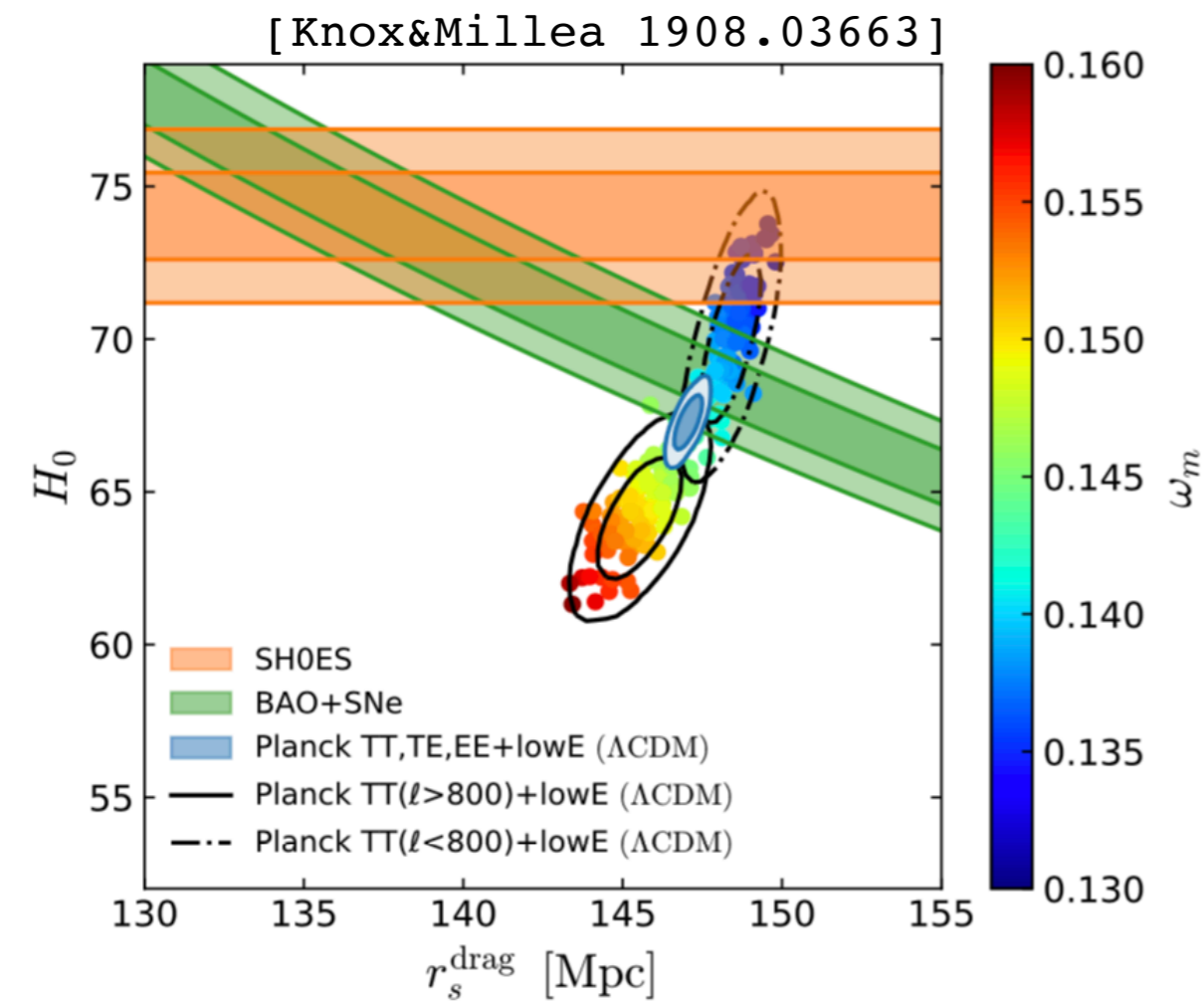


SW XIV Cargèse  
8–14 May 2022

# The Hubble tension

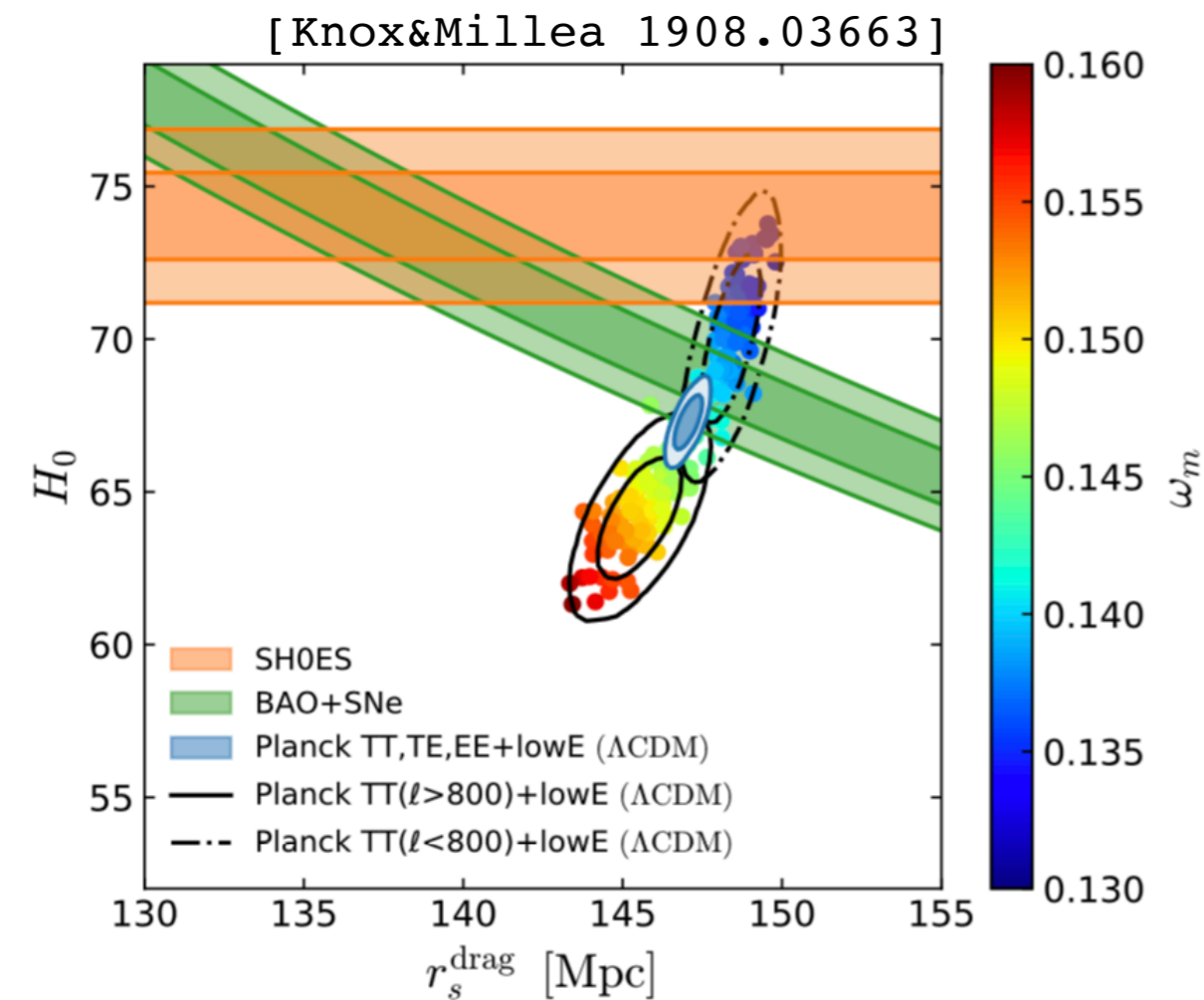


# The Hubble tension



(i) model-dependent statement:  
 SH0ES + Planck 5 sigma discrepant

# The Hubble tension



(i) model-dependent statement:  
 SH0ES + Planck 5 sigma discrepant

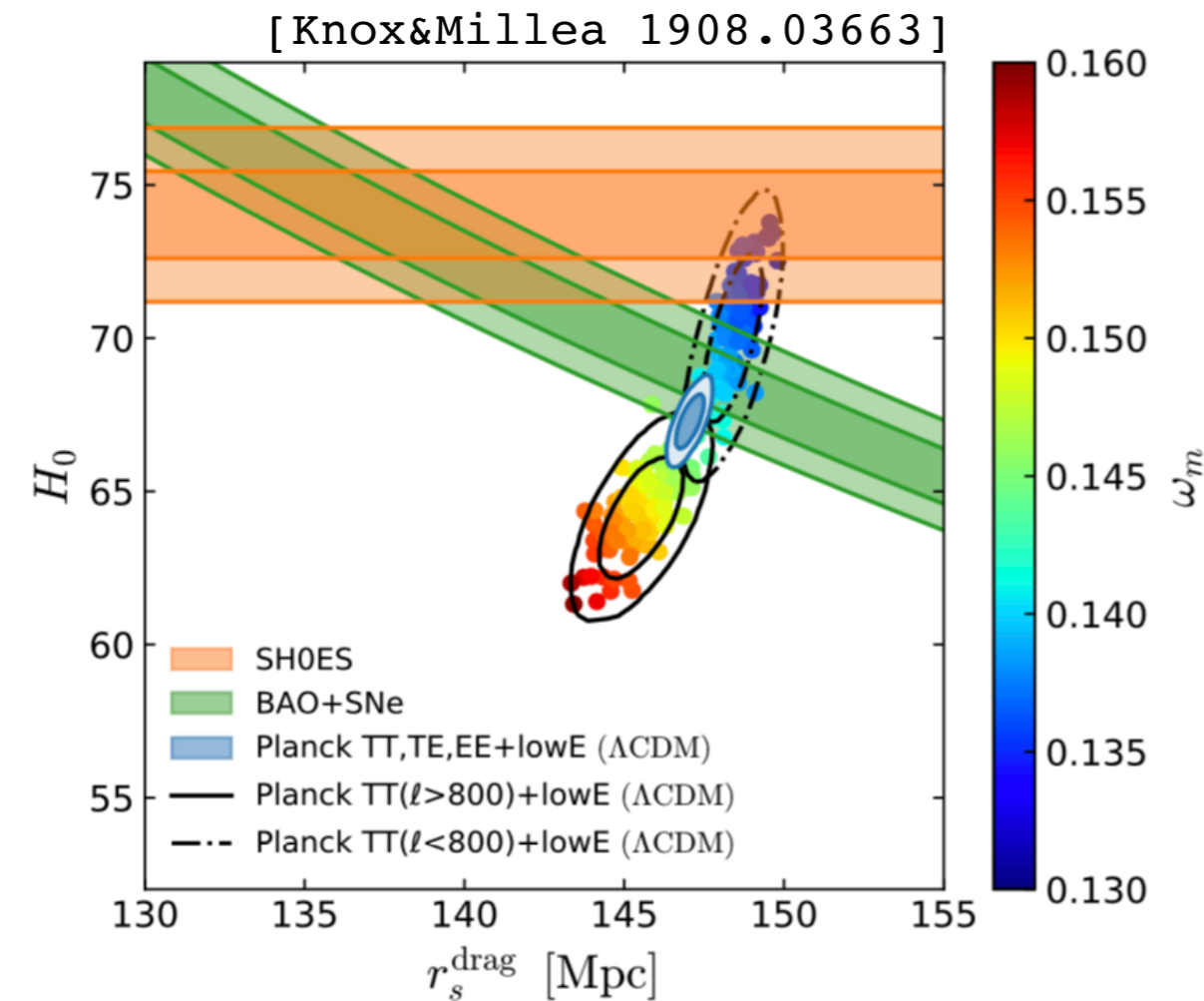
(ii) model-independent statement:

BAO + SNe:  $H_0 r_s \simeq \text{const}$

$$H_0 \nearrow \rightarrow r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} \searrow$$

[Aylor++1811.00537] [Arendse++1909.07986]

# The Hubble tension



(i) model-dependent statement:  
 SH0ES + Planck 5 sigma discrepant

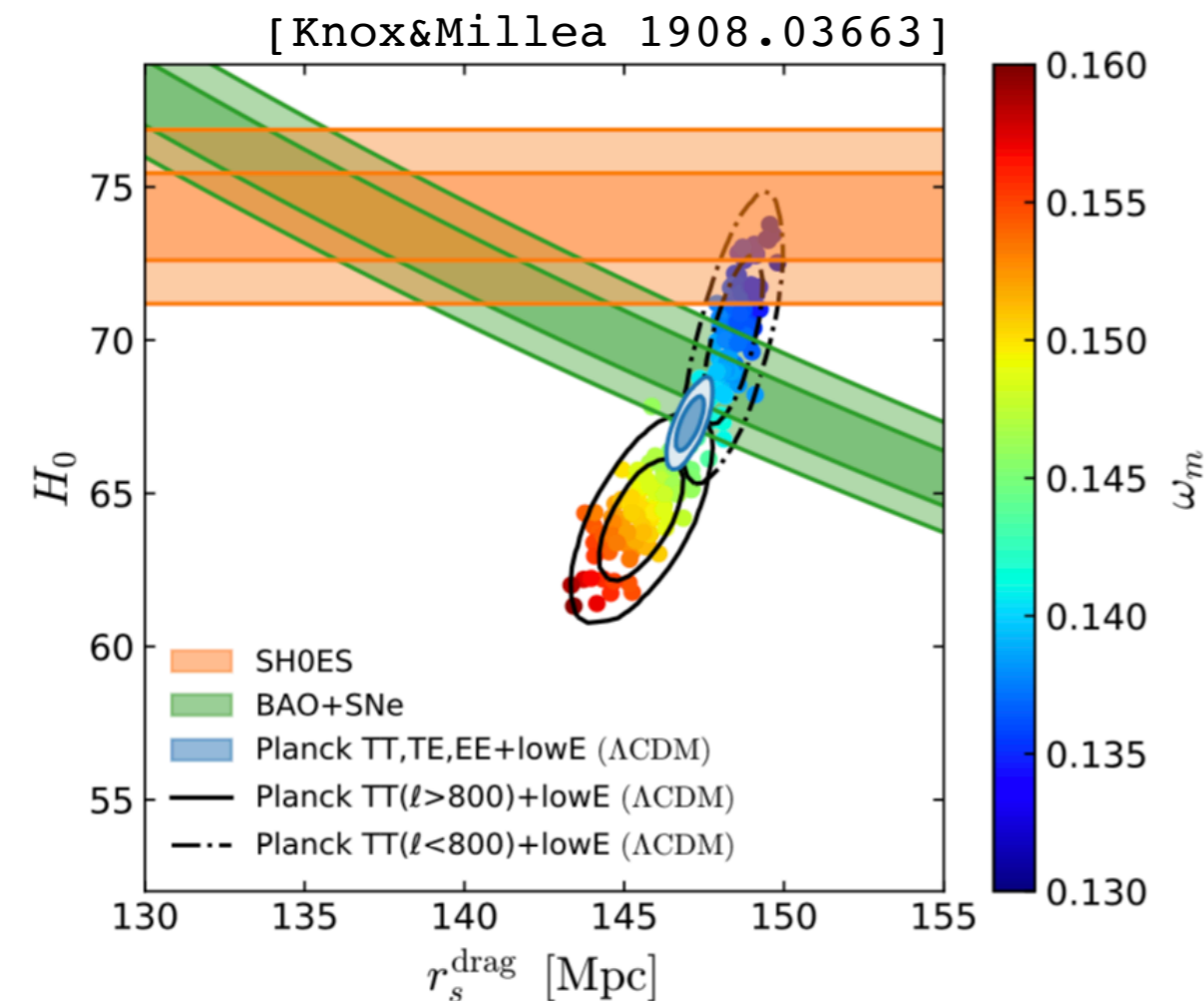
(ii) model-independent statement:  
 BAO + SNe:  $H_0 r_s \simeq \text{const}$

$H_0 \nearrow \rightarrow r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} \searrow$

[Aylor++1811.00537] [Arendse++1909.07986]

► Resolving the tension requires lowering the sound horizon by  $\sim 8$  Mpc.

# The Hubble tension



(i) model-dependent statement:  
**SH0ES + Planck** 5 sigma discrepant

(ii) model-independent statement:  
**BAO + SNe:**  $H_0 r_s \simeq \text{const}$

$H_0 \nearrow \rightarrow r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} \searrow$

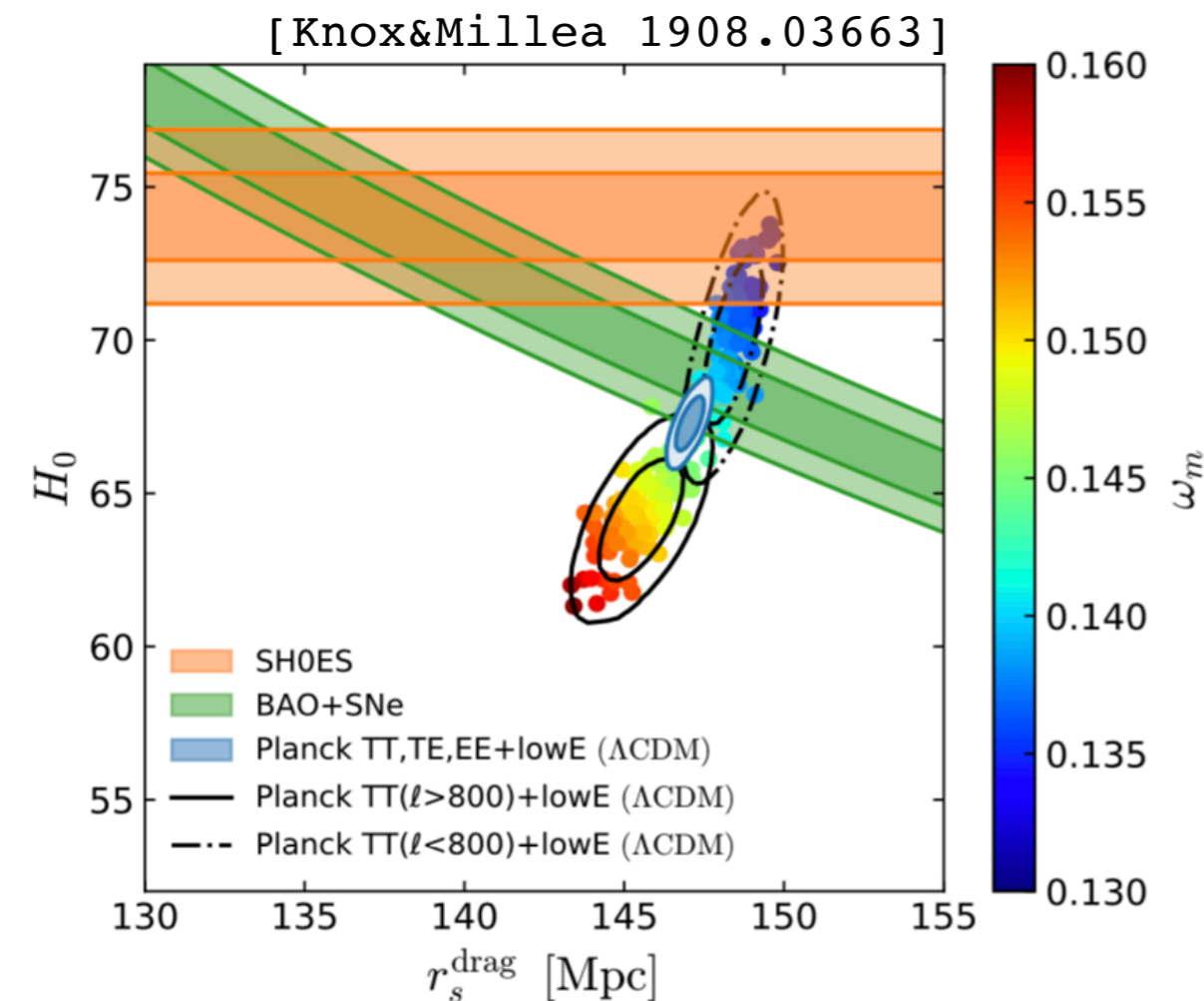
[Aylor++1811.00537] [Arendse++1909.07986]

- Resolving the tension requires lowering the sound horizon by  $\sim 8$  Mpc.
- This clearly suggests new physics pre recombination in redshift window:

Modify history of universe when  
highly constrained!

$$1000 < z < 25000$$

# The Hubble tension



(i) model-dependent statement:  
**SH0ES + Planck** 5 sigma discrepant

(ii) model-independent statement:  
**BAO + SNe:**  $H_0 r_s \simeq \text{const}$

$H_0 \nearrow \rightarrow r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} \searrow$

[Aylor++1811.00537] [Arendse++1909.07986]

- Resolving the tension requires lowering the sound horizon by  $\sim 8$  Mpc.
- This clearly suggests new physics pre recombination in redshift window:  
 Modify history of universe when highly constrained!
- Challenge: The new physics should preserve good fit to CMB observables.

$$1000 < z < 25000$$

# How to shorten sound horizon?

# How to shorten sound horizon?

- ▶ Energy injection **before** recombination (but not too early!).

# How to shorten sound horizon?

► Energy injection **before** recombination (but not too early!).

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3 + \Omega_r(1+z)^4 + \Omega_X(z)} \quad \leftarrow \text{new component } (\sim 10\%)$$

→ increases  $H(z)$  prior to recombination

$$\rightarrow r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} \quad \searrow$$

# How to shorten sound horizon?

► Energy injection **before** recombination (but not too early!).

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3 + \Omega_r(1+z)^4 + \Omega_X(z)} \quad \leftarrow \text{new component } (\sim 10\%)$$

→ increases  $H(z)$  prior to recombination

$$\rightarrow r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} \quad \searrow$$

► Balanced by increasing  $H_0$

# How to shorten sound horizon?

- ▶ Energy injection **before** recombination (but not too early!).

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3 + \Omega_r(1+z)^4 + \Omega_X(z)} \quad \leftarrow \text{new component (\sim 10\%)}$$

→ increases  $H(z)$  prior to recombination

$$r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz \quad \searrow$$

- ▶ Balanced by increasing  $H_0$
- ▶ Subsequently, the new component has to decay at least as fast as radiation.

# How to shorten sound horizon?

- ▶ Energy injection **before** recombination (but not too early!).

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3 + \Omega_r(1+z)^4 + \Omega_X(z)}$$

← new component (~10%)

→ increases  $H(z)$  prior to recombination

$$r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz$$

↓

- ▶ Balanced by increasing  $H_0$
- ▶ Subsequently, the new component has to decay at least as fast as radiation.
- ▶ Canonical example: **Dark Radiation (DR)**

# How to shorten sound horizon?

- Energy injection **before** recombination (but not too early!).

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3 + \Omega_r(1+z)^4 + \Omega_X(z)}$$

← new component (~10%)

→ increases  $H(z)$  prior to recombination

$$r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz$$

↓

- Balanced by increasing  $H_0$
- Subsequently, the new component has to decay at least as fast as radiation.
- Canonical example: **Dark Radiation (DR)**

Dark radiation:  $\Omega_X(t) = \Omega_{\text{DR}} a(t)^{-4}$

Simplest implementation: promote  $N_{\text{eff}}$  to free parameter

# How to shorten sound horizon?

- Energy injection **before** recombination (but not too early!).

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3 + \Omega_r(1+z)^4 + \Omega_X(z)}$$

← new component (~10%)

→ increases  $H(z)$  prior to recombination

$$r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz$$

↓

- Balanced by increasing  $H_0$
- Subsequently, the new component has to decay at least as fast as radiation.
- Canonical example: **Dark Radiation (DR)**

Dark radiation:  $\Omega_X(t) = \Omega_{\text{DR}} a(t)^{-4}$

Simplest implementation: promote  $N_{\text{eff}}$  to free parameter

- **Result:** Tension only reduced slightly, still ~4 sigma

[Planck 2018 + BAO (+LSS)  
+ Pantheon + BBN]

# How to shorten sound horizon?

- Energy injection **before** recombination (but not too early!).

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3 + \Omega_r(1+z)^4 + \Omega_X(z)} \quad \leftarrow \text{new component } (\sim 10\%)$$

► increases  $H(z)$  prior to recombination

$$r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz \quad \searrow$$

- Balanced by increasing  $H_0$
- Subsequently, the new component has to decay at least as fast as radiation.
- Canonical example: **Dark Radiation (DR)**

Dark radiation:  $\Omega_X(t) = \Omega_{\text{DR}} a(t)^{-4}$

Simplest implementation: promote  $N_{\text{eff}}$  to free parameter

- **Result:** Tension only reduced slightly, still  $\sim 4$  sigma [Planck 2018 + BAO (+LSS) + Pantheon + BBN]
- **Problem:** Too much diffusion damping on small scales.
- Generalisations where DR constituents becomes non-relativistic around eV scale and annihilate are more promising (see Majoron and “step” proposal).

# Early Dark Energy

[Karwal et al., 2016]

[Smith et al., 2019]

# Early Dark Energy

[Karwal et al., 2016]

[Smith et al., 2019]

2nd order  
transition

$$\phi^{2n}$$

$$n \geq 2$$

(ii)



(i)

(i) frozen due to Hubble  
friction until matter-rad. eq.

(ii) coherent oscillations

~~$\phi^2$~~

~~$\phi^4$~~

$\phi^6$

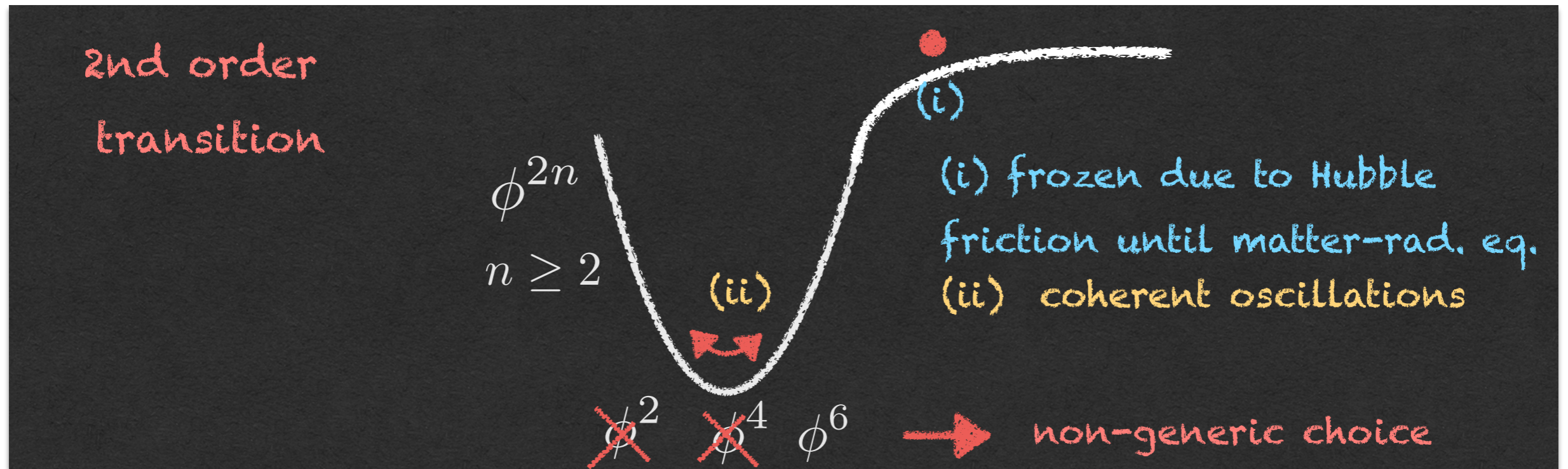


non-generic choice

# Early Dark Energy

[Karwal et al., 2016]

[Smith et al., 2019]



## ► Phenomenology requires:

- Flattening of potential at high field values.
- Oscillations in anharmonic potential.
- Ultralight effective mass :  $m \sim 10^{-27} \text{eV}$

cycle-averaged:

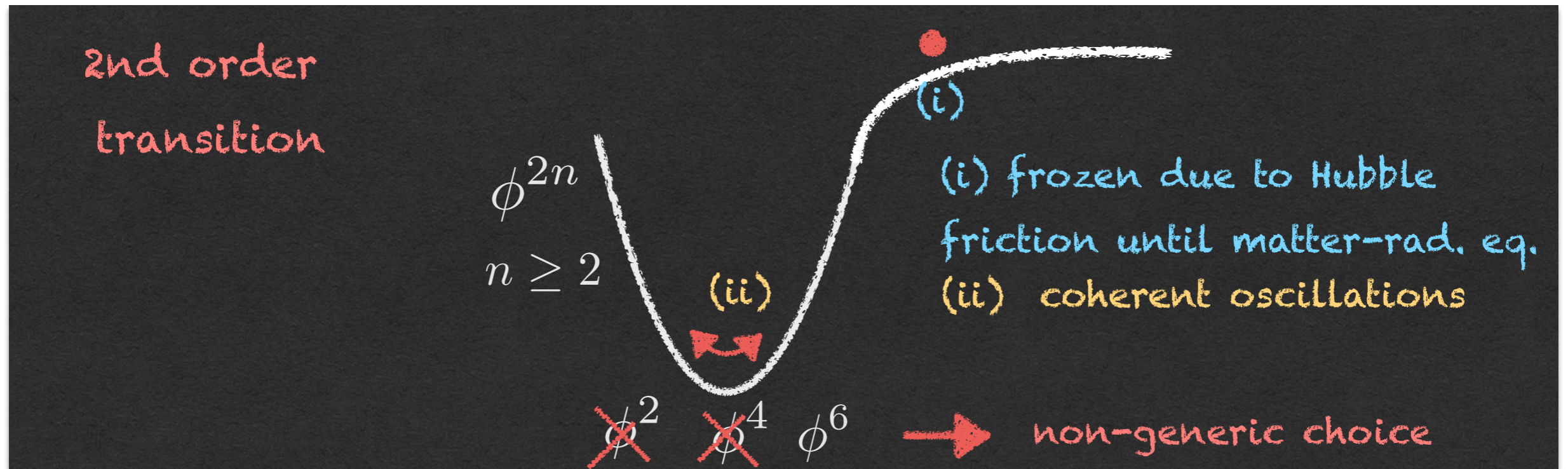
$$\Omega_X(t) \simeq \begin{cases} \Omega_{\text{EDE}} \\ \Omega_{\text{EDE}} [a(t_*)/a(t)]^\alpha \end{cases}$$

where  $\alpha = 3 \left( 1 + \frac{n-1}{n+1} \right)$

# Early Dark Energy

[Karwal et al., 2016]

[Smith et al., 2019]



## ► Phenomenology requires:

- Flattening of potential at high field values.
- Oscillations in anharmonic potential.
- Ultralight effective mass :  $m \sim 10^{-27} \text{eV}$

## ► Brings tension down to $\sim 2.5$ sigma (1p-EDE).

[Planck 2018 + BAO (+LSS) + Pantheon + BBN]

cycle-averaged:

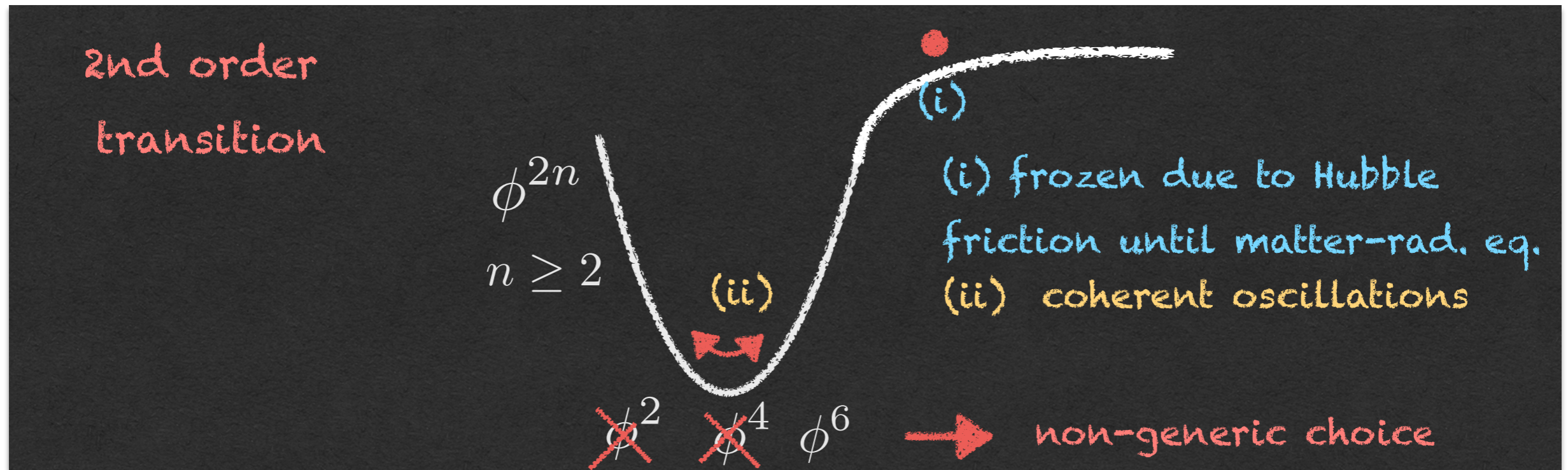
$$\Omega_X(t) \simeq \begin{cases} \Omega_{\text{EDE}} \\ \Omega_{\text{EDE}} [a(t_*)/a(t)]^\alpha \end{cases}$$

where  $\alpha = 3 \left( 1 + \frac{n-1}{n+1} \right)$

# Early Dark Energy

[Karwal et al., 2016]

[Smith et al., 2019]



## ► Phenomenology requires:

- Flattening of potential at high field values.
- Oscillations in anharmonic potential.
- Ultralight effective mass :  $m \sim 10^{-27} \text{eV}$

## ► Brings tension down to $\sim 2.5$ sigma (1p-EDE).

[Planck 2018 + BAO (+LSS) + Pantheon + BBN]

## ► Challenges:

- How to justify choice  $\phi^6$ ?
- Not resolving S8 tension (does not make it much worse though).

[Smith++, 2009.10740] [Amico++, 2006.12420]

[Murgia++, 2009.10733] [Hill++, 2003.07355]

cycle-averaged:

$$\Omega_X(t) \simeq \begin{cases} \Omega_{\text{EDE}} \\ \Omega_{\text{EDE}} [a(t_*)/a(t)]^\alpha \end{cases}$$

where  $\alpha = 3 \left( 1 + \frac{n-1}{n+1} \right)$

# First order decay scenario

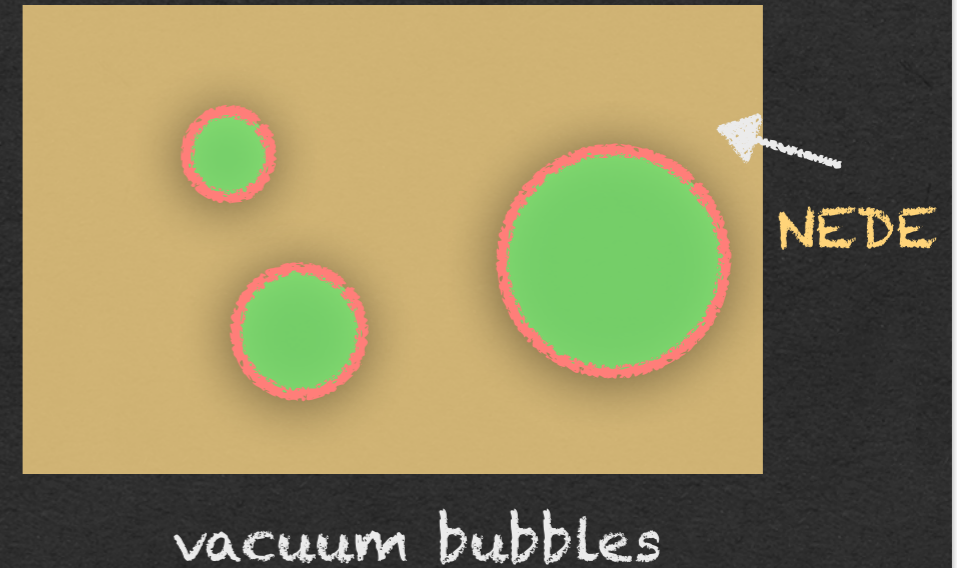
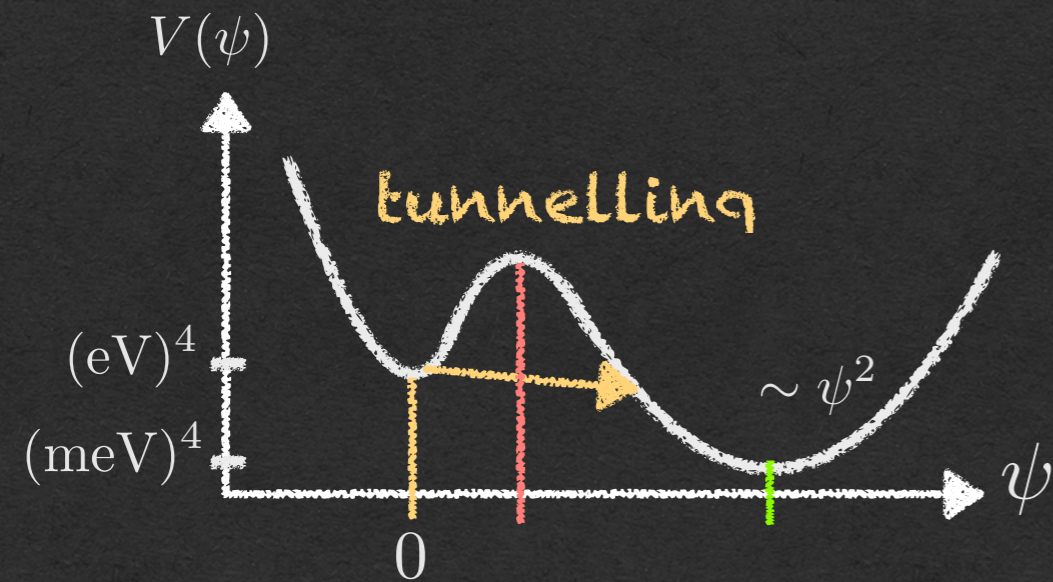
# First order decay scenario

► **Question:** Can we find a model more motivated from particle physics?

# First order decay scenario

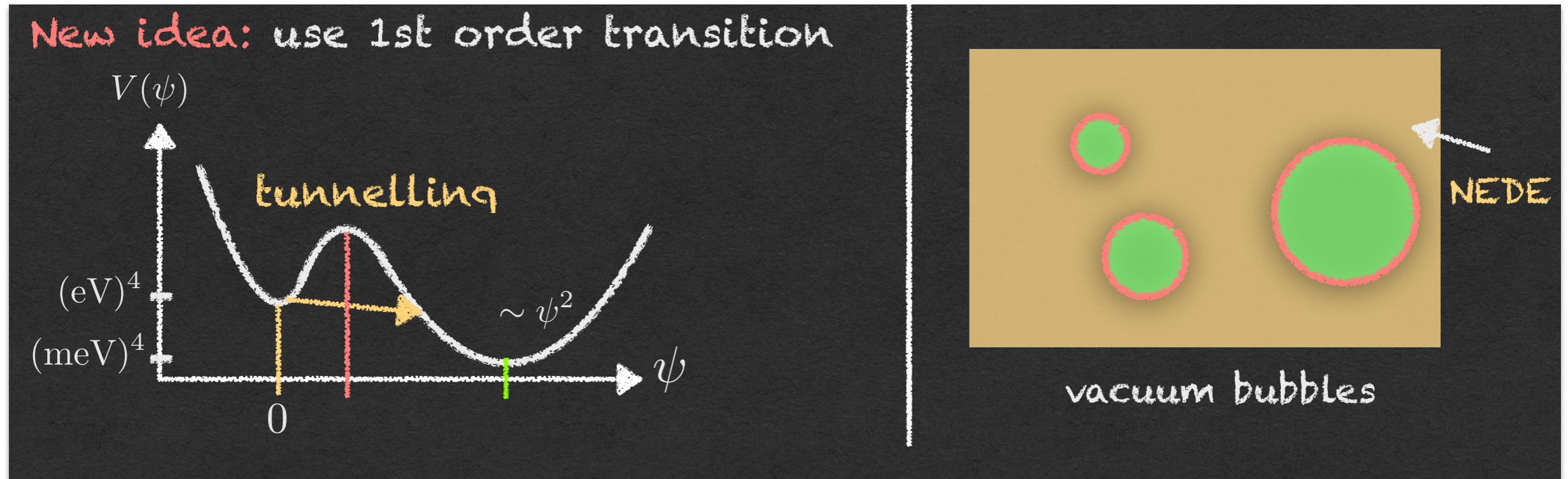
► **Question:** Can we find a model more motivated from particle physics?

**New idea:** use 1st order transition



# First order decay scenario

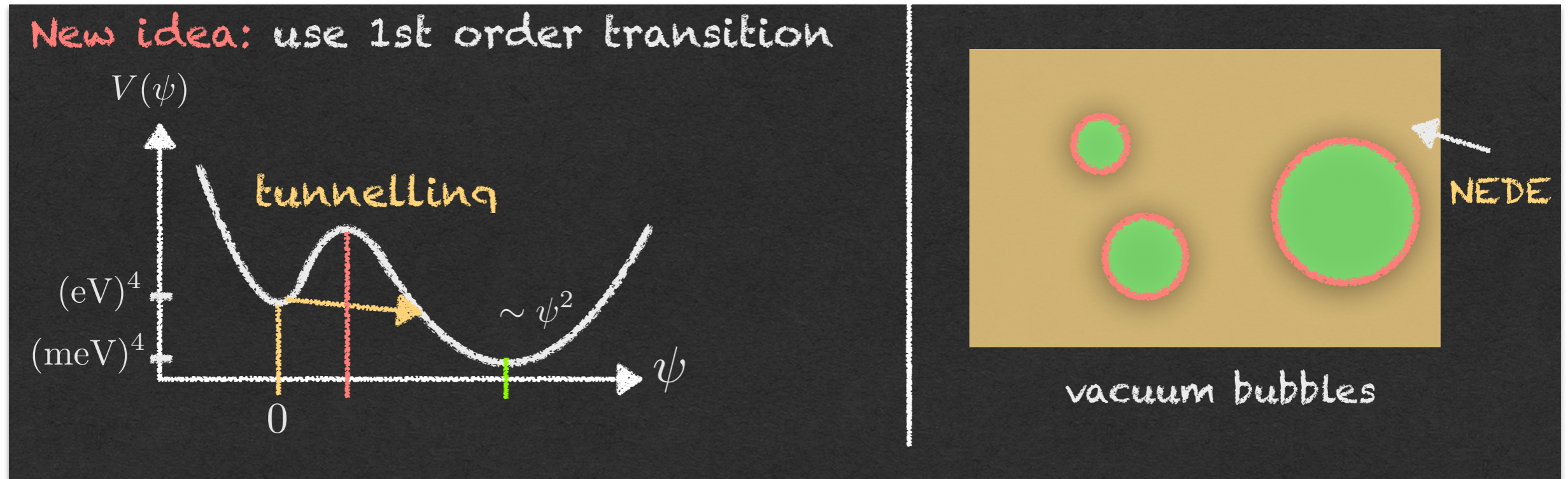
► **Question:** Can we find a model more motivated from particle physics?



► **Hubble tension:** EDE/NEDE provided by (decaying) false vacuum energy.

# First order decay scenario

► **Question:** Can we find a model more motivated from particle physics?

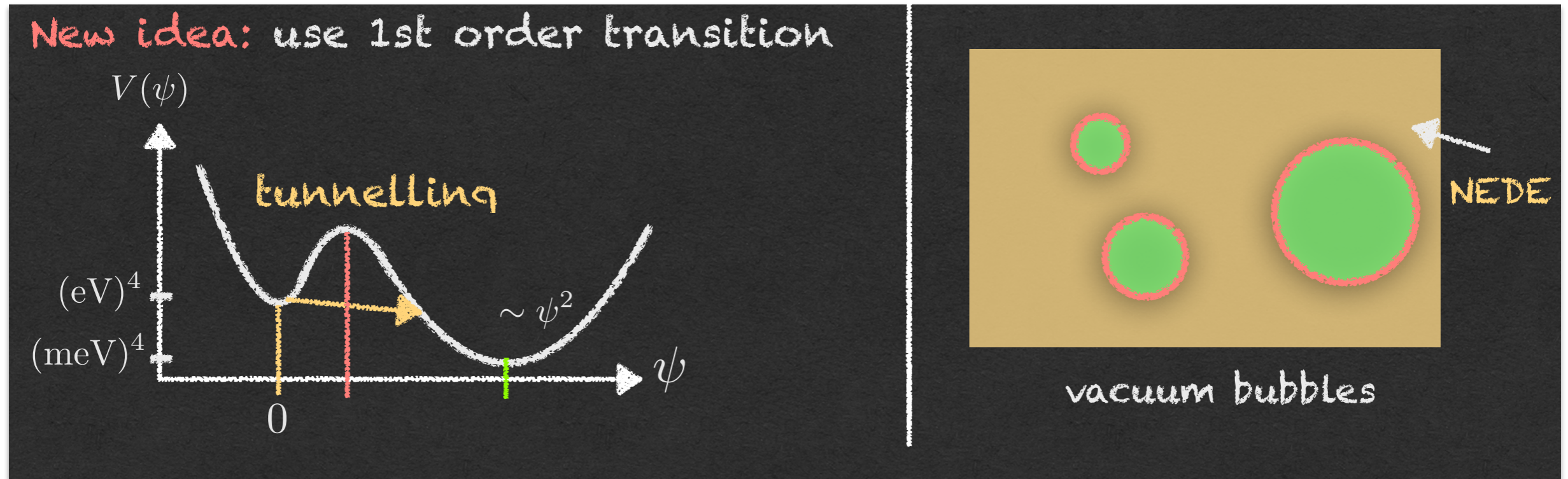


► **Hubble tension:** EDE/NEDE provided by (decaying) false vacuum energy.

► This idea faces challenges:

# First order decay scenario

► **Question:** Can we find a model more motivated from particle physics?



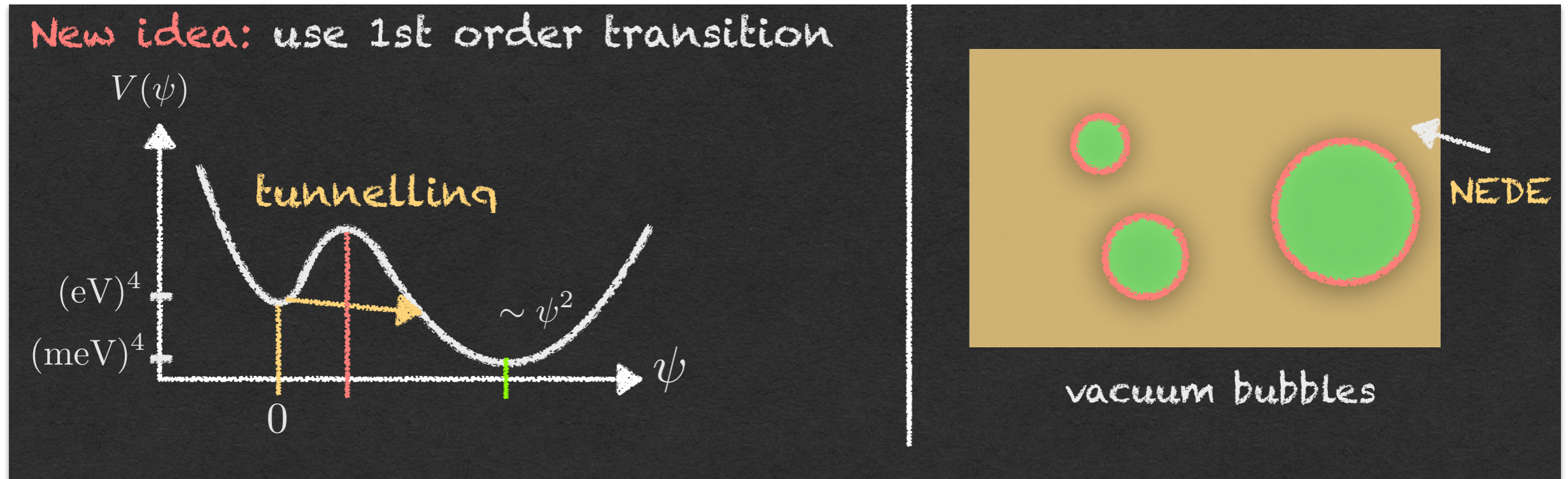
► **Hubble tension:** EDE/NEDE provided by (decaying) false vacuum energy.

► This idea faces challenges:

1. Decay should happen around matter–radiation equality (lesson from EDE).

# First order decay scenario

► **Question:** Can we find a model more motivated from particle physics?



► **Hubble tension:** EDE/NEDE provided by (decaying) false vacuum energy.

► This idea faces challenges:

1. Decay should happen around matter–radiation equality (lesson from EDE).
2. Bubble percolation has to be extremely efficient to avoid anisotropies.  
(prevented from growing to cosmological size).

Cold New Early Dark Energy

# Cold New Early Dark Energy

► Introduce a **trigger field** to synchronise decay.

# Cold New Early Dark Energy

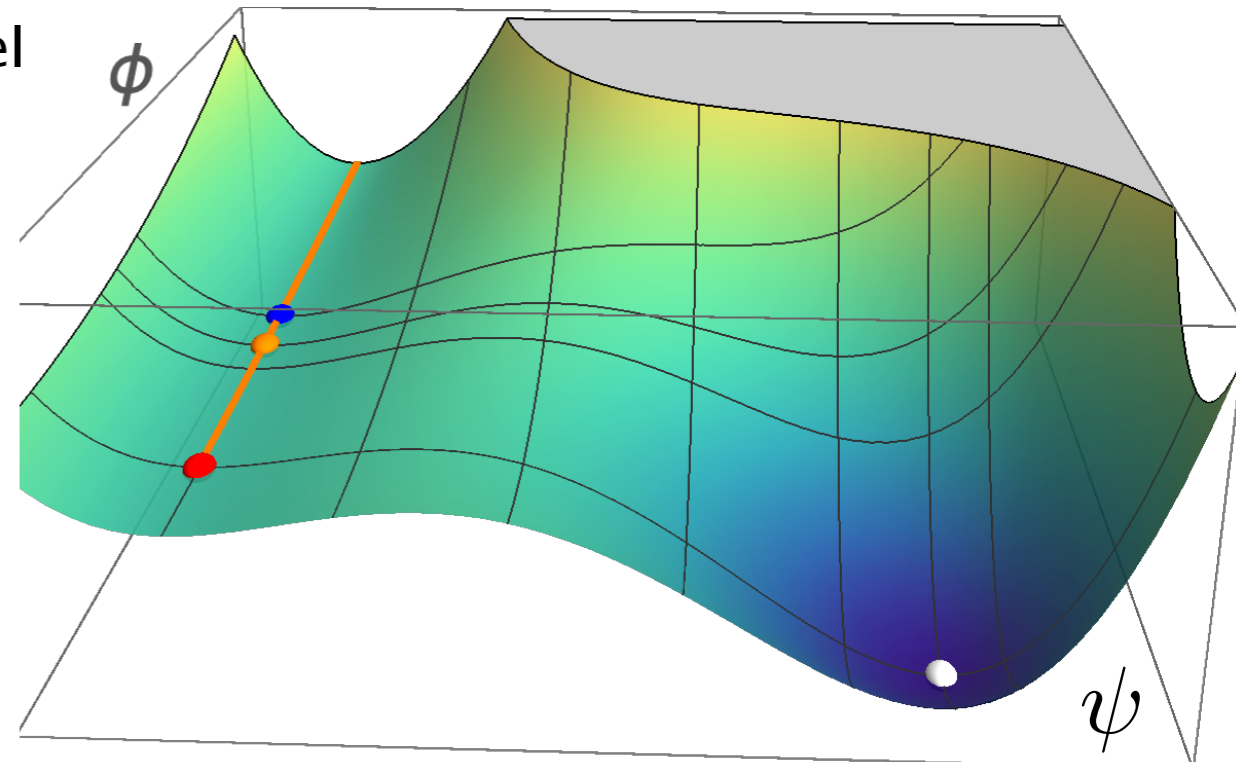
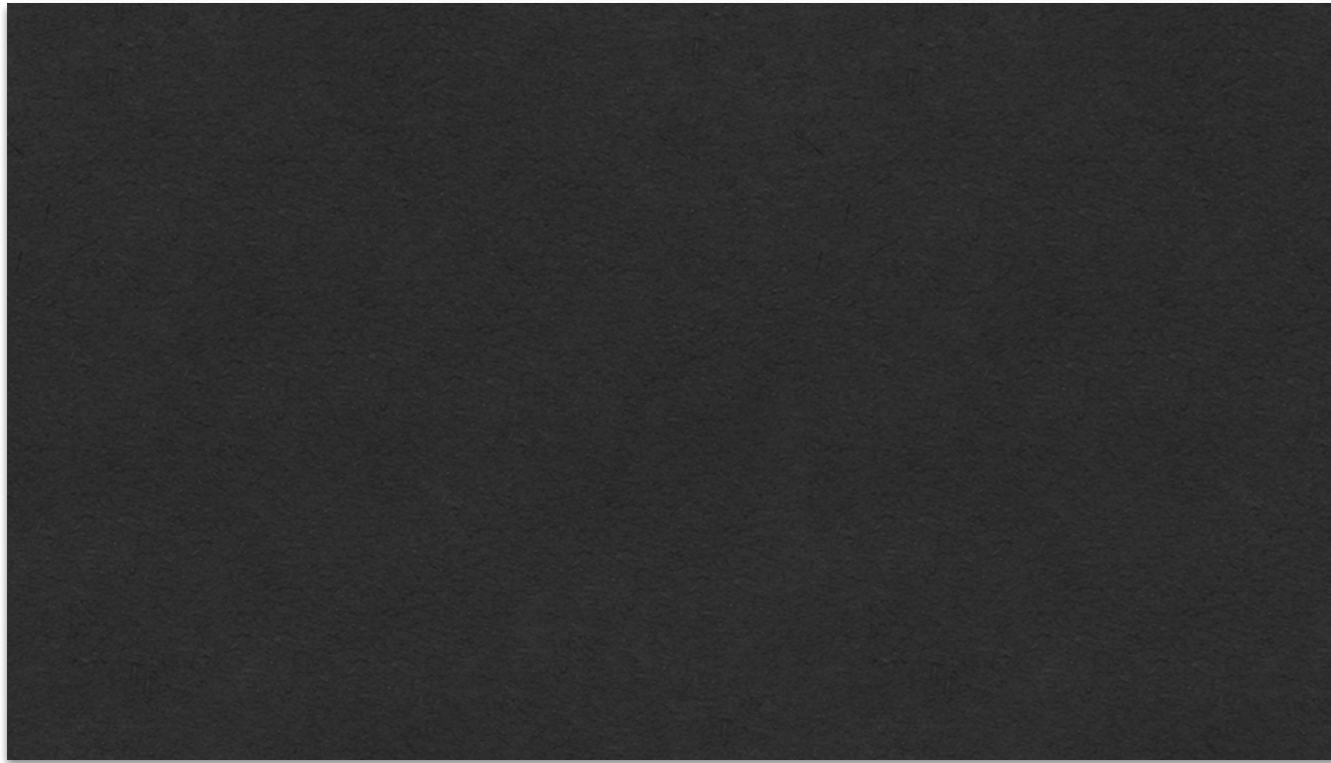
- ▶ Introduce a **trigger field** to synchronise decay.
- ▶ eV scale adaption of first-order inflationary model

[Linde, 1990][Adams, Freese, 1990]

# Cold New Early Dark Energy

- ▶ Introduce a **trigger field** to synchronise decay.
- ▶ eV scale adaption of first-order inflationary model

[Linde, 1990][Adams, Freese, 1990]



# Cold New Early Dark Energy

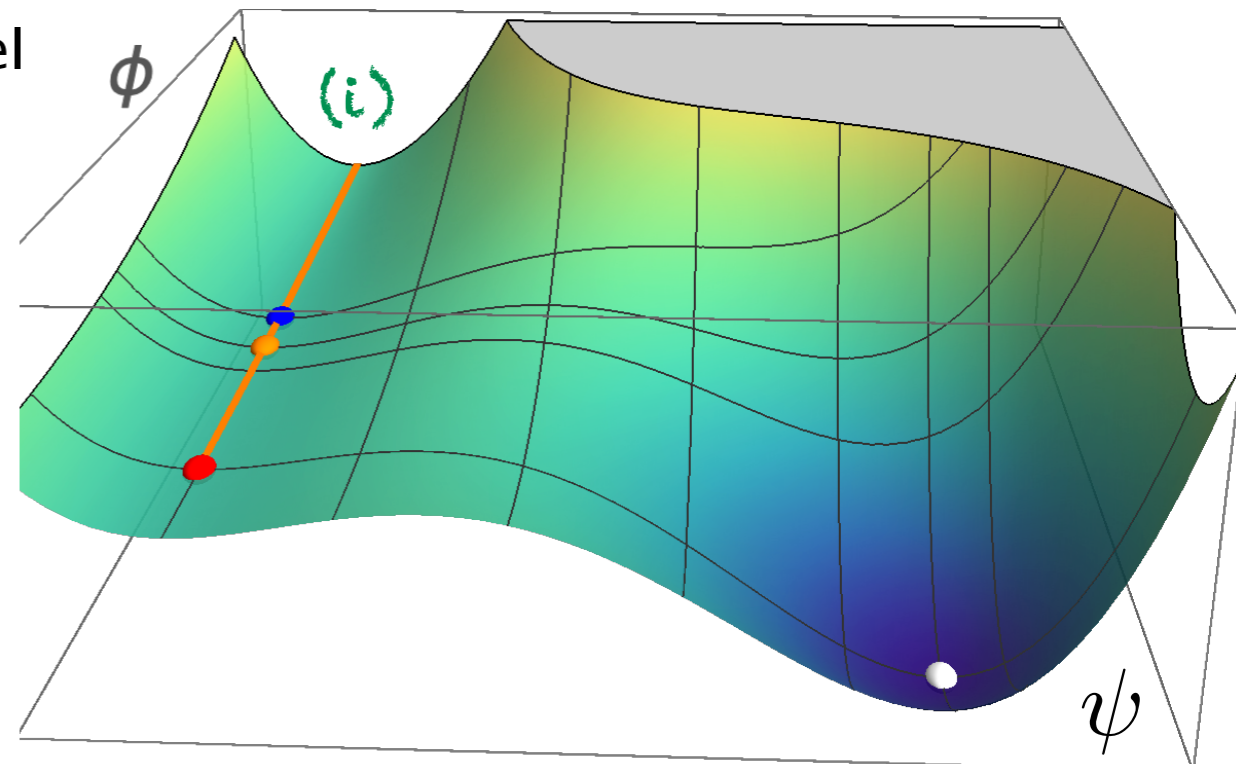
- Introduce a **trigger field** to synchronise decay.
- eV scale adaption of first-order inflationary model

[Linde, 1990][Adams, Freese, 1990]

(i)  $\phi \simeq \phi_{ini}$

Field stuck in false minimum

$\Gamma/H^4 \ll 1$  suppressed tunnelling rate



# Cold New Early Dark Energy

- Introduce a **trigger field** to synchronise decay.
- eV scale adaption of first-order inflationary model

[Linde, 1990][Adams, Freese, 1990]

(i)  $\phi \simeq \phi_{ini}$

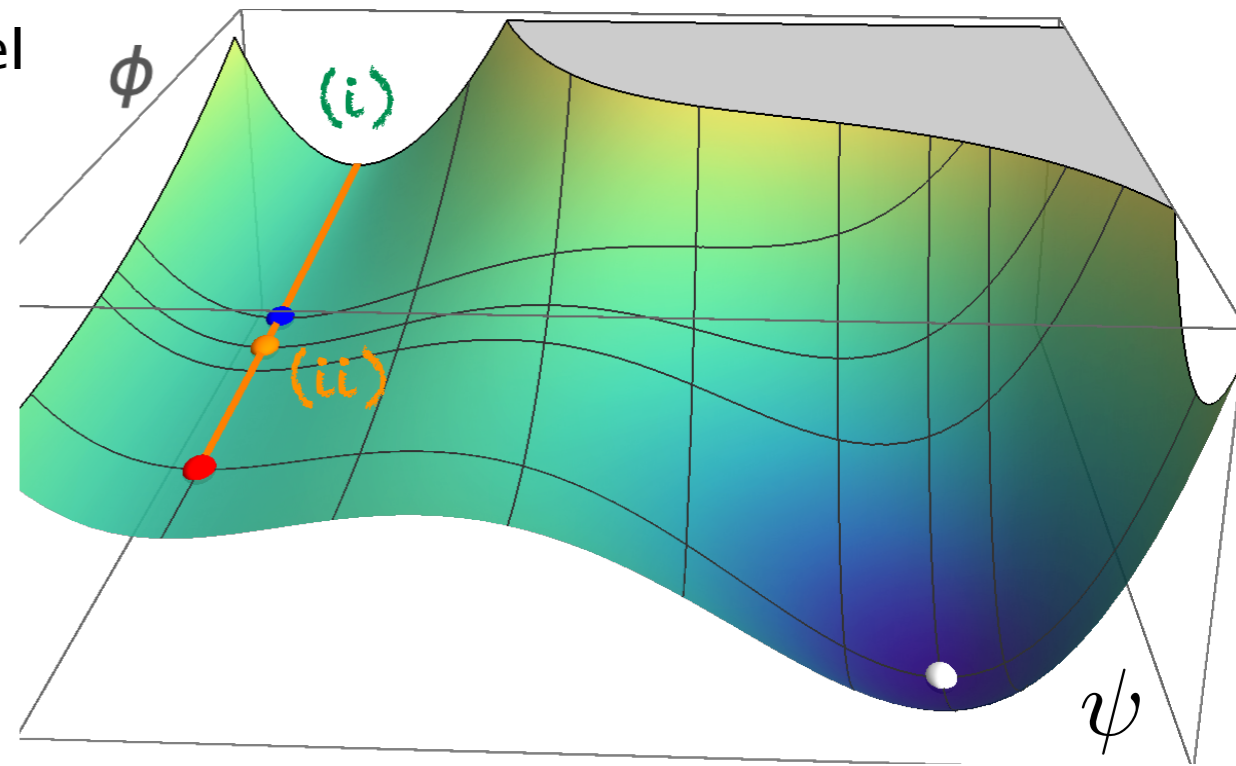
Field stuck in false minimum

$\Gamma/H^4 \ll 1$  suppressed tunnelling rate

(ii)  $\phi$  starts evolving

after **orange** dot:  $\Gamma/H^4 \gtrsim 1$

→ strong nucleation event



# Cold New Early Dark Energy

- Introduce a **trigger field** to synchronise decay.
- eV scale adaption of first-order inflationary model

[Linde, 1990][Adams, Freese, 1990]

(i)  $\phi \simeq \phi_{ini}$

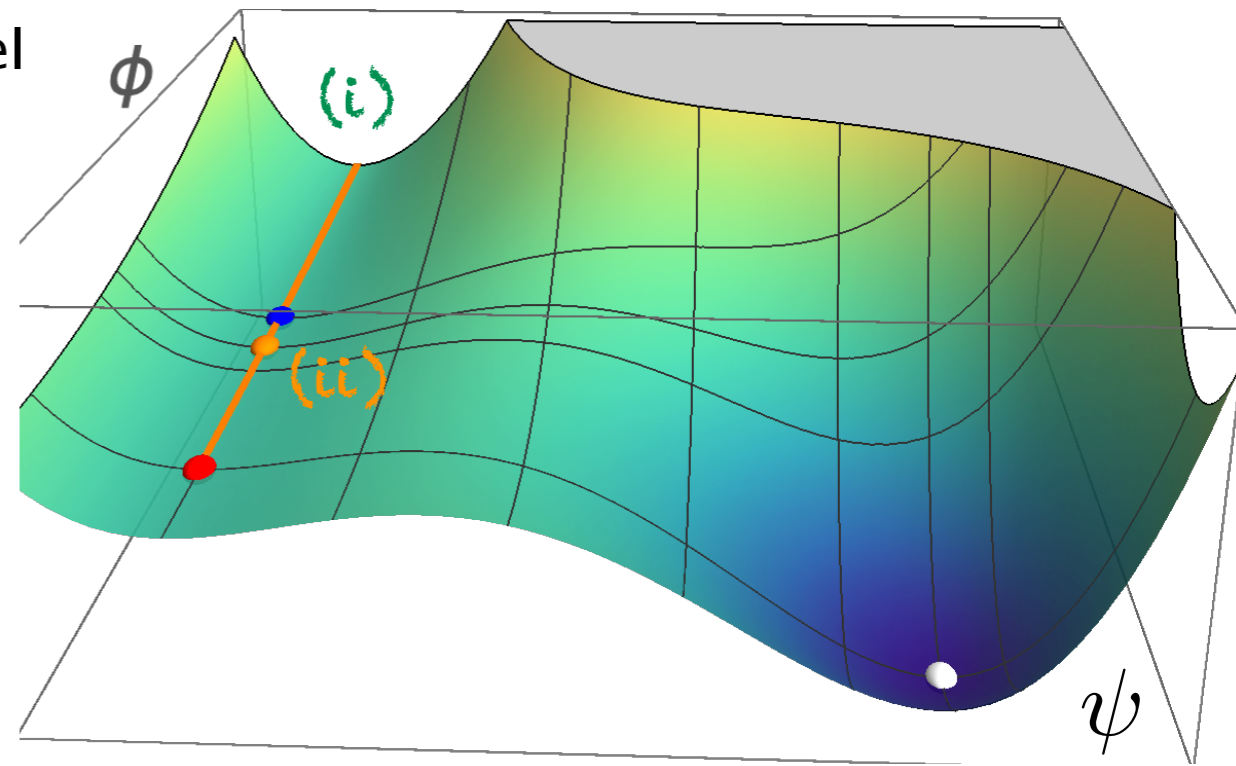
Field stuck in false minimum

$\Gamma/H^4 \ll 1$  suppressed tunnelling rate

(ii)  $\phi$  starts evolving

after **orange** dot:  $\Gamma/H^4 \gtrsim 1$

→ **strong nucleation event**



$$V(\psi, \phi) = \frac{\lambda}{4} \psi^4 + \frac{1}{2} M^2 \psi^2 - \frac{1}{3} \alpha M \psi^3 + \frac{1}{2} m^2 \phi^2 + \frac{1}{2} \tilde{\lambda} \phi^2 \psi^2 \quad \alpha = \mathcal{O}(1)$$

# Cold New Early Dark Energy

- Introduce a **trigger field** to synchronise decay.
- eV scale adaption of first-order inflationary model

[Linde, 1990][Adams, Freese, 1990]

(i)  $\phi \simeq \phi_{ini}$

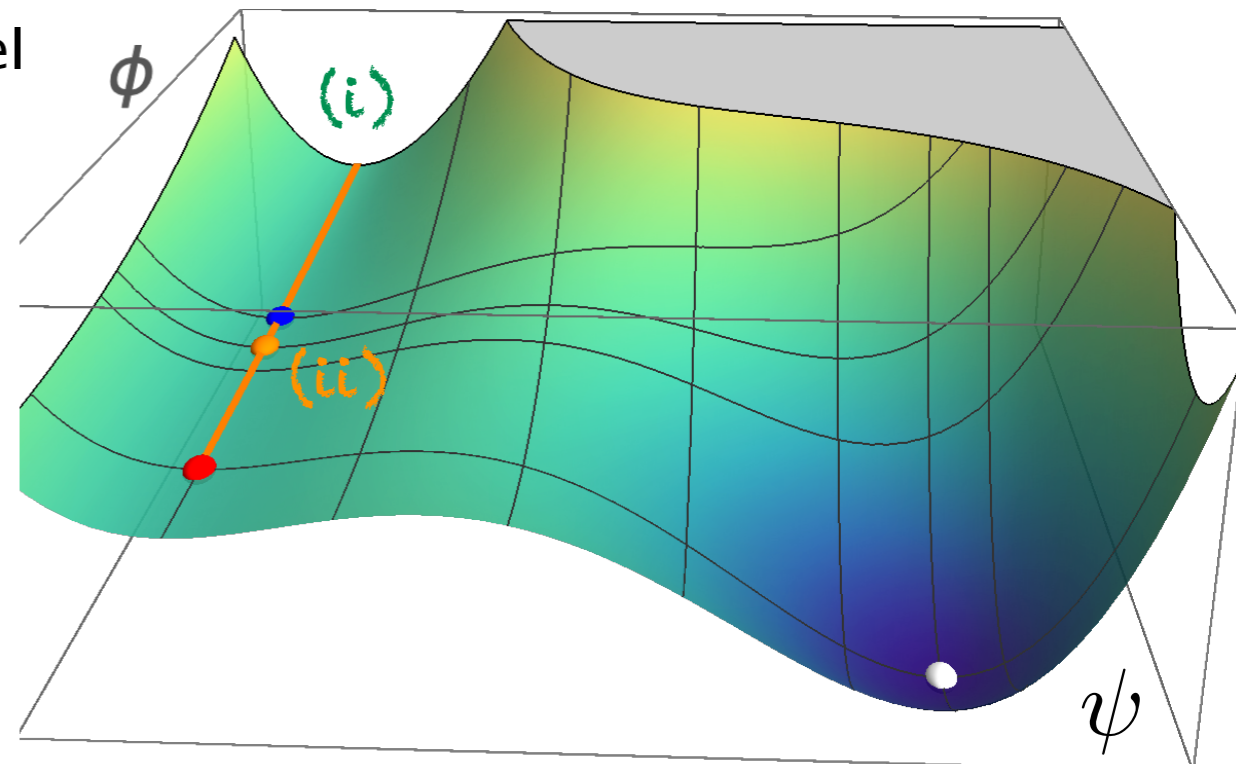
Field stuck in false minimum

$\Gamma/H^4 \ll 1$  suppressed tunnelling rate

(ii)  $\phi$  starts evolving

after **orange** dot:  $\Gamma/H^4 \gtrsim 1$

→ **strong nucleation event**



$$V(\psi, \phi) = \frac{\lambda}{4} \psi^4 + \frac{1}{2} M^2 \psi^2 - \frac{1}{3} \alpha M \psi^3 + \frac{1}{2} m^2 \phi^2 + \frac{1}{2} \tilde{\lambda} \phi^2 \psi^2 \quad \alpha = \mathcal{O}(1)$$

hierarchy:  $M \sim \text{eV} \gg m \sim 10^{-27} \text{eV}$  → **ultra-light physics**

radiative stability:  $\tilde{\lambda} \lesssim 10^3 m^2 / M^2 \ll 1$       weak coupling:  $\lambda < 0.1$

# Cold New Early Dark Energy

► Introduce a **trigger field** to synchronise decay.

► eV scale adaption of first-order inflationary model

[Linde, 1990] [Adams, Freese, 1990]

(i)  $\phi \simeq \phi_{ini}$

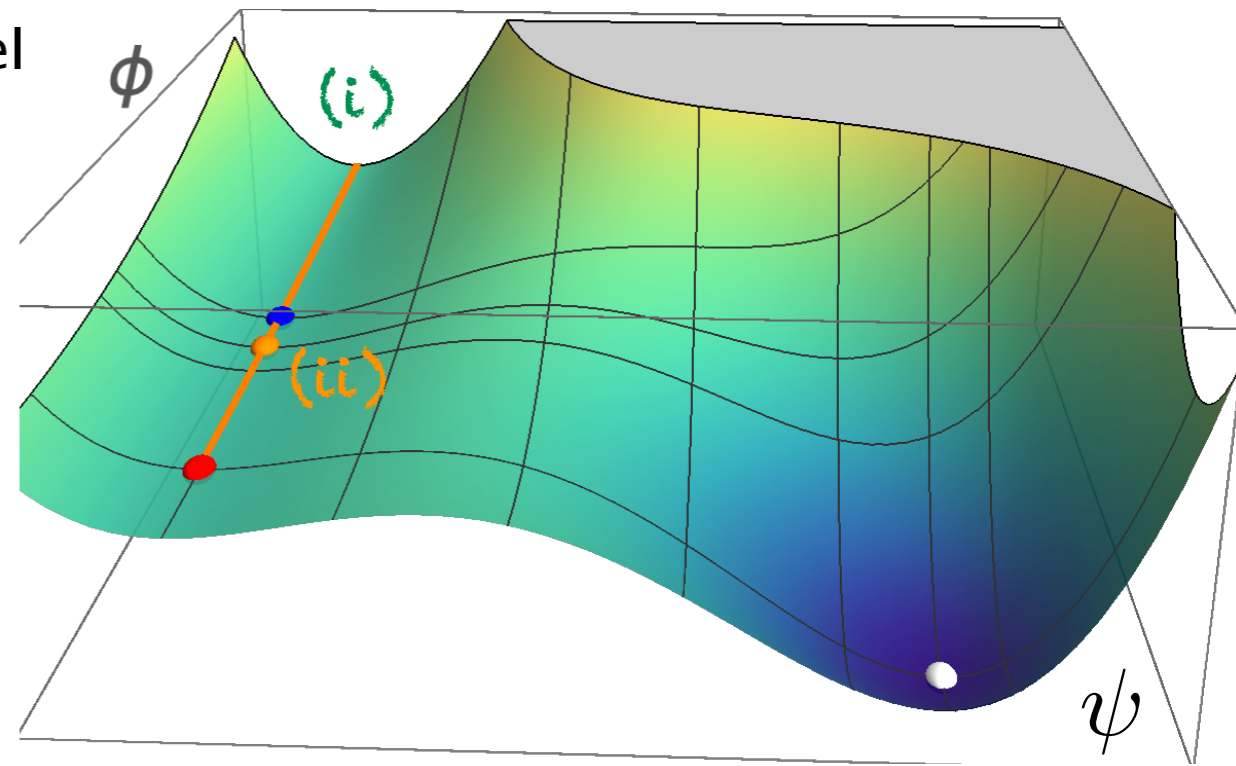
Field stuck in false minimum

$\Gamma/H^4 \ll 1$  suppressed tunnelling rate

(ii)  $\phi$  starts evolving

after **orange** dot:  $\Gamma/H^4 \gtrsim 1$

→ **strong nucleation event**



$$V(\psi, \phi) = \frac{\lambda}{4} \psi^4 + \frac{1}{2} M^2 \psi^2 - \frac{1}{3} \alpha M \psi^3 + \frac{1}{2} m^2 \phi^2 + \frac{1}{2} \tilde{\lambda} \phi^2 \psi^2 \quad \alpha = \mathcal{O}(1)$$

hierarchy:  $M \sim \text{eV} \gg m \sim 10^{-27} \text{eV}$  → **ultra-light physics**

radiative stability:  $\tilde{\lambda} \lesssim 10^3 m^2 / M^2 \ll 1$  weak coupling:  $\lambda < 0.1$

► Central assumption:

$$\text{inv. duration: } \bar{\beta} = dS_E/dt \simeq \dot{\Gamma}/\Gamma \gg H$$

# NEDE Phenomenology

# NEDE Phenomenology

► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.

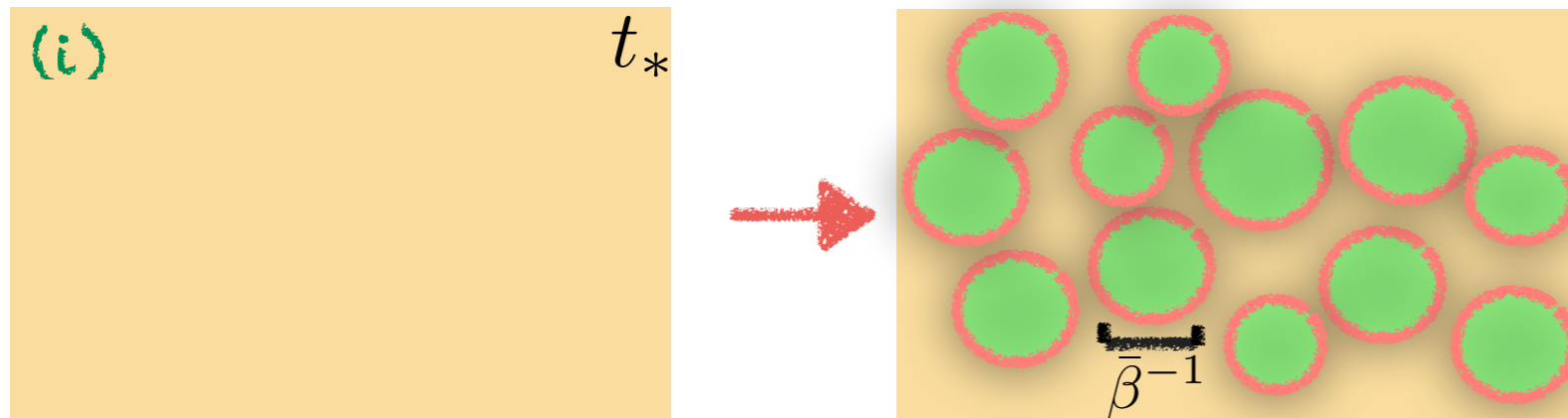
# NEDE Phenomenology

► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.



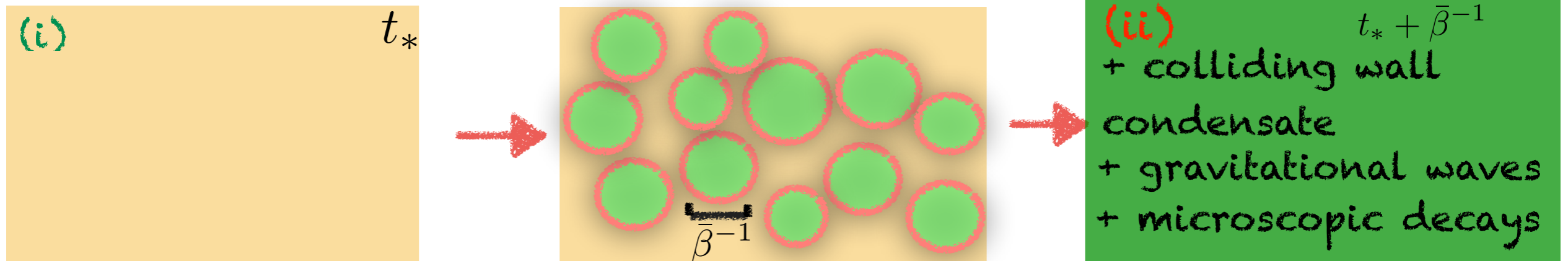
# NEDE Phenomenology

► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.



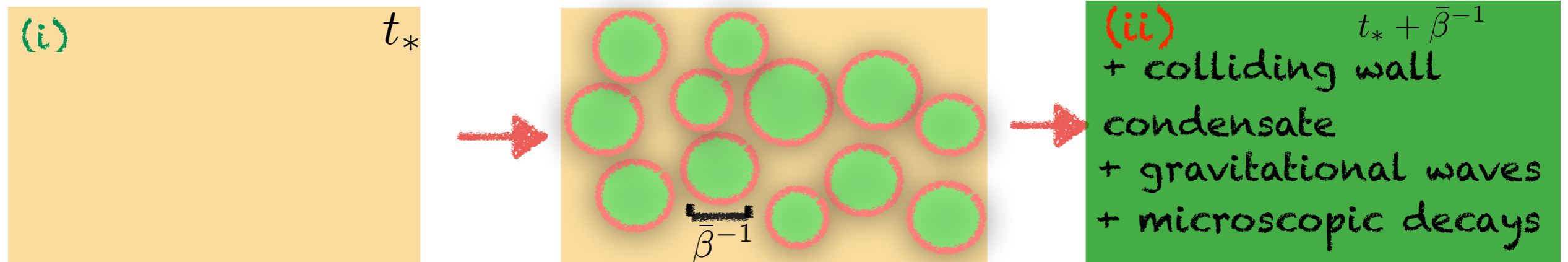
# NEDE Phenomenology

► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.



# NEDE Phenomenology

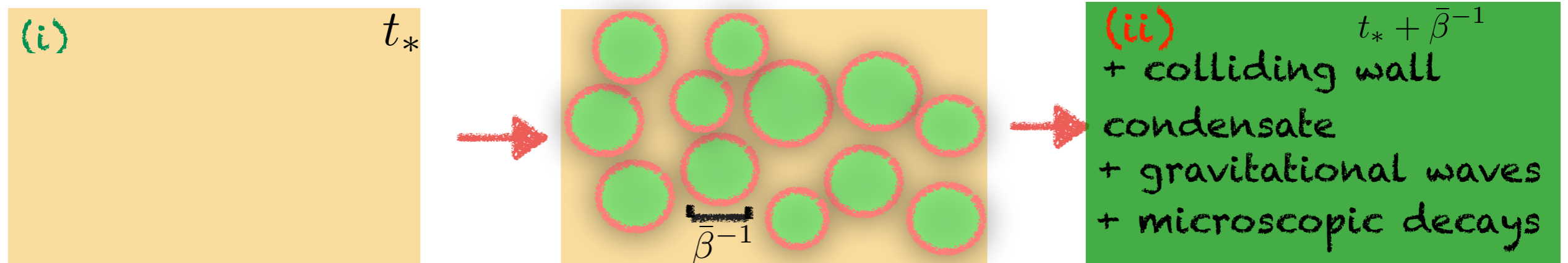
► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.



► **After PT:** Mixture of **radiation** and **small scale anisotropic stress**

# NEDE Phenomenology

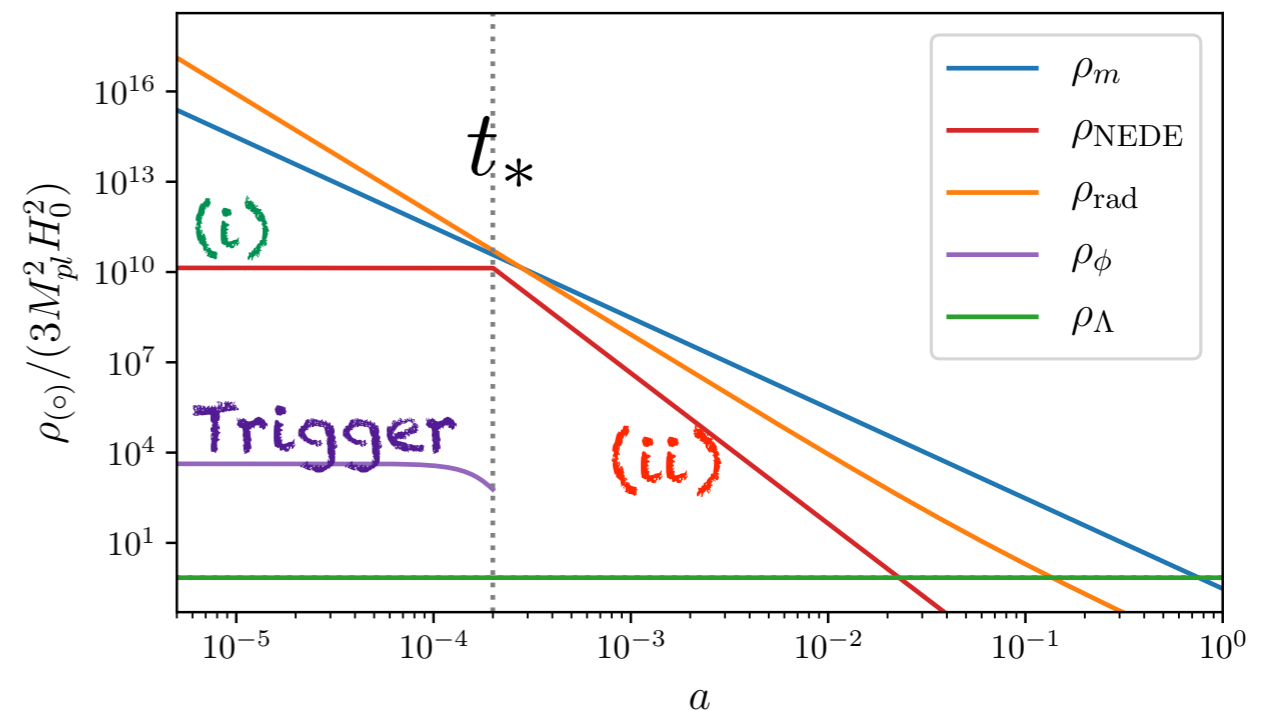
► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.



► **After PT:** Mixture of **radiation** and **small scale anisotropic stress**

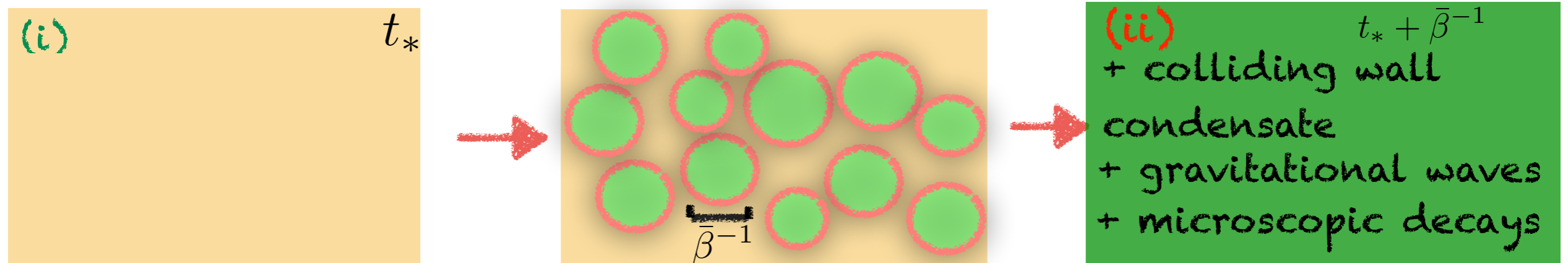
► NEDE as **phenomenological model:**

- (i) ● **Before** transition: NEDE plays role of CC.
- **Sudden** triggered transition at time:  $t_*$
- (ii) ● **After** transition: NEDE is described by decaying dark fluid with e.o.s.p.:



# NEDE Phenomenology

► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.

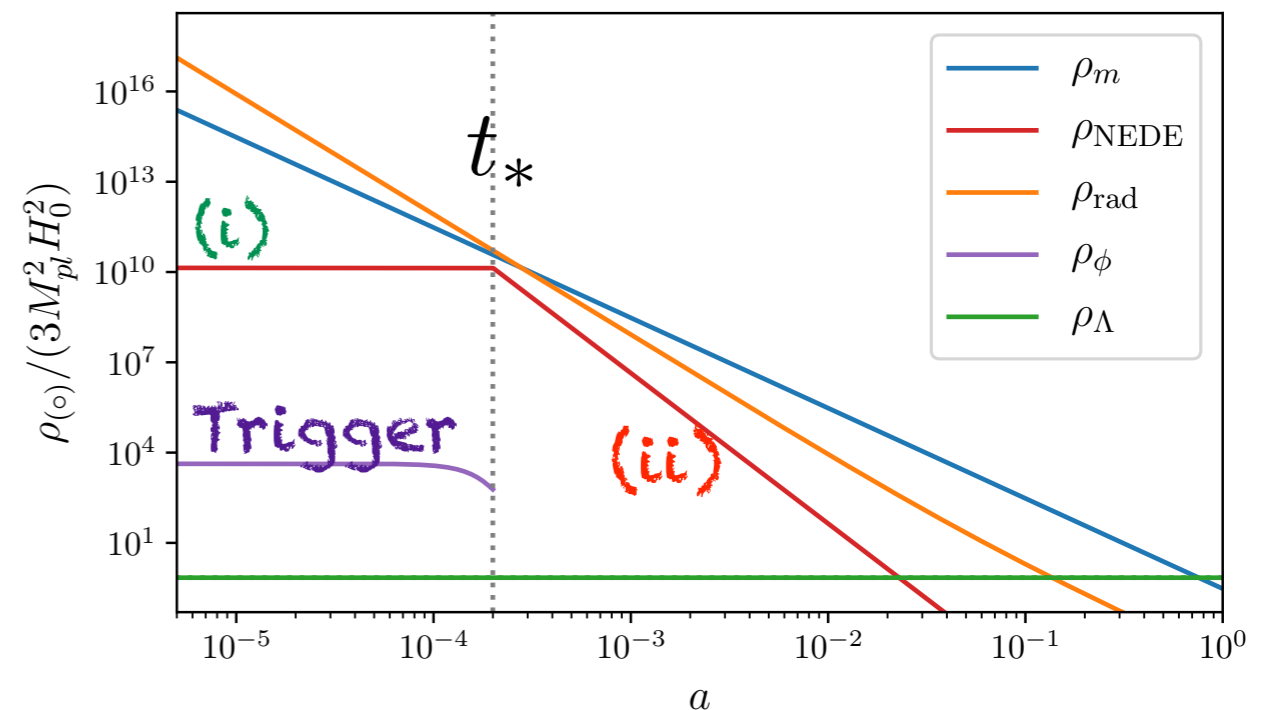


► **After PT:** Mixture of **radiation** and **small scale anisotropic stress**

► NEDE as **phenomenological model:**

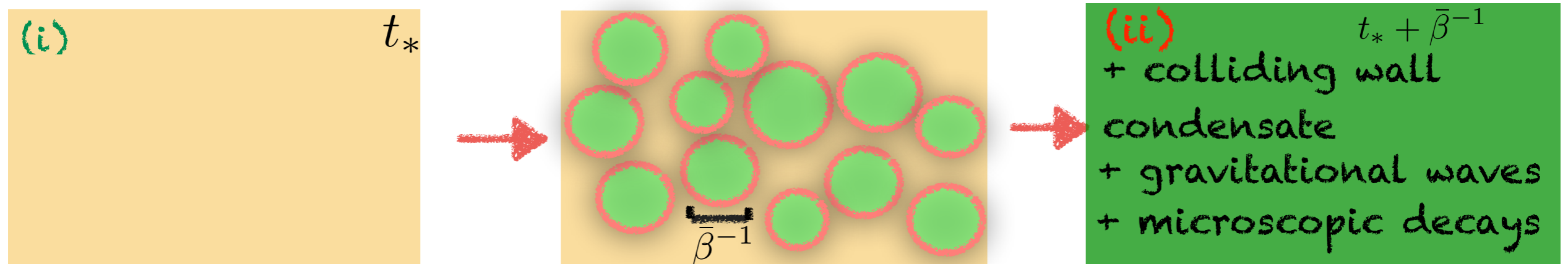
- (i) ● **Before** transition: NEDE plays role of CC.
- **Sudden** triggered transition at time:  $t_*$
- (ii) ● **After** transition: NEDE is described by decaying dark fluid with e.o.s.p.:

$$1/3 < w_{\text{NEDE}}(t) < 1$$



# NEDE Phenomenology

► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.



► **After PT:** Mixture of **radiation** and **small scale anisotropic stress**

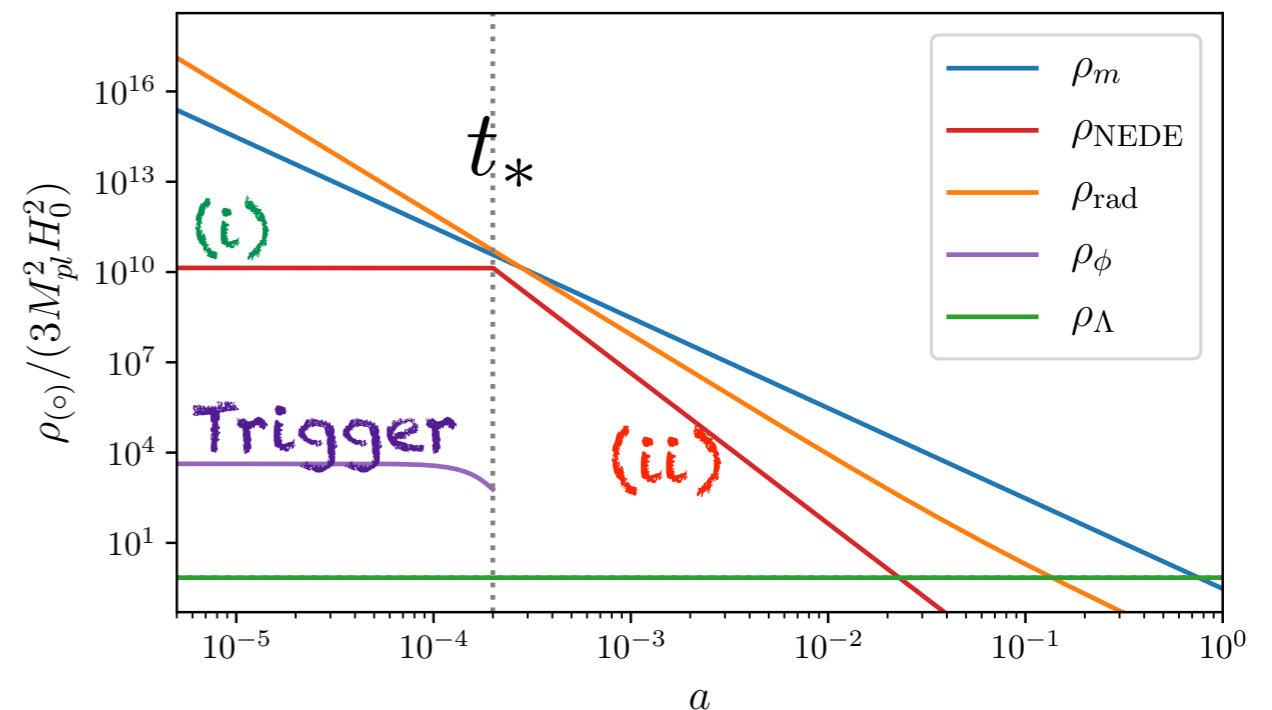
► NEDE as **phenomenological model:**

(i) ● **Before** transition: NEDE plays role of CC.

● **Sudden** triggered transition at time:  $t_*$

(ii) ● **After** transition: NEDE is described by decaying dark fluid with e.o.s.p.:

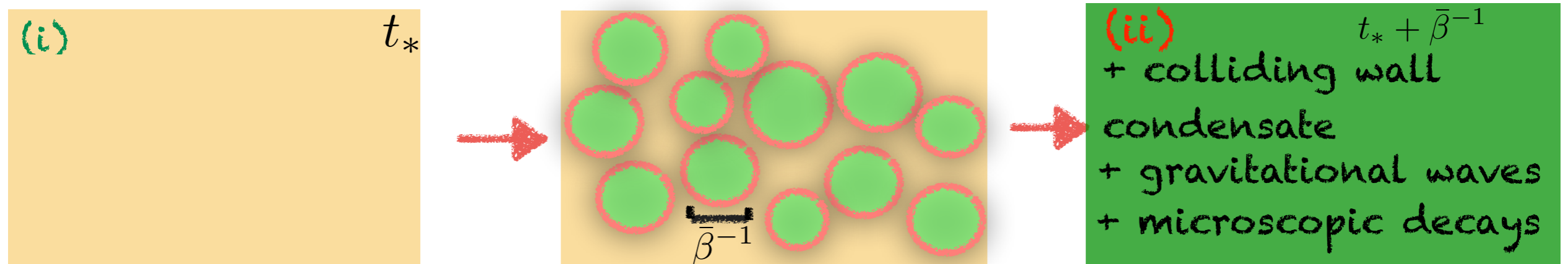
$$1/3 < w_{\text{NEDE}}(t) < 1$$



► Perturbations in decaying NEDE fluid seeded by trigger field perturbations.

# NEDE Phenomenology

► **Important result:** Phase transition is an **instantaneous** process on cosmological scales.



► **After PT:** Mixture of **radiation** and **small scale anisotropic stress**

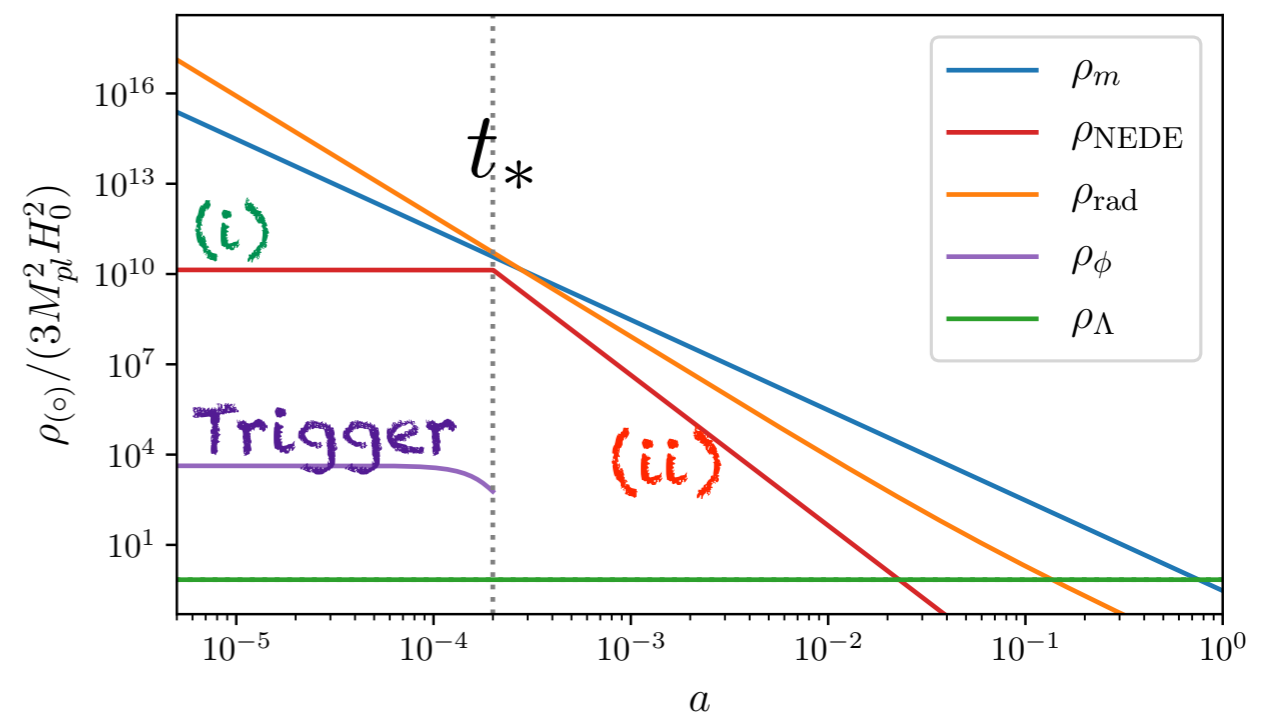
► NEDE as **phenomenological model:**

(i) ● **Before** transition: NEDE plays role of CC.

● **Sudden** triggered transition at time:  $t_*$

(ii) ● **After** transition: NEDE is described by decaying dark fluid with e.o.s.p.:

$$1/3 < w_{\text{NEDE}}(t) < 1$$



► Perturbations in decaying NEDE fluid seeded by trigger field perturbations.

► Description applies to other **triggered decay scenarios** too (e.g. hybrid NEDE).

# Cosmological parameter extraction

# Cosmological parameter extraction

► Consider **simplest** implementation of NEDE:

- ◆ parameters: trigger mass & fraction NEDE  $f_{\text{NEDE}}$
- ◆ fix  $w_{\text{NEDE}} = c_s^2 = 2/3$  (relaxed later)

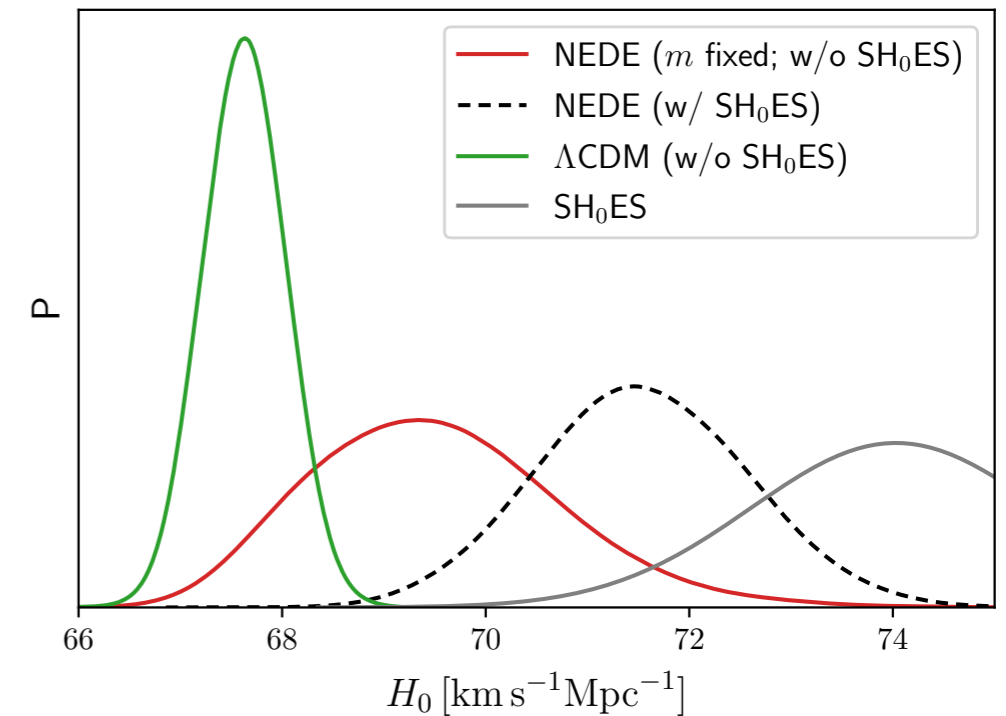
# Cosmological parameter extraction

► Consider **simplest** implementation of NEDE:

- ◆ parameters: trigger mass & fraction NEDE  $f_{\text{NEDE}}$
- ◆ fix  $w_{\text{NEDE}} = c_s^2 = 2/3$  (relaxed later)

Parameters	w/o SH <sub>0</sub> ES	w/ SH <sub>0</sub> ES
	$f_{\text{NEDE}}$ $m(\text{fixed})$	$f_{\text{NEDE}}$ $m$
$H_0$ [km/s/Mpc]	$69.6^{+1.0}_{-1.1}$	$71.4 \pm 1.0$
Evidence $f_{\text{NEDE}} \neq 0$	$\simeq 2\sigma$	$\simeq 4\sigma$
Hubble tension	$\simeq 2.5\sigma$	$(\simeq 1.5\sigma)$
S8 tension	$\simeq 2.7\sigma$	$\simeq 2.8\sigma$
$\Delta\chi^2$	-2.6	-15.6

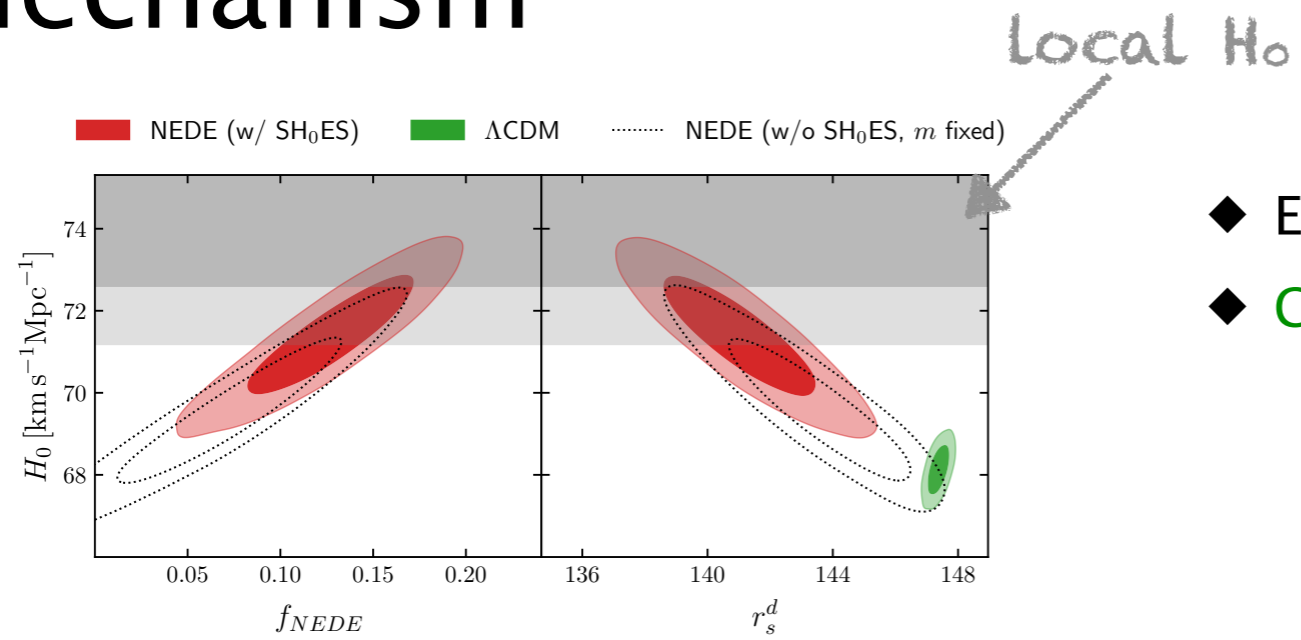
[Planck 2018, BAO, Pantheon, SH0ES 2019]



[FN&Sloth:2006.06686]

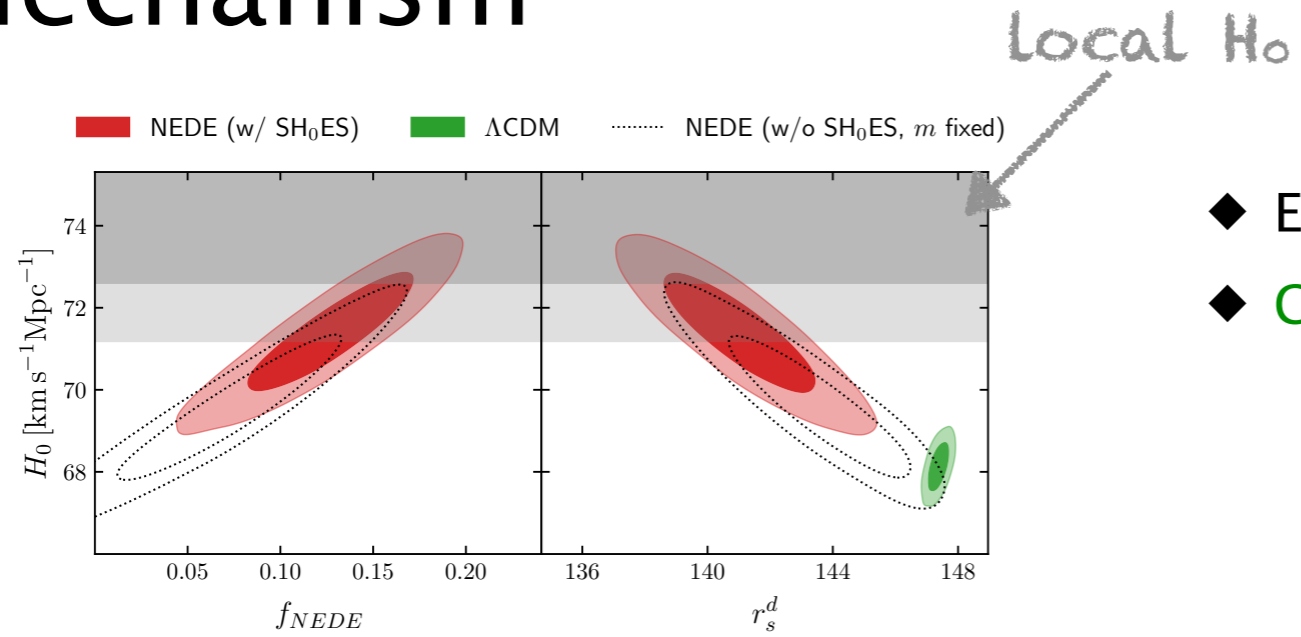
# Mechanism

# Mechanism

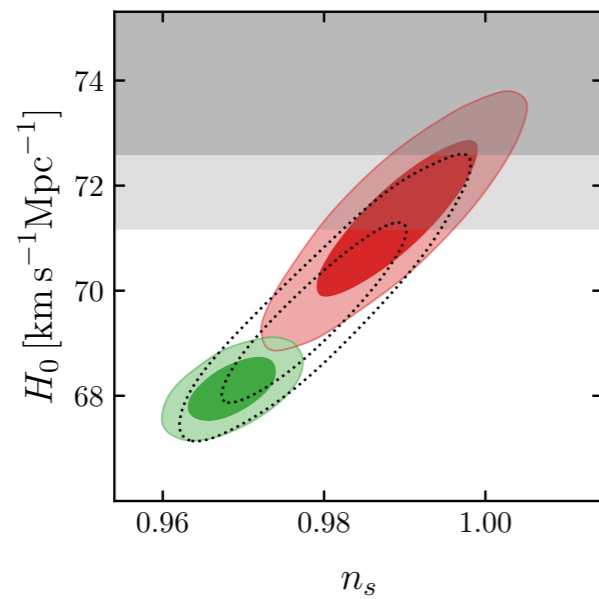


- ◆ Energy injection reduces sound horizon.
- ◆ **Compensated** by larger  $H_0$

# Mechanism

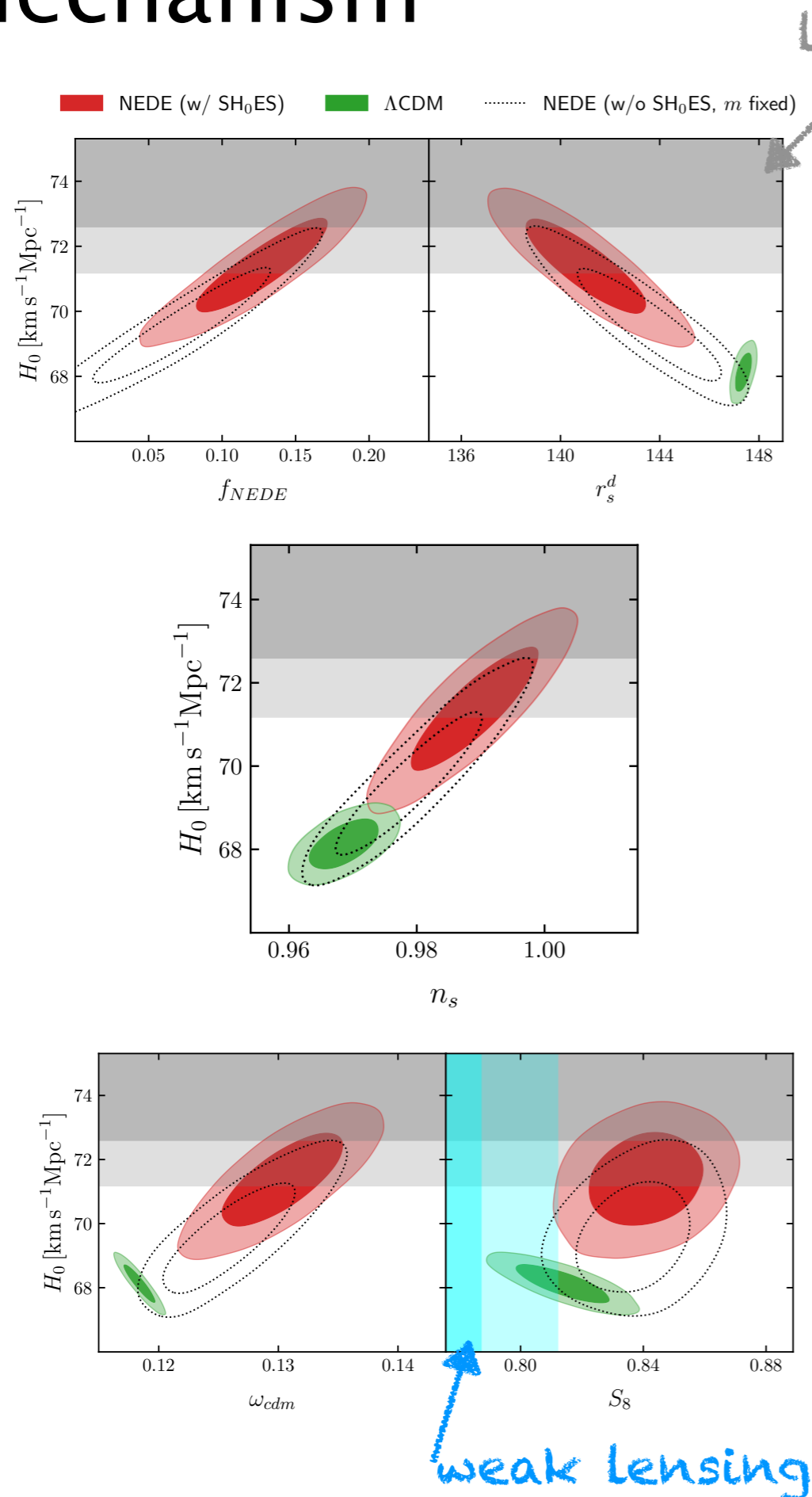


- ◆ Energy injection reduces sound horizon.
- ◆ **Compensated** by larger  $H_0$



- ◆ Enhanced diffusion damping.
- ◆ **Compensated** by larger  $n_s$

# Mechanism



- ◆ Energy injection reduces sound horizon.
- ◆ **Compensated** by larger  $H_0$

- ◆ Enhanced diffusion damping.
- ◆ **Compensated** by larger  $n_s$

- ◆ Quicker decay of Weyl potential due to NEDE perturbations and delayed matter domination.
- ◆ **Compensated** by large  $\omega_{\text{cdm}}$
- ◆ Small residual effect on small scales.
- ◆ Still:  $S_8$  tension comparable to  $\Lambda$ CDM (2.5  $\rightarrow$  2.8 sigma)



# The $H_0$ Olympics: A fair ranking of proposed models

Nils Schöneberg,<sup>a</sup> Guillermo Franco Abellán,<sup>b</sup> Andrea Pérez Sánchez,<sup>a</sup> Samuel J. Witte,<sup>c</sup> Vivian Poulin,<sup>b</sup> and Julien Lesgourgues<sup>a</sup>

deals with non-Gaussian posteriors

Model	$\Delta N_{\text{param}}$	$M_B$	Gaussian Tension	$Q_{\text{DMAP}}$ Tension		$\Delta\chi^2$	$\Delta\text{AIC}$		Finalist
$\Lambda\text{CDM}$	0	$-19.416 \pm 0.012$	$4.4\sigma$	$4.5\sigma$	$\times$	0.00	0.00	$\times$	$\times$
Majoron	3	$-19.380 \pm 0.027$	$3.0\sigma$	$2.9\sigma$	✓	-13.74	-7.74	✓	✓ ②
primordial B	1	$-19.390 \pm 0.018$	$3.5\sigma$	$3.5\sigma$	$\times$	-10.83	-8.83	✓	✓ ③
varying $m_e$	1	$-19.391 \pm 0.034$	$2.9\sigma$	$3.2\sigma$	$\times$	-9.87	-7.87	✓	✓ ③
varying $m_e + \Omega_k$	2	$-19.368 \pm 0.048$	$2.0\sigma$	$1.7\sigma$	✓	-16.11	-12.11	✓	✓ ①
EDE	3	$-19.390 \pm 0.016$	$3.6\sigma$	$1.6\sigma$	✓	-20.80	-14.80	✓	✓ ②
NEDE	3	$-19.380 \pm 0.021$	$3.2\sigma$	$2.0\sigma$	✓	-17.70	-11.70	✓	✓ ②

[Planck 2018 + BAO + Pantheon (+ SH0ES)]

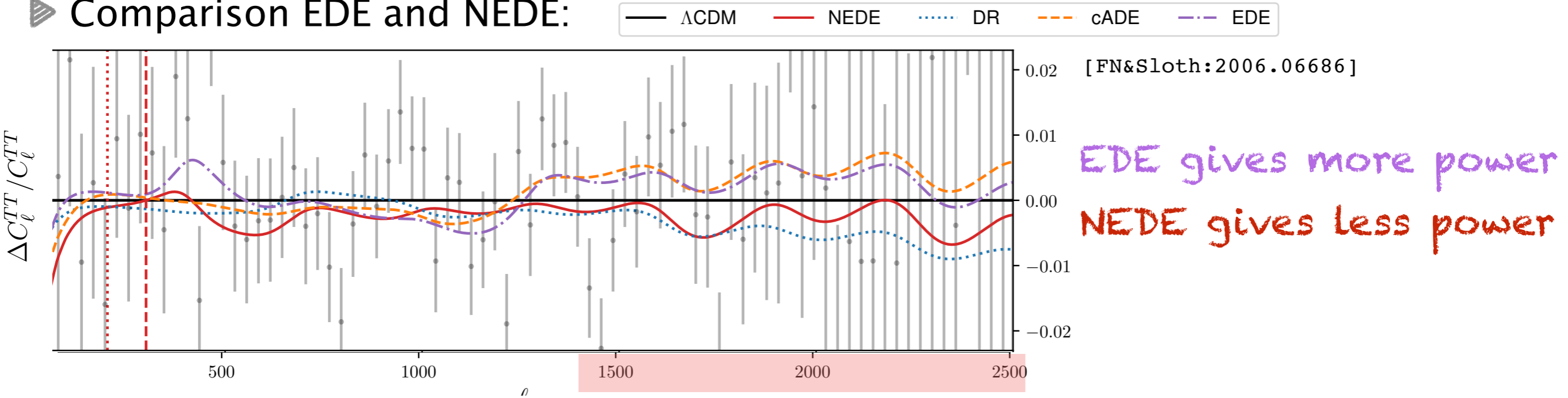
# The $H_0$ Olympics: A fair ranking of proposed models

Nils Schöneberg,<sup>a</sup> Guillermo Franco Abellán,<sup>b</sup> Andrea Pérez Sánchez,<sup>a</sup> Samuel J. Witte,<sup>c</sup> Vivian Poulin,<sup>b</sup> and Julien Lesgourgues<sup>a</sup>

Model	$\Delta N_{\text{param}}$	$M_B$	Gaussian Tension	$Q_{\text{DMAP}}$ Tension		$\Delta\chi^2$	$\Delta\text{AIC}$		Finalist
$\Lambda\text{CDM}$	0	$-19.416 \pm 0.012$	$4.4\sigma$	$4.5\sigma$	<i>X</i>	0.00	0.00	<i>X</i>	<i>X</i>
Majoron	3	$-19.380 \pm 0.027$	$3.0\sigma$	$2.9\sigma$	✓	-13.74	-7.74	✓	✓ ②
primordial B	1	$-19.390 \pm 0.018$	$3.5\sigma$	$3.5\sigma$	<i>X</i>	-10.83	-8.83	✓	✓ ③
varying $m_e$	1	$-19.391 \pm 0.034$	$2.9\sigma$	$3.2\sigma$	<i>X</i>	-9.87	-7.87	✓	✓ ③
varying $m_e + \Omega_k$	2	$-19.368 \pm 0.048$	$2.0\sigma$	$1.7\sigma$	✓	-16.11	-12.11	✓	✓ ①
EDE	3	$-19.390 \pm 0.016$	$3.6\sigma$	$1.6\sigma$	✓	-20.80	-14.80	✓	✓ ②
NEDE	3	$-19.380 \pm 0.021$	$3.2\sigma$	$2.0\sigma$	✓	-17.70	-11.70	✓	✓ ②

[Planck 2018 + BAO + Pantheon (+ SH0ES)]

## ► Comparison EDE and NEDE:



# The $H_0$ Olympics: A fair ranking of proposed models

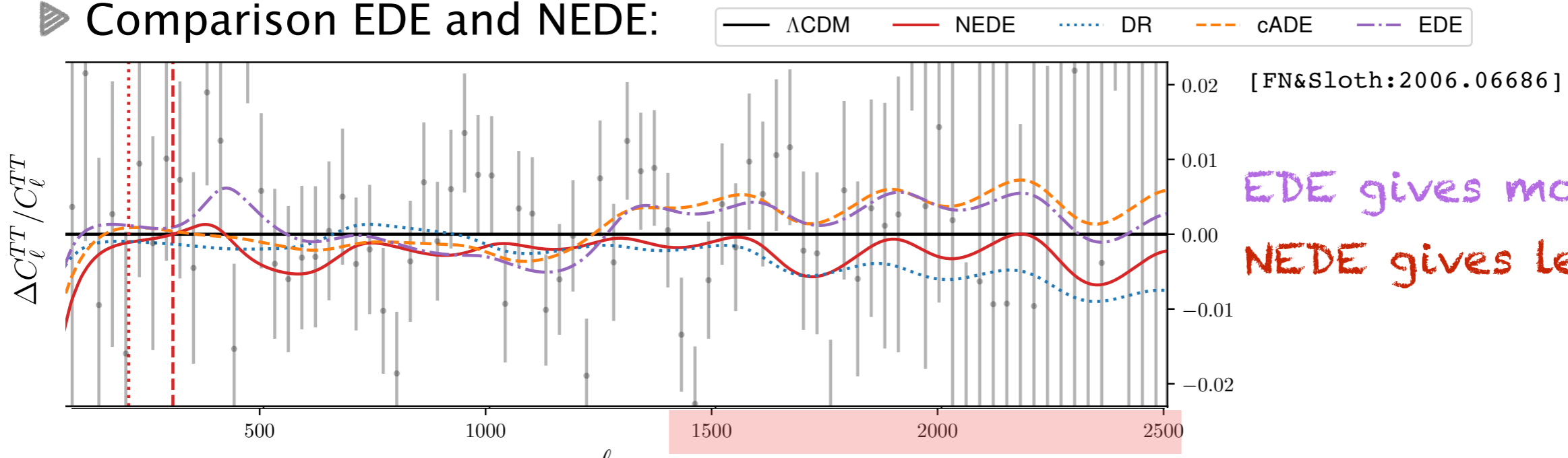
Nils Schöneberg,<sup>a</sup> Guillermo Franco Abellán,<sup>b</sup> Andrea Pérez Sánchez,<sup>a</sup> Samuel J. Witte,<sup>c</sup> Vivian Poulin,<sup>b</sup> and Julien Lesgourgues<sup>a</sup>

deals with non-Gaussian posteriors

Model	$\Delta N_{\text{param}}$	$M_B$	Gaussian Tension	$Q_{\text{DMAP}}$ Tension		$\Delta\chi^2$	$\Delta\text{AIC}$		Finalist
$\Lambda\text{CDM}$	0	$-19.416 \pm 0.012$	$4.4\sigma$	$4.5\sigma$	$\times$	0.00	0.00	$\times$	$\times$
Majoron	3	$-19.380 \pm 0.027$	$3.0\sigma$	$2.9\sigma$	✓	-13.74	-7.74	✓	✓ ②
primordial B	1	$-19.390 \pm 0.018$	$3.5\sigma$	$3.5\sigma$	$\times$	-10.83	-8.83	✓	✓ ③
varying $m_e$	1	$-19.391 \pm 0.034$	$2.9\sigma$	$3.2\sigma$	$\times$	-9.87	-7.87	✓	✓ ③
varying $m_e + \Omega_k$	2	$-19.368 \pm 0.048$	$2.0\sigma$	$1.7\sigma$	✓	-16.11	-12.11	✓	✓ ①
EDE	3	$-19.390 \pm 0.016$	$3.6\sigma$	$1.6\sigma$	✓	-20.80	-14.80	✓	✓ ②
NEDE	3	$-19.380 \pm 0.021$	$3.2\sigma$	$2.0\sigma$	✓	-17.70	-11.70	✓	✓ ②

[Planck 2018 + BAO + Pantheon (+ SH0ES)]

## ► Comparison EDE and NEDE:



EDE gives more power  
NEDE gives less power

- More precise high- $\ell$  TT (and EE) data can help distinguish EDE and NEDE.
- A special role will be played by LSS data.
- Key message: perturbation sector matters

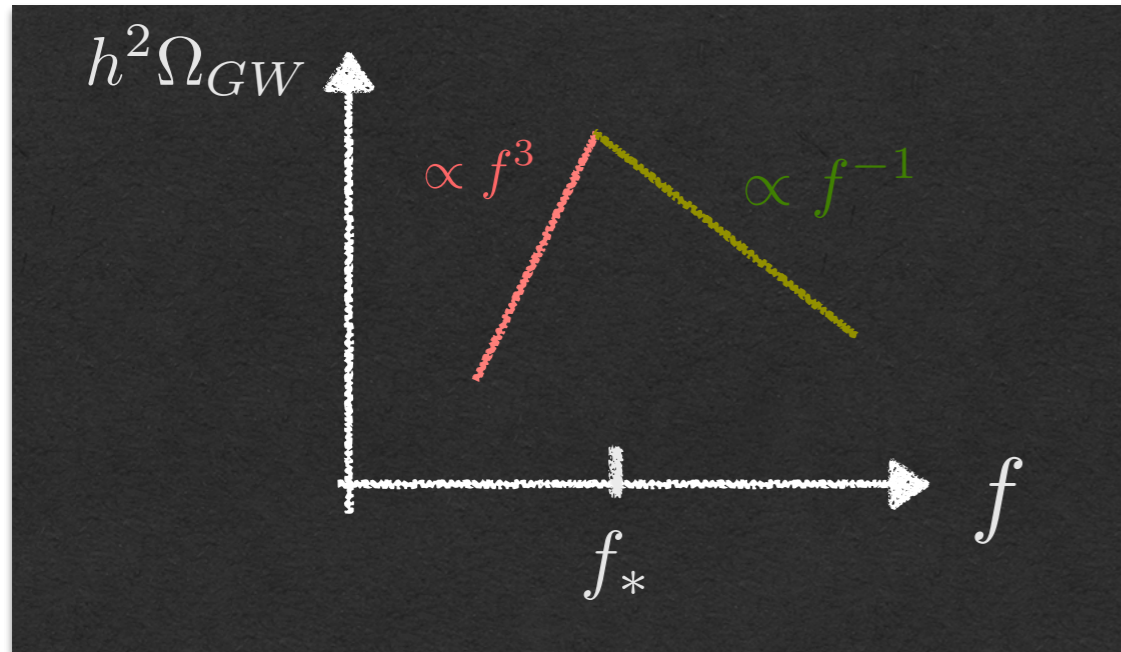
# Gravitational Waves

# Gravitational Waves

- ▶ First order phase transitions (PT) act as source of gravitational waves.

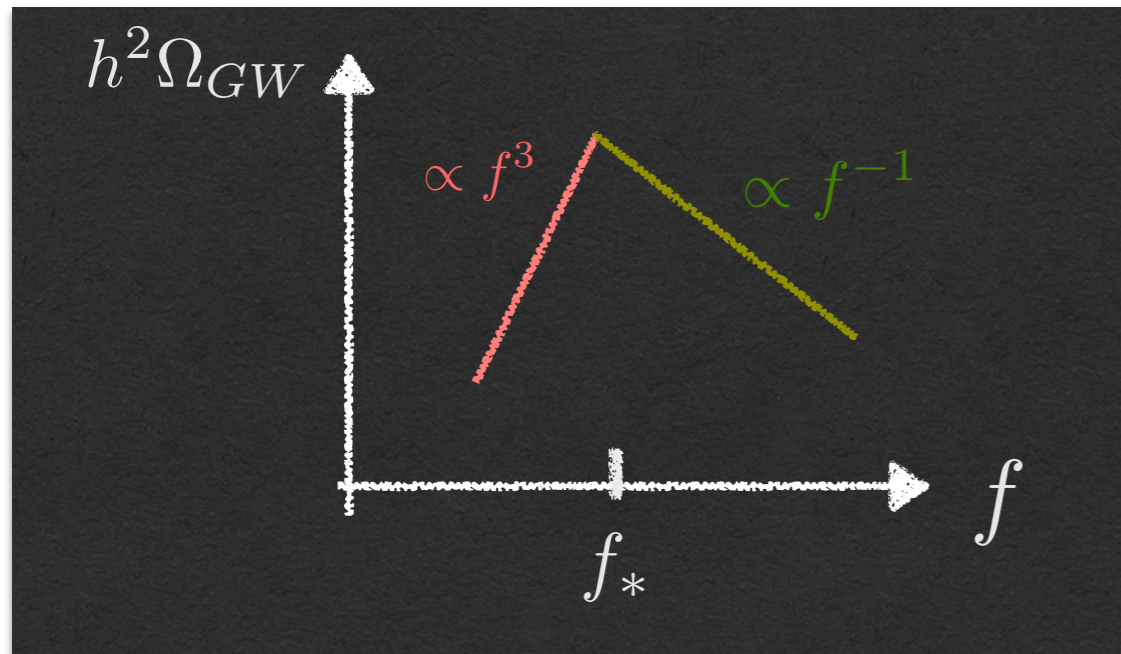
# Gravitational Waves

- First order phase transitions (PT) act as source of gravitational waves.



# Gravitational Waves

- First order phase transitions (PT) act as source of gravitational waves.



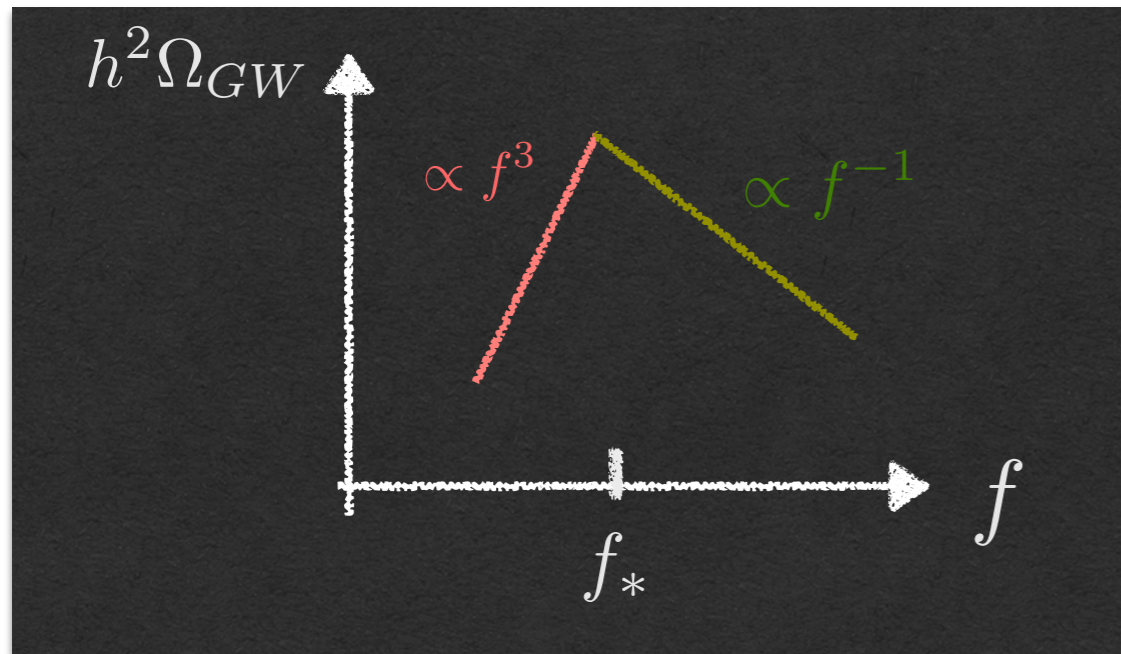
1/f regime:

$$h^2 \Omega_{GW} \sim 10^{-15} H \bar{\beta}^{-1} \left( \frac{10^{-9} \text{Hz}}{f} \right)$$

single dial

# Gravitational Waves

- First order phase transitions (PT) act as source of gravitational waves.



1/f regime:

$$h^2 \Omega_{GW} \sim 10^{-15} H \bar{\beta}^{-1} \left( \frac{10^{-9} \text{Hz}}{f} \right)$$

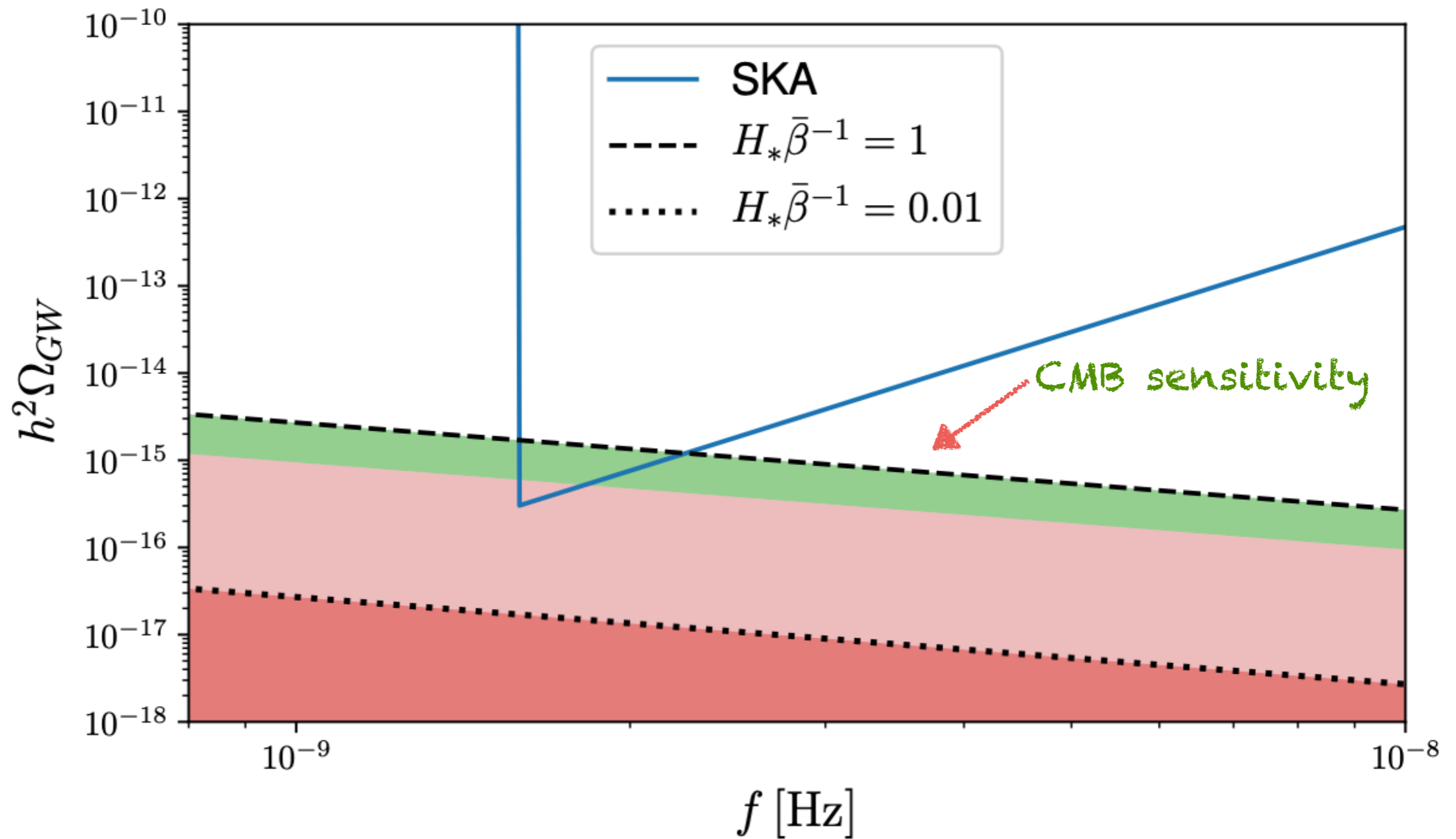
single dial

- Moderate prospects of detection with **pulsar timing arrays**.

Square Kilometer Array, sensitivity:  $h^2 \Omega_{GW} \sim 10^{-15}$

→ window for detection:  $0.1 < H \bar{\beta}^{-1} \lesssim 1$

- First order phase transitions (PT) act as source of gravitational waves.



- Marginally compatible with Square Kilometre Array

# Summary

- ▶ H0 and S8 tension exciting opportunity to probe the dark sector.
- ▶ EDE looks promising, although the potential appears fine-tuned  
+ no solution to S8 tension → look for new particle physics models!
- ▶ NEDE brings H0 tension down to 2.5 sigma.
- ▶ **Unique** signature: gravitational waves.
- ▶ Three phenomenological challenges remain:
  - 1. **S8 tension:** → Dark sector interactions
  - 2. New **coincidence problem** → Martin's talk
  - 3. Unknown systematics → Stay open about new developments.
- ▶ Further theoretical work
  - Find new microscopic scenarios for NEDE (see Martin's talk).
  - Relate NEDE fluid to microscopic parameters.
  - Multi-axion system: small masses protected by approximately broken shift symmetry.

