

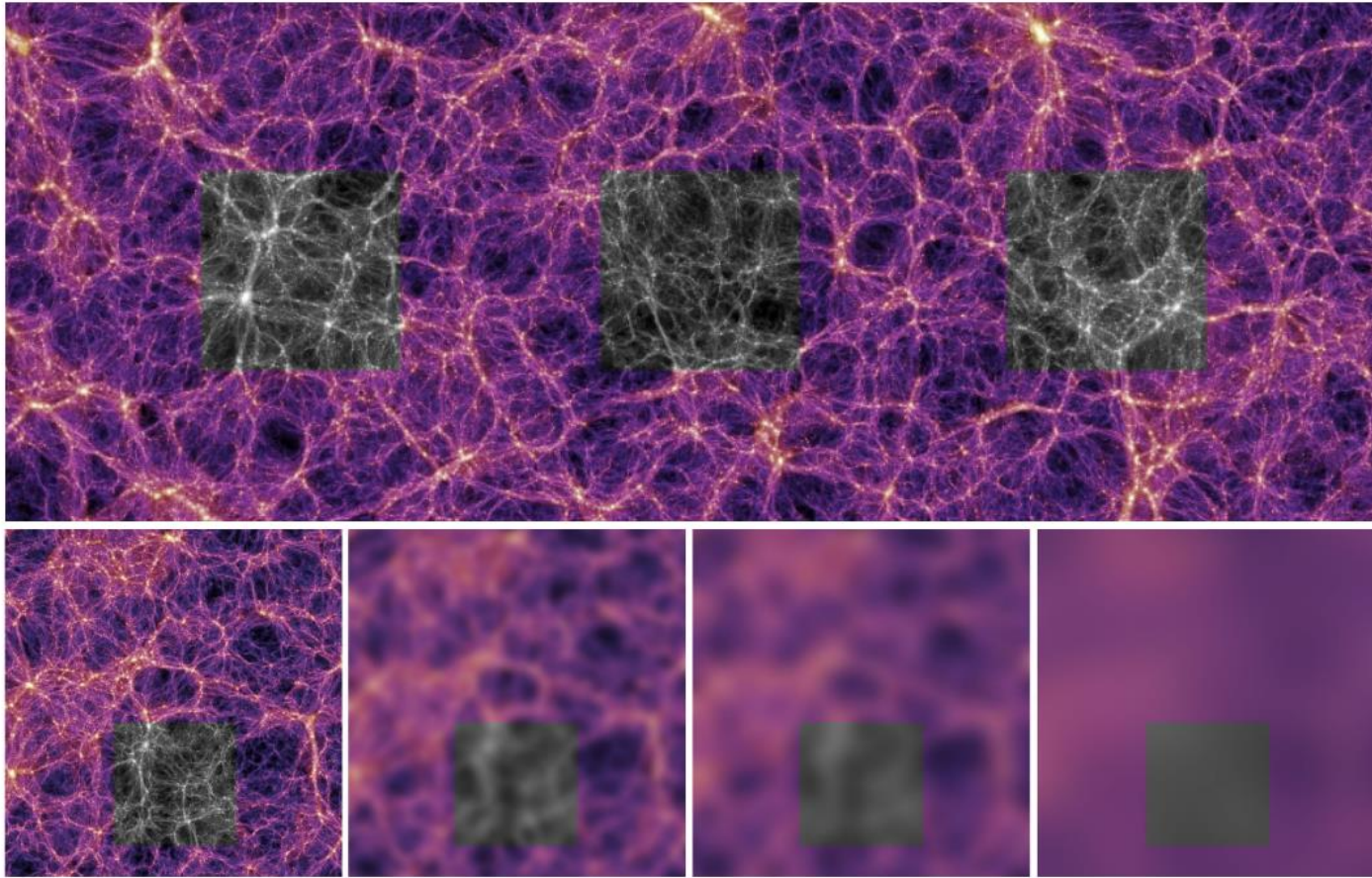
Relativistic structure formation: extreme objects

Jan J. Ostrowski, NCBJ, Warsaw

Hot topics in Modern Cosmology

Spontaneous Workshop XIV

May 8th – 14th, 2022, IESC Cargese, France



Clarkson, Rept. Prog. Phys. 74 (2011)

Averaging Coarse-graining of structure, such that small-scale effects are hidden to reveal large scale geometry and dynamics.

Backreaction Gravity gravitates, so local gravitational inhomogeneities may affect the cosmological dynamics. How this is calculated depends on the degree of coarse graining.

Fitting How do we appropriately fit an idealized model to observations made from one location in a lumpy universe, given that this ‘background’ does not in fact exist?

Averaging/smoothing does not commute with evaluating the inverse metric,
connections and curvature

$$\begin{array}{ccccccc}
 g_{ab}^n & \longrightarrow & \langle \rangle & \longrightarrow & g_{ab}^{n-1} & \longrightarrow & \langle \rangle & \longrightarrow & g_{ab}^1 & \longrightarrow & \langle \rangle & \longrightarrow & g_{ab}^0 \\
 \downarrow & & & & \downarrow & & & & \downarrow & & & & \downarrow \\
 G_{ab}^n = \kappa T_{ab} & \longrightarrow & \langle \rangle & \longrightarrow & \langle G_{ab}(g_{ab}^n) \rangle \neq G_{ab}(g_{ab}^{n-1}) & \longrightarrow & \langle \rangle & \longrightarrow & \dots & \longrightarrow & \langle \rangle & \longrightarrow & \langle G_{ab}(g_{ab}^n) \rangle \neq G_{ab}(g_{ab}^0)
 \end{array}$$

- There is no overall agreement on how to construct the cosmological background
- We can test different backgrounds with the large scale structure formation
- We can try to say something about the large scale structures without or with a minimal reference to the background

The silent universe approach

- Initial conditions compatible with the concordance model
- Non-linear evolution of density in 1+3 formalism; silent universe equations : no rotation, no gravitational waves, no heat flux
- Numerical evaluation of the mass function of cosmological structures
- Poisson statistics of the most massive objects in the Universe

Silent universe equations:

$$\dot{\rho} = -\rho \Theta$$

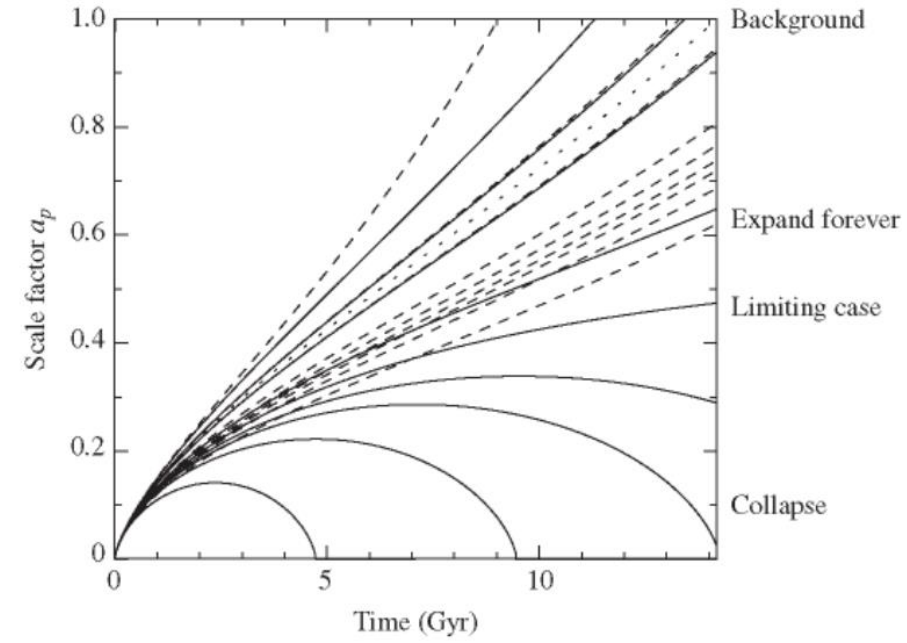
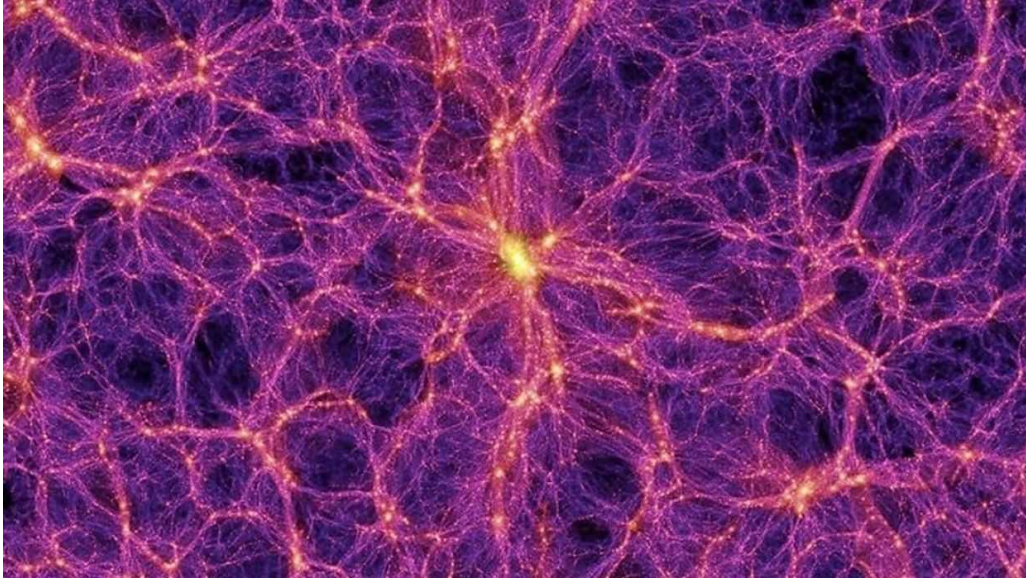
$$\dot{\Theta} = -\frac{1}{3}\Theta^2 - \frac{1}{2}\kappa\rho - 6\sigma^2 + \Lambda$$

$$\dot{\sigma} = -\frac{2}{3}\Theta\sigma + \sigma^2 - \mathcal{W}$$

$$\dot{\mathcal{W}} = -\Theta\mathcal{W} - \frac{1}{2}\kappa\rho\sigma - 3\sigma\mathcal{W}$$

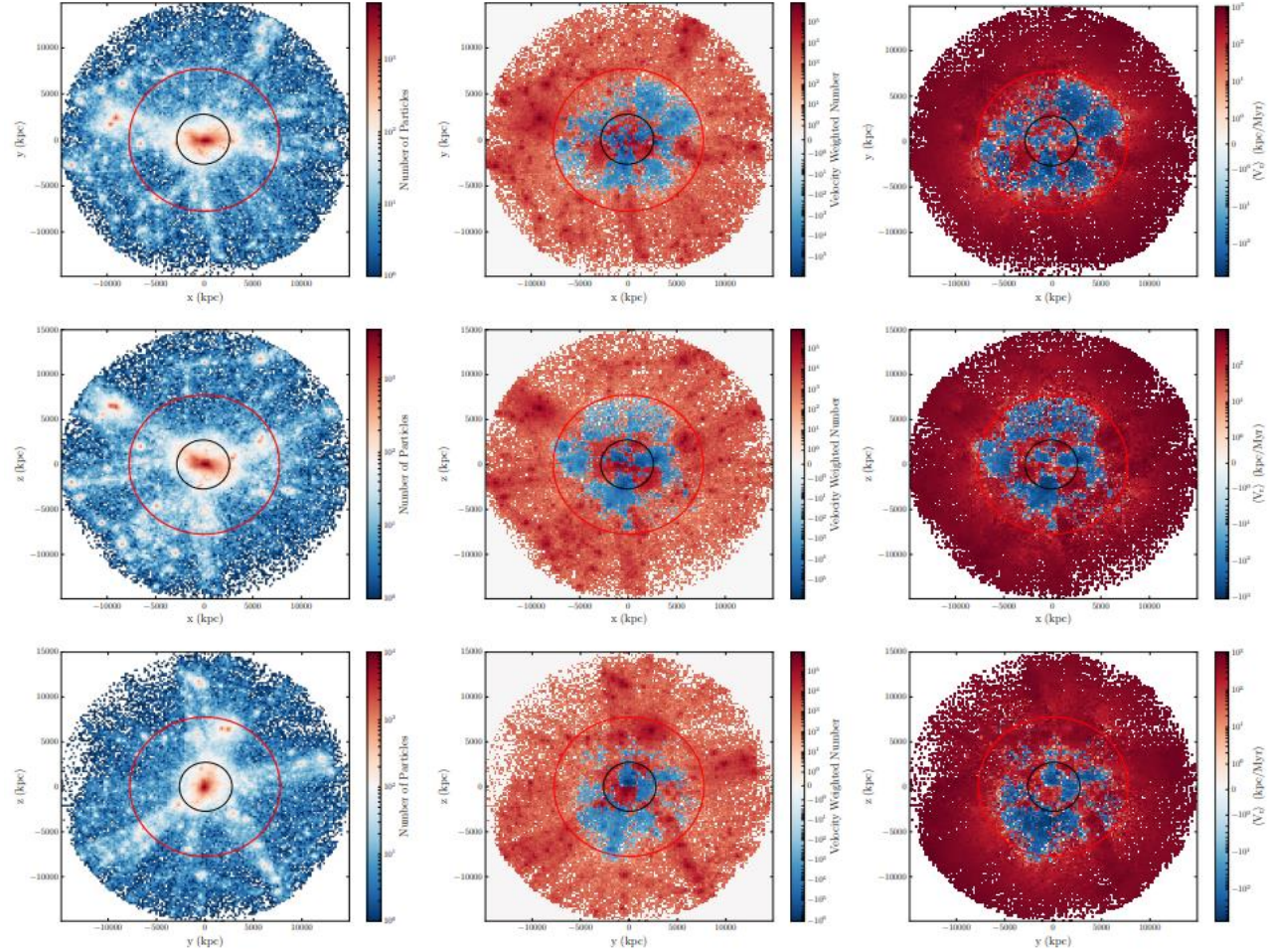
Mass function of galaxy clusters:

$$\frac{\frac{dN}{d\ln M}}{M} = \frac{\varrho_0}{M} f(\sigma_M) \left| \frac{d\ln \sigma_M}{d\ln M} \right|$$



Millenium project, Press & Schechter 1974

G. Korkidis et al.: Turnaround radius of galaxy clusters in N-body simulations



Maximum radius of structure: $R_{max} = \left(\frac{3GM}{\Lambda} \right)^{1/3}$

Spherically symmetric metric

$$ds^2 = -dt^2 + \frac{(R')^2}{1+2E}dr^2 + R^2d\Omega^2$$

$$\dot{R}^2 - \frac{2GM}{R} = 2E \qquad 4\pi G\rho = \frac{M'}{R^2R'}$$

$$E < 0 \quad , \quad R = -\frac{M}{2E}(1 - \cos\mu) \quad , \quad \mu - \sin\mu = \frac{(-2E)^{3/2}}{M}(t - t_B)$$

$$E = 0 \quad , \quad R = \left(\frac{9}{2}M(t - t_B)^2\right)^{1/3}$$

$$E > 0 \quad , \quad R = \frac{M}{2E}(\cosh\mu - 1) \quad , \quad \sinh\mu - \mu = \frac{(2E)^{3/2}}{M}(t - t_B) \quad ,$$

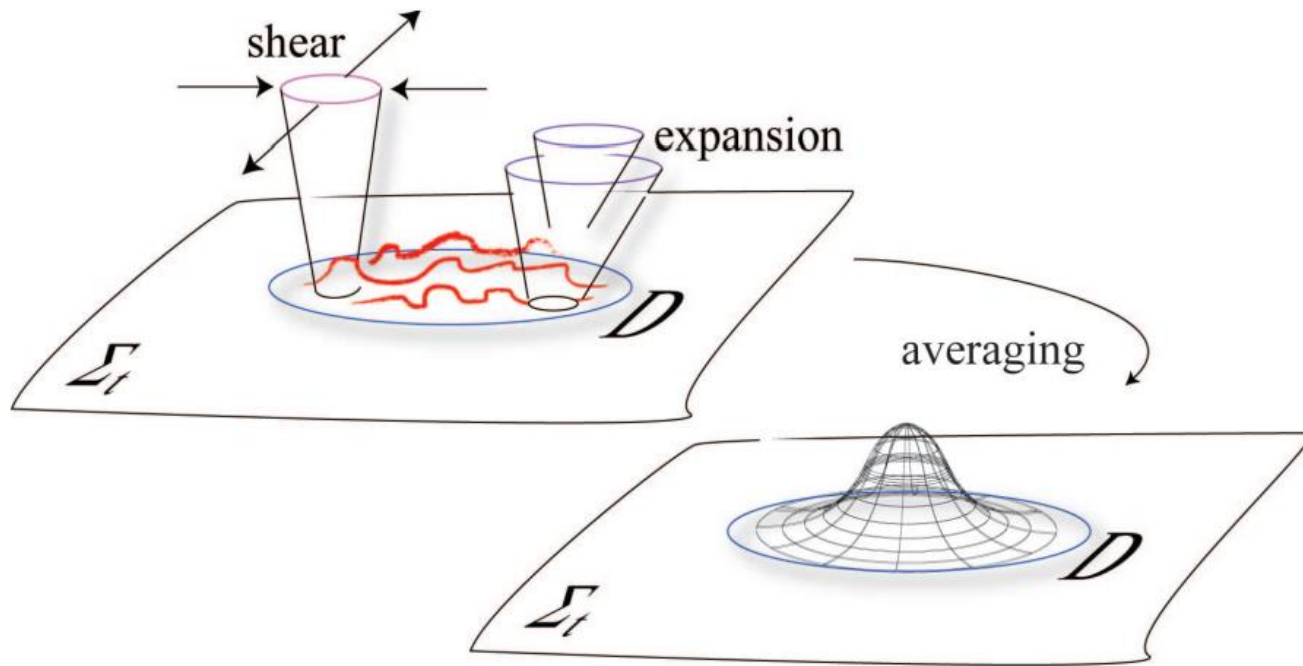
Hamiltonian constrain in the co-moving, synchronous gauge:

$$\frac{1}{3}\Theta^2 = 8\pi G\rho + \sigma^2 - \frac{1}{2}\mathcal{R} + \Lambda.$$

Maximum volume condition is equivalent to the zero averaged

$$\begin{aligned} 3H_{\mathcal{D}} := \langle \theta \rangle_{\mathcal{D}} &= \frac{4\pi}{3V} \int_0^{r_{\mathcal{D}}} \frac{\partial_r (\dot{R}R^2)}{\sqrt{1+2E}} dr \\ V &= \frac{4\pi}{3} \int_0^{r_{\mathcal{D}}} \frac{\partial_r (R^3)}{\sqrt{1+2E}} dr \end{aligned}$$

$$3H_{\mathcal{D}} = \frac{\partial_t V}{V}$$



Buchert and Carfora, CQG (2008)

$$a_{\mathcal{D}} = \left(\frac{V_{\mathcal{D}}(t)}{V_{\mathcal{D}_i}} \right)^{\frac{1}{3}} \qquad H_{\mathcal{D}} := \frac{\dot{a}_{\mathcal{D}}}{a_{\mathcal{D}}}$$

$$3 \frac{\ddot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} + 4\pi G \langle \rho \rangle_{\mathcal{D}} = \mathcal{Q}_{\mathcal{D}} \quad , \quad H_{\mathcal{D}}^2 - \frac{8\pi G}{3} \langle \rho \rangle_{\mathcal{D}} = - \frac{\langle \mathcal{R} \rangle_{\mathcal{D}} + \mathcal{Q}_{\mathcal{D}}}{6}$$

$$\mathcal{Q}_{\mathcal{D}} = 2 \langle \text{II} \rangle_{\mathcal{D}} - \frac{2}{3} \langle \text{I} \rangle_{\mathcal{D}}^2 \qquad \text{I} = \text{tr} \left(\Theta_{ij} \right) \quad , \quad \text{II} = \frac{1}{2} \left((\text{tr} \Theta_{ij})^2 - \text{tr} \left((\Theta_{ij})^2 \right) \right) \quad , \quad \text{III} = \det \left(\Theta_{ij} \right)$$

Relativistic Zel'dovich approximation

$${}^{\text{RZA}}Q_{\mathcal{D}} = \frac{\dot{\xi}^2 (\gamma_1 + \xi \gamma_2 + \xi^2 \gamma_3)}{(1 + \xi \langle \text{I}_{\mathbf{i}} \rangle_{\mathcal{I}} + \xi^2 \langle \text{II}_{\mathbf{i}} \rangle_{\mathcal{I}} + \xi^3 \langle \text{III}_{\mathbf{i}} \rangle_{\mathcal{I}})^2}, \quad \text{with :}$$

$$\begin{cases} \gamma_1 := 2 \langle \text{II}_{\mathbf{i}} \rangle_{\mathcal{I}} - \frac{2}{3} \langle \text{I}_{\mathbf{i}} \rangle_{\mathcal{I}}^2 \\ \gamma_2 := 6 \langle \text{III}_{\mathbf{i}} \rangle_{\mathcal{I}} - \frac{2}{3} \langle \text{II}_{\mathbf{i}} \rangle_{\mathcal{I}} \langle \text{I}_{\mathbf{i}} \rangle_{\mathcal{I}} \\ \gamma_3 := 2 \langle \text{I}_{\mathbf{i}} \rangle_{\mathcal{I}} \langle \text{III}_{\mathbf{i}} \rangle_{\mathcal{I}} - \frac{2}{3} \langle \text{II}_{\mathbf{i}} \rangle_{\mathcal{I}}^2 \end{cases}$$

$$\xi(t) := [q(t) - q(t_{\mathbf{i}})]/\dot{q}(t_{\mathbf{i}})$$

$$\ddot{q}(t) + 2H\dot{q}(t) - 4\pi\rho_H q(t) = 0$$

$$\langle \mathcal{A} \rangle_{\mathcal{D}} = \frac{\langle \mathcal{A}J \rangle_{\mathcal{I}}}{\langle J \rangle_{\mathcal{I}}}, \quad \langle J \rangle_{\mathcal{I}} = a_{\mathcal{D}}^3$$

$$\mathcal{J}(t, \mathbf{X}) := 1 + \xi(t)\text{I}_{\mathbf{i}} + \xi^2(t)\text{II}_{\mathbf{i}} + \xi^3(t)\text{III}_{\mathbf{i}}$$

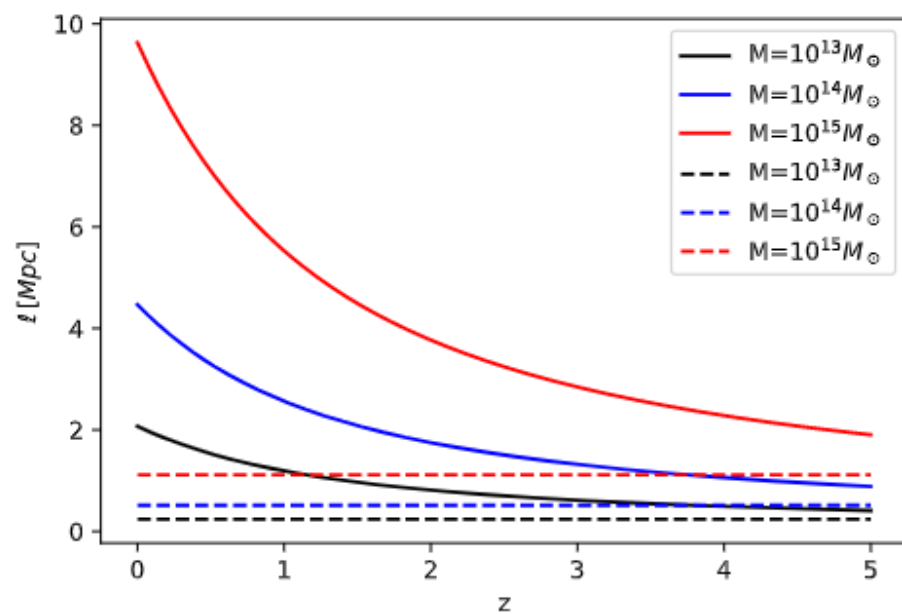
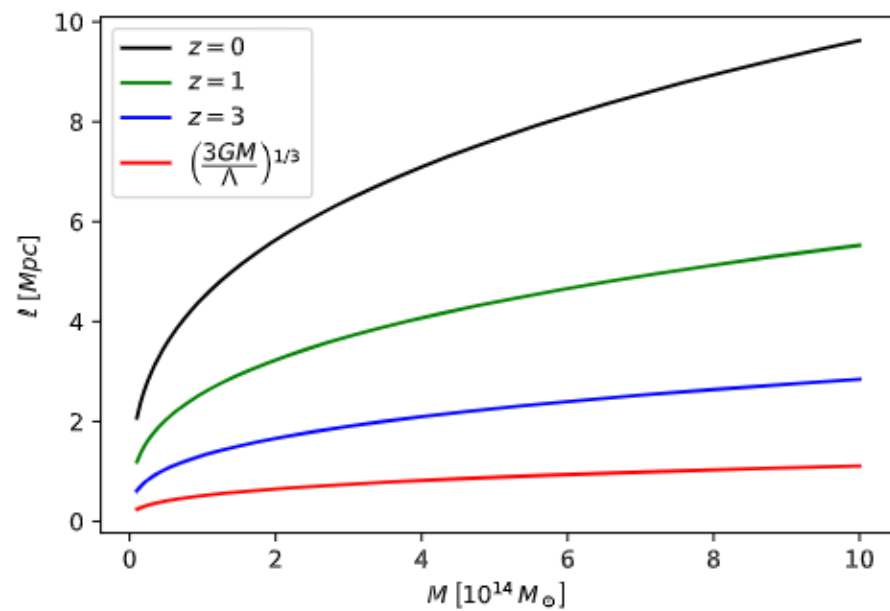
$$J \approx {}^{\text{RZA}}J := a^3 \mathcal{J}$$

Using the condition for the turnaround:

$$\text{Turnaround:} \quad 3H_{\mathcal{D}} = \langle \Theta \rangle_{\mathcal{D}} = \frac{\dot{V}_{\mathcal{D}}}{V_{\mathcal{D}}} = 0 \implies \dot{V}_{\mathcal{D}} = 0$$

averaged equations and the Relativistic Zel'dovich approximation as a closure condition, we get:

$$V_{\mathcal{D}}(t_{ta}) = \frac{M}{\rho_H \left(1 + 3H \left(\frac{\dot{q}}{q} \right)^{-1} \right)}$$



Conclusions

- Non-linear, relativistic approach to the cosmological structure formation can guide us towards solving the averaging problem
- Extreme cosmological structures provide a useful test for different cosmological backgrounds and/or alternative theories of gravity
- Further details can be found in: *Bolejko & Ostrowski PRD (2020)*, *Roukema & Ostrowski JCAP (2020)*, *Ostrowski AppB (2020)*, *Ostrowski & Delgado Gaspar JCAP (2022)*