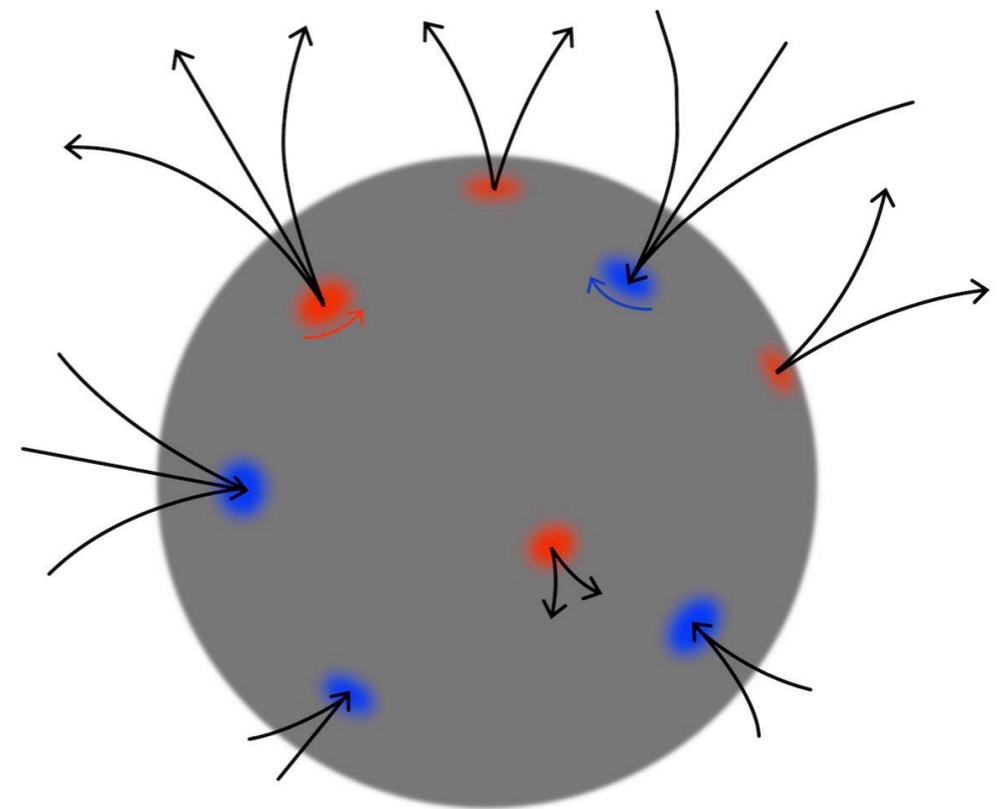


Vortexes and Black Holes

Michael Zantedeschi - LMU MPI
Spontaneous workshop XIV
13th May 2022



G. Dvali, F. Kühnel and MZ arXiv:2112.08354

Main Message:

Properties of black holes are not unique to gravity

There are localized and self-sustained configurations outside of gravity that possess properties analogous to black holes:

1. Entropy-area law $S_{\text{BH}} = \text{Area}/G_{\text{Gold}}$
This is maximal entropy compatible with unitarity
Objects saturating this bound will be called saturons

G.Dvali arXiv:2003.05546

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2. Hawking thermal emission $T = 1/R$
3. Long time information retrieval (Page time) $t \sim S_{\text{BH}} R \sim R^3/G_{\text{Gold}}$
4. Maximal spin and halt of Hawking emission $J \lesssim S_{\text{BH}}$ *G. Dvali, F. Kühnel and MZ*
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arXiv:2112.08354
- 2, 3 and 4 are a consequence of 1
 - Black holes are saturons under identification $G_{\text{Gold}} = G_{\text{N}}$

Saturn configurations can be found also not in gravity,
e.g.

- Bose-Einstein Condensates
- Gross-Neveu model (Otari Shakelashvili in previous talk)
- Color-glass condensate
- Renormalizable QFTs in 3+1 dim. (example shown in this talk)

Goal: provide a deeper understanding of black holes studying
systems easier to control

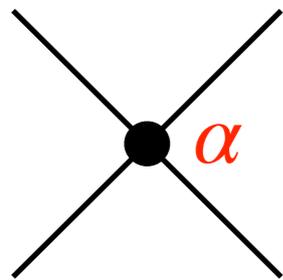
Possible to come up with new predictions?

An example of Saturon

Consider scalar field Φ_i^j in the adjoint representation of $SU(N)$ theory

$$i, j = 1, \dots, N \rightarrow \text{"flavour" indices}, \quad \text{Tr } \Phi = 0, \quad \Phi = \Phi^\dagger$$

$$\mathcal{L} = \text{Tr } \partial_\mu \Phi \partial^\mu \Phi - \alpha \text{Tr} \left(f \Phi - \Phi^2 + \frac{1}{N} \text{Tr} \Phi^2 \right)^2$$



Coupling

Goldstone decay constant

Collective coupling is controlled by unitarity

$$\alpha N \lesssim 1$$

NB The model is renormalizable!

NNB Double scaling limit $N \rightarrow \infty, \alpha \rightarrow 0$

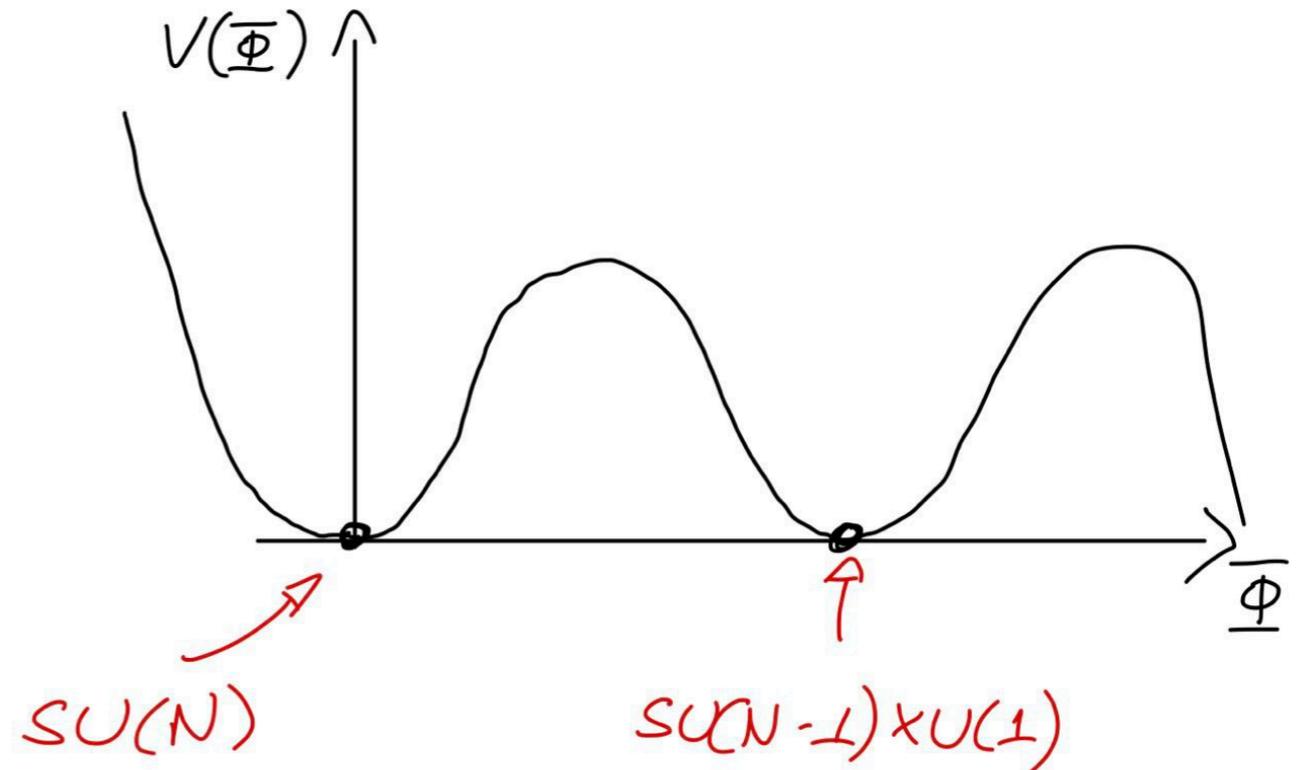
An example of Saturon

Vacuum structure - many degenerate vacua

$$V(\Phi) = 0$$

↓

$$f\Phi - \Phi^2 + \frac{1}{N}\text{Tr}\Phi^2 = 0$$



1) $SU(N)$ symmetric vacuum:

$$\langle \Phi \rangle = 0, \quad \text{mass gap } m = \sqrt{\alpha} f$$

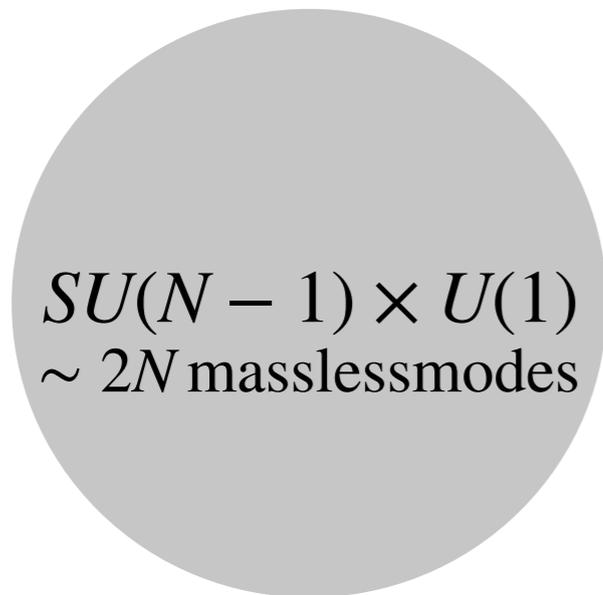
2) $SU(N-1) \times U(1)$ symmetric vacuum

$$\langle \Phi \rangle = \frac{f}{N} \text{diag} (N-1, -1, \dots, -1)$$

$\sim 2N$ gapless modes

An example of Saturnon

Bubble configuration



$SU(N)$
mass gap $m = \sqrt{\alpha} f$

$$\Phi = \frac{\rho(r)}{f} \langle \Phi \rangle$$

Bubble is highly degenerate and can store **information in Goldstone modes**

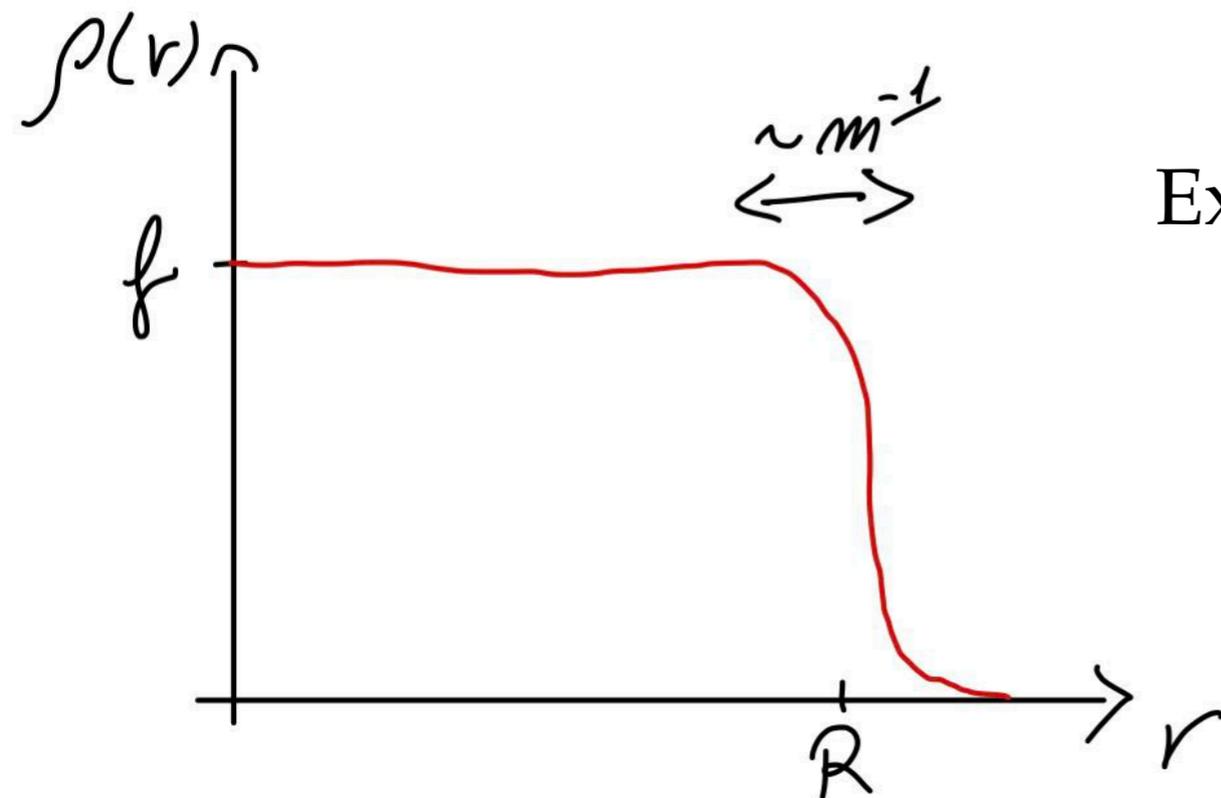


Exciting Goldstones stationarizes the bubble
(adding information)

$$\Phi \rightarrow U^\dagger \Phi U = \frac{\rho(r)}{f} e^{i\omega t \hat{T}} \langle \Phi \rangle e^{-i\omega t \hat{T}}$$

\hat{T} being one of the broken generators

Similar construction in non-topological solitons



An example of Saturon

Configuration saturates unitarity bound i.e., it becomes a **saturon** when

$$\omega \sim m \sim 1/R \quad \text{and} \quad n_{\text{Gold}} \sim N \simeq 1/\alpha$$



$$S \simeq \frac{\text{Area}}{G_{\text{Gold}}} \simeq N \simeq \frac{1}{\alpha}, \quad M \simeq \frac{1}{R} N$$

Analogy with black holes $G_{\text{Gold}} = G_{\text{N}} \quad (f = M_{\text{pl}})$

The saturated bubble also has a thermal spectrum (Hawking) and a notion of information horizon

See Dvali + Dvali, Kaikov, Valbuena '21

An example of Saturnon

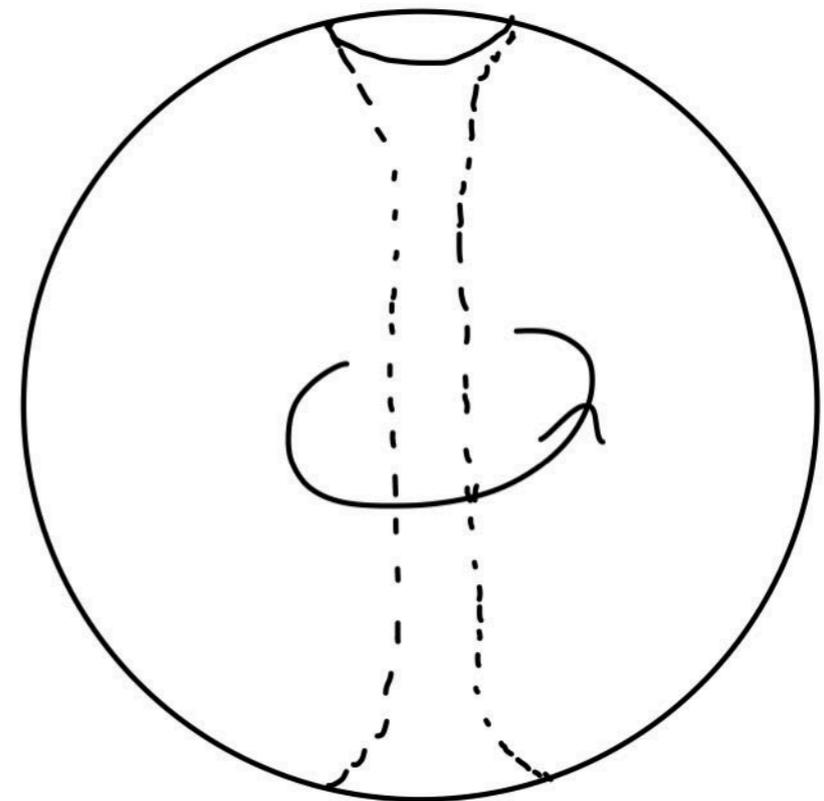
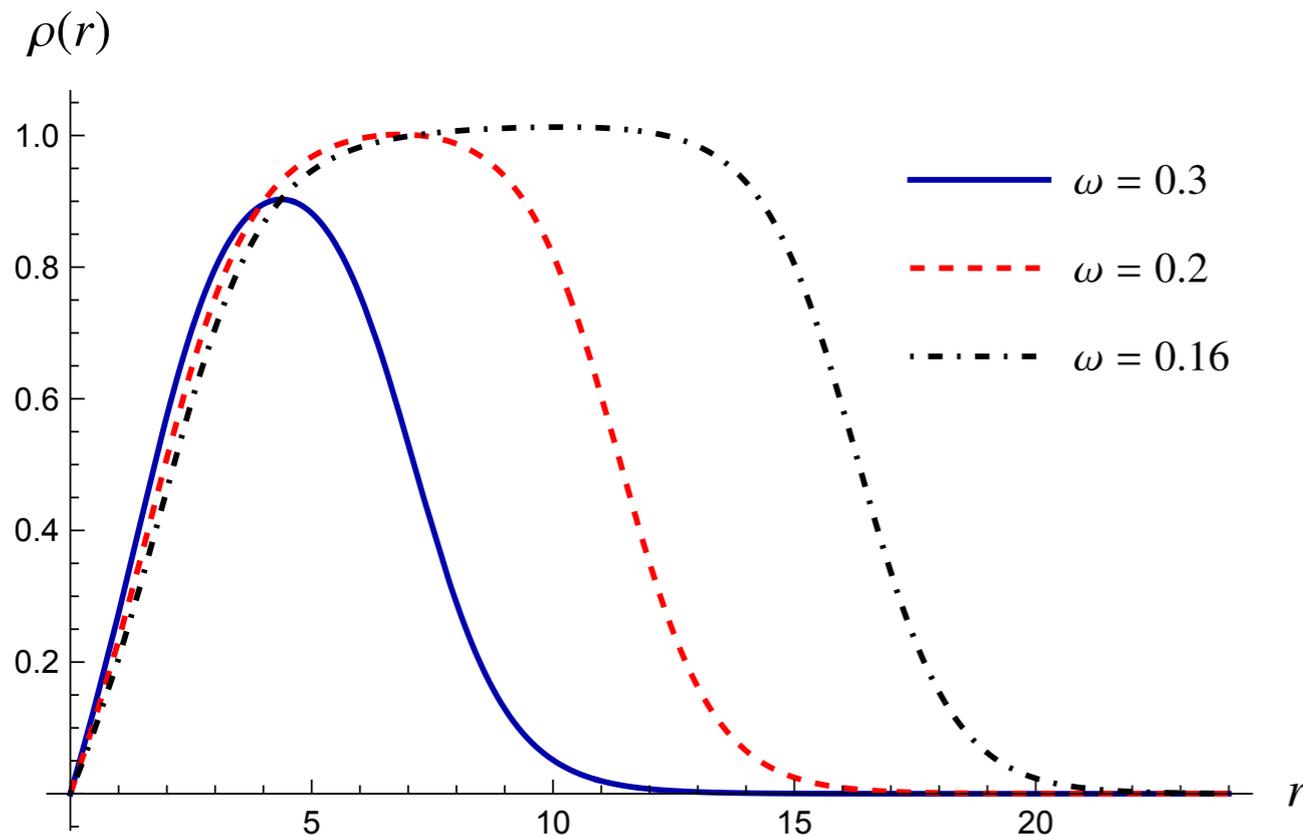
There is a way to spin a saturnon bubble in an axial-symmetric way:

Vorticity

$$\Phi = \frac{\rho(r)}{f} e^{i(\omega t + \kappa \varphi)} \hat{T} \langle \Phi \rangle e^{-i(\omega t + \kappa \varphi)} \hat{T}$$

winding number = $\kappa = 0, \pm 1, \pm 2, \dots$
 $\varphi = \text{polar angle}$

Angular momentum $J \simeq \kappa n_{\text{Gold}}$



Multiplicity in broken generators

Similar study for Q-ball see Volkov, Wohnert '02

An example of Saturnon

Angular momentum $J \simeq \kappa n_{\text{Gold}} \sim \kappa S$ @ saturation

Requiring vortex energy smaller than bubble energy

$$E_{\text{spin}} \lesssim M_{\text{bubble}} \rightarrow \kappa \sim \mathcal{O}(1)$$

An example of Saturon

Angular momentum $J \simeq \kappa n_{\text{Gold}} \sim \kappa S$ @ saturation

Requiring vortex energy smaller than bubble energy

$$E_{\text{spin}} \lesssim M_{\text{bubble}} \rightarrow \kappa \sim \mathcal{O}(1)$$

Saturon bubble

$$J_{\text{saturon}} \lesssim M^2 G_{\text{Gold}} \\ \kappa \sim \mathcal{O}(1)$$

Black hole

$$J_{\text{black hole}} \lesssim M^2 G_{\text{N}}$$

- Saturon and black hole obey the same bound on spin
- For maximally spin, topology prevents Hawking-like emission

Non-gravitational Saturon

- Entropy Area law
- Thermal spectrum
- Information horizon
- Page time
- Maximal spin bound
- **Vorticity**

Black hole



Non-gravitational Saturon

- Entropy Area law
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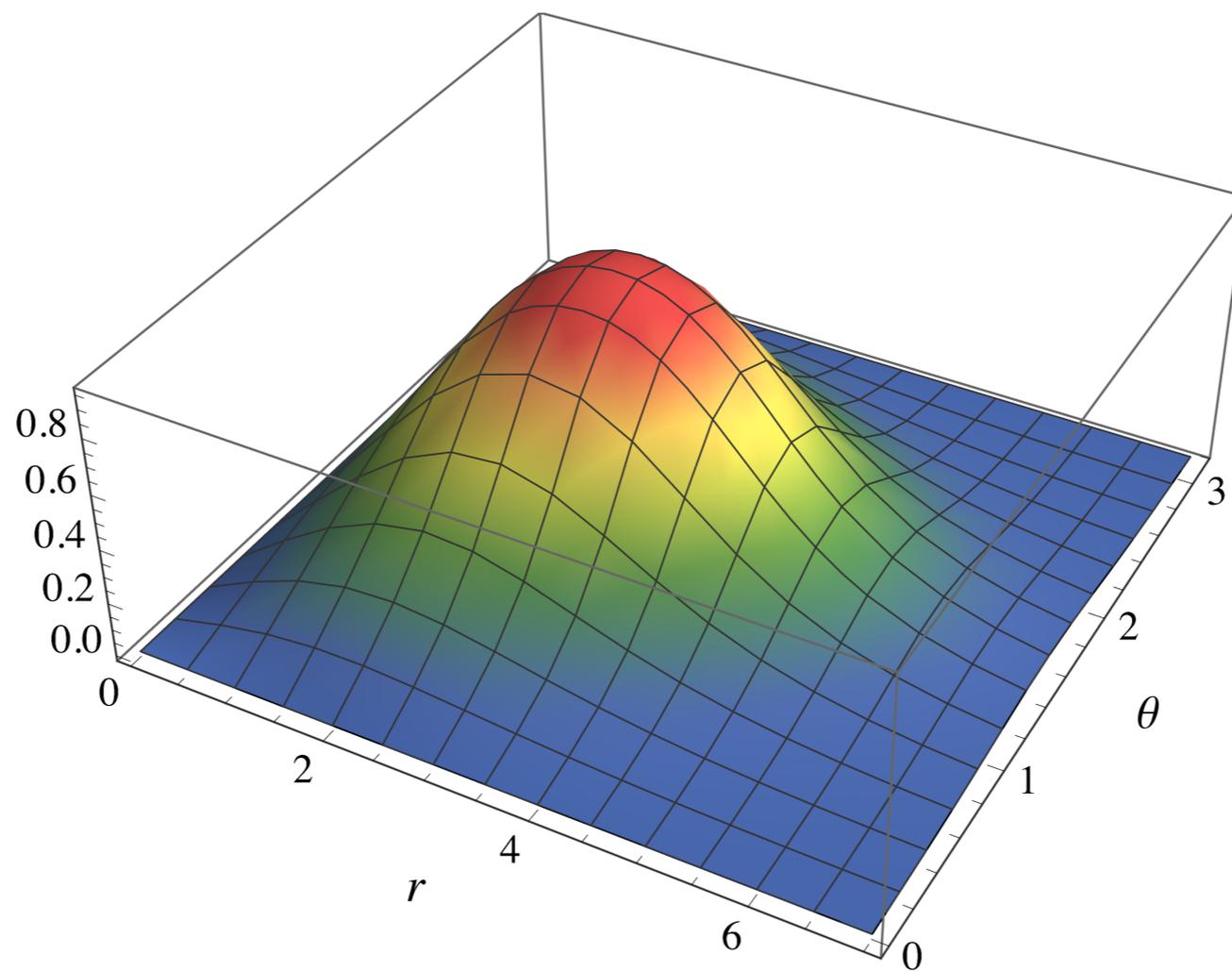
Black hole



Question: *Is vorticity an inherent property of (close-to-)extremal black holes?*

If so, these structures offer a natural support for magnetic field, and can bear pheno consequences

Thank you!



Conjecture: *highly rotating black holes correspond to graviton condensates endowed with vorticity.*

- Motivated also by the black hole N-portrait
- Degeneracy of Goldstone vacuum, responsible for maximal entropy, is in one to one correspondence with the degeneracy of vortex configurations (share same generators)
- This would give a microscopic rationale for the maximal spin bound of black holes
- There can exist non-stationary non-axisymmetric configurations with zero net vorticity but multiple vortices
- Phenomenological consequences possible

Indeed it can be that BHs are very special objects!

Magnetic trapping

Global vortex, when interacting with charged matter under arbitrary $U(1)_{\text{gauge}}$ traps the associated magnetic field

Dvali, Senjanovic '93

Consider two oppositely $U(1)_{\text{gauge}}$ charged field

$$\chi_+ = \rho_+ e^{i\theta_+}, \quad \chi_- = \rho_- e^{i\theta_-}$$

Effectively coupled to the vortex order parameter $\psi = \rho e^{i\theta}$ as

$$\psi\chi_+\chi_- + \text{h.c.} = 2\rho\rho_+\rho_- \cos(\theta + \theta_+ + \theta_-)$$

Asymptotically for the gauge field

$$A_\mu = \frac{1}{eq} \frac{\langle \rho_+ \rangle^2 \partial_\mu \theta_+ - \langle \rho_- \rangle^2 \partial_\mu \theta_-}{\langle \rho_- \rangle^2 + \langle \rho_+ \rangle^2}$$

Magnetic trapping

$$\text{Flux} = \int dx^\mu A_\mu = \frac{2\pi}{eq} \left[\kappa_+ + \kappa \frac{\langle \rho_- \rangle^2}{\langle \rho_- \rangle^2 + \langle \rho_+ \rangle^2} \right]$$

$$\kappa_\pm = \frac{1}{2\pi} \int dx^\mu \partial_\mu \theta_\pm \quad \kappa_+ + \kappa_- = -\kappa$$

E.g., charged components of neutral plasma or charged dark matter

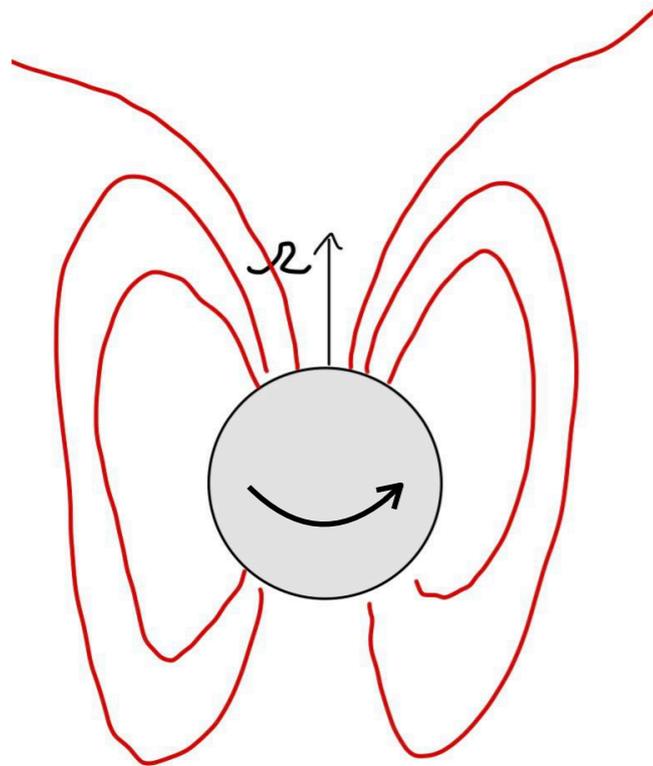
Fractional flux has further stabilising properties

If $U(1)_{\text{gauge}} = U(1)_{\text{em}}$ natural support for magnetic field

Black hole pierced by magnetic field lines

Intermezzo: Blandford Znajek emission

Highly rotating BHs endowed with a magnetosphere produce extremely powerful jets (BZ)



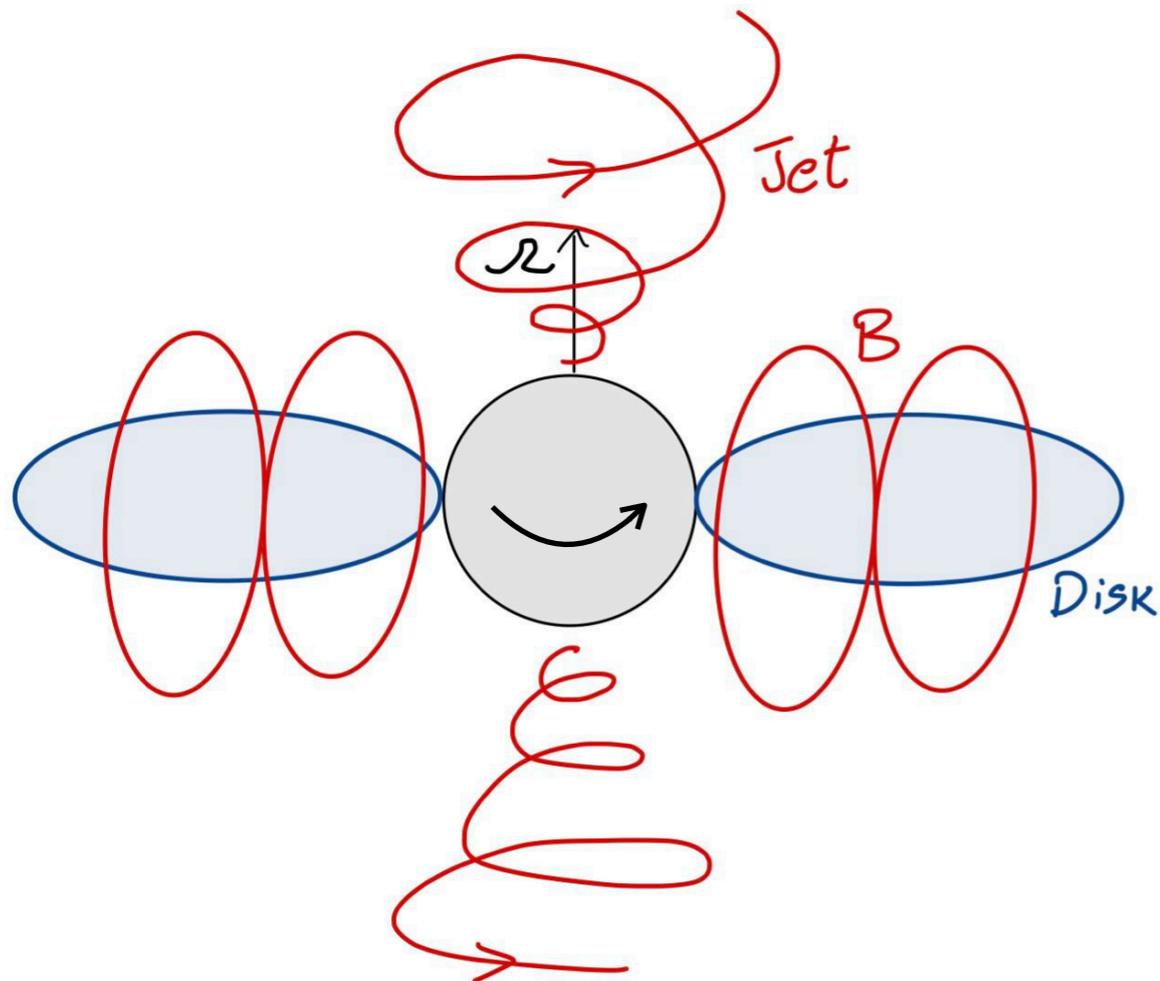
$$P_{\text{BZ}} \sim \text{Flux}^2 \Omega^2$$

Classically, a BH cannot have magnetic hairs.
The mechanism providing magnetic field remains to present day a mystery (although BZ-like emissions are observed).

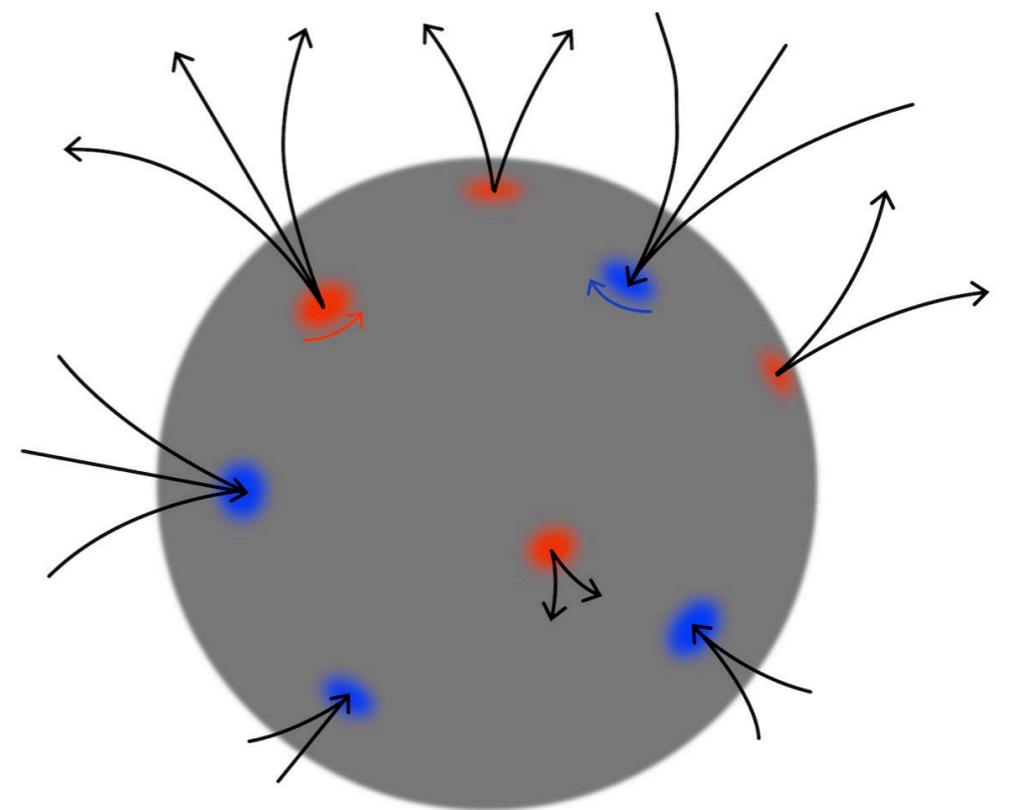
Intermezzo: Blandford Znajek emission

Highly rotating BHs endowed with a magnetosphere produce extremely powerful jets (BZ)

Traditionally it is assumed the presence of an accreting disk with very specific coherent magnetization which can temporarily source a magnetic field on the BH



A possible source could also be vorticity



Magnetic trapping

$$\text{Flux} = \int dx^\mu A_\mu = \frac{2\pi}{eq} \left[\kappa_+ + \kappa \frac{\langle \rho_- \rangle^2}{\langle \rho_- \rangle^2 + \langle \rho_+ \rangle^2} \right]$$

Jet emission (e.g., à la Blandford Znajek) can take place **without the need of an accreting magnetized disk** providing a **smoking gun** for the scenario

Example: milli-charged dark matter

Consider the axial-symmetric solution found before and $eq \sim 10^{-39}$

$$P_{\text{BZ}} \sim \text{Flux}^2 \Omega^2 \sim P_{\text{M}_{87}} \sim 10^{44} \text{erg s}^{-1} \quad \text{for maximally rotating BH}$$

$$\text{Reminder: } \kappa_+ + \kappa_- = -\kappa \sim \mathcal{O}(1)$$

Maximal flux

Maximal magnetic field for BH is restricted by $E_{\text{magnetic}} < M_{\text{BH}}$

$$B \sim \text{Flux}/R_{\text{BH}}^2 \lesssim M_{\text{pl}}/R_{\text{BH}} \sim M_{\text{pl}}^2/\sqrt{S}$$

Which in turn implies a bound on maximal flux

$$\text{Flux} \sim \frac{1}{eq} \lesssim \sqrt{S}$$

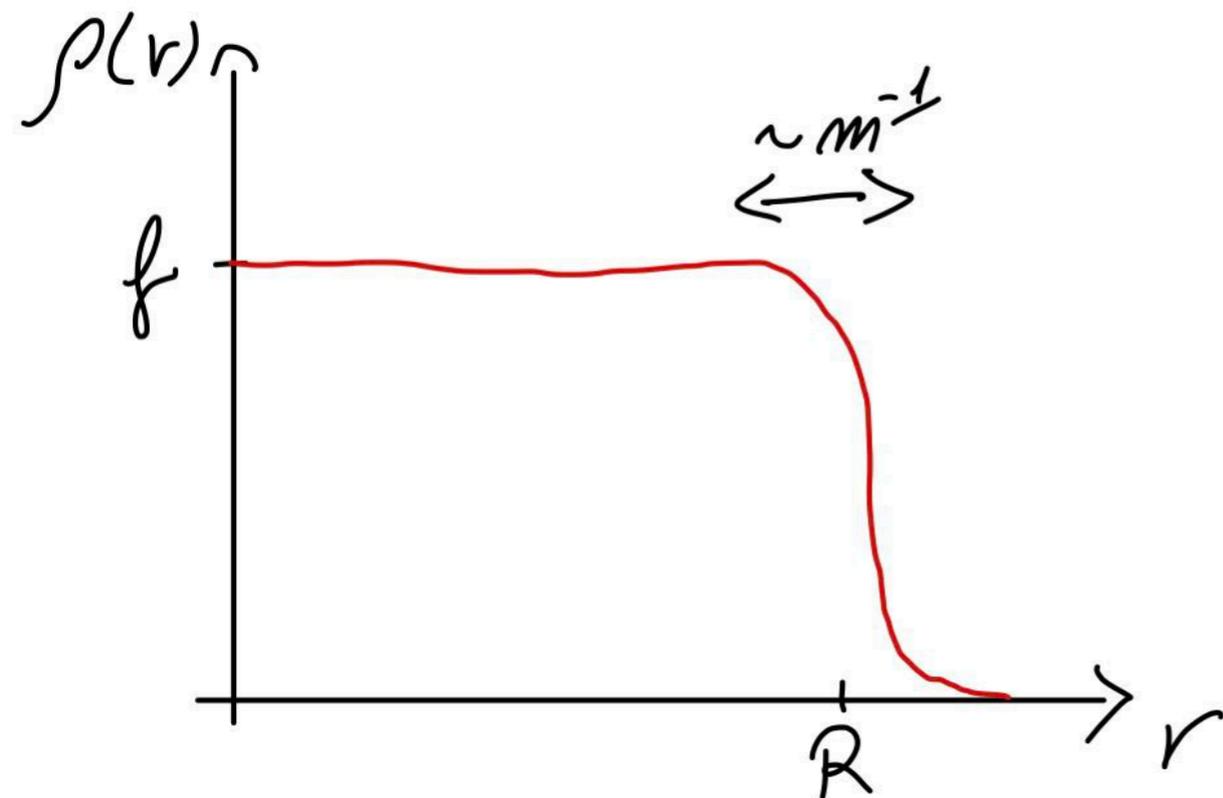
For M87 case $\sqrt{S} \sim 10^{47}$ (we used $eq \sim 10^{-39}$)

An example of Saturnon

$$d_r^2 \rho + \frac{2}{r} d_r \rho + \rho [\omega^2 - \alpha(\rho - f)(2\rho - f)] = 0$$

$$Q = n_{\text{Gold}} \sim \frac{1}{\alpha} \left(\frac{m}{\omega} \right)^5 \quad R \sim \frac{m}{\omega^2} \quad E \sim \omega n_{\text{Gold}}$$

$$n_{\text{states}} \simeq \left(1 + \frac{2N}{n_{\text{Gold}}} \right)^{n_{\text{Gold}}} \left(1 + \frac{n_{\text{Gold}}}{2N} \right)^N \quad S = \log n_{\text{states}}$$



Where the number of states is counted as number of configurations of form

$$|\text{Pattern}\rangle = |n_{\omega}^1, n_{\omega}^2, \dots, n_{\omega}^{2N}\rangle$$

under the constraint

$$\sum_a^{2N} n_{\omega}^a = n_{\text{Gold}}$$

An example of Saturon

Properties

$$n_{\text{Gold}} \sim \omega f^2 R^3 \sim \frac{1}{\alpha} \left(\frac{m}{\omega} \right)^5 \quad R \sim \frac{m}{\omega^2} \quad E \sim \omega n_{\text{Gold}}$$

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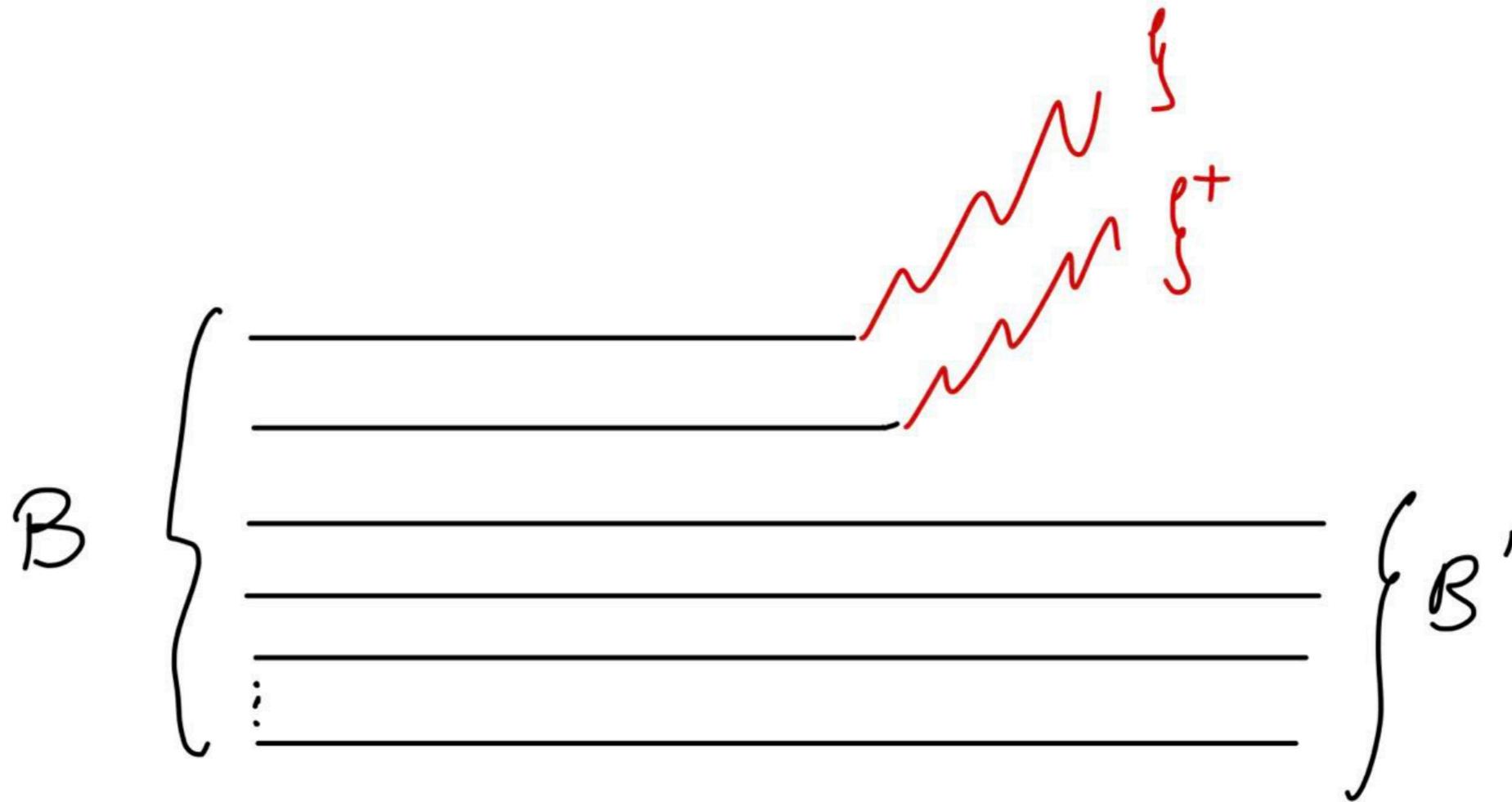
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Remember: At saturation

$$\omega \sim m \sim 1/R \quad \text{and} \quad n_{\text{Gold}} \sim N$$

Hawking emission

Following Hawking couple massless field ξ

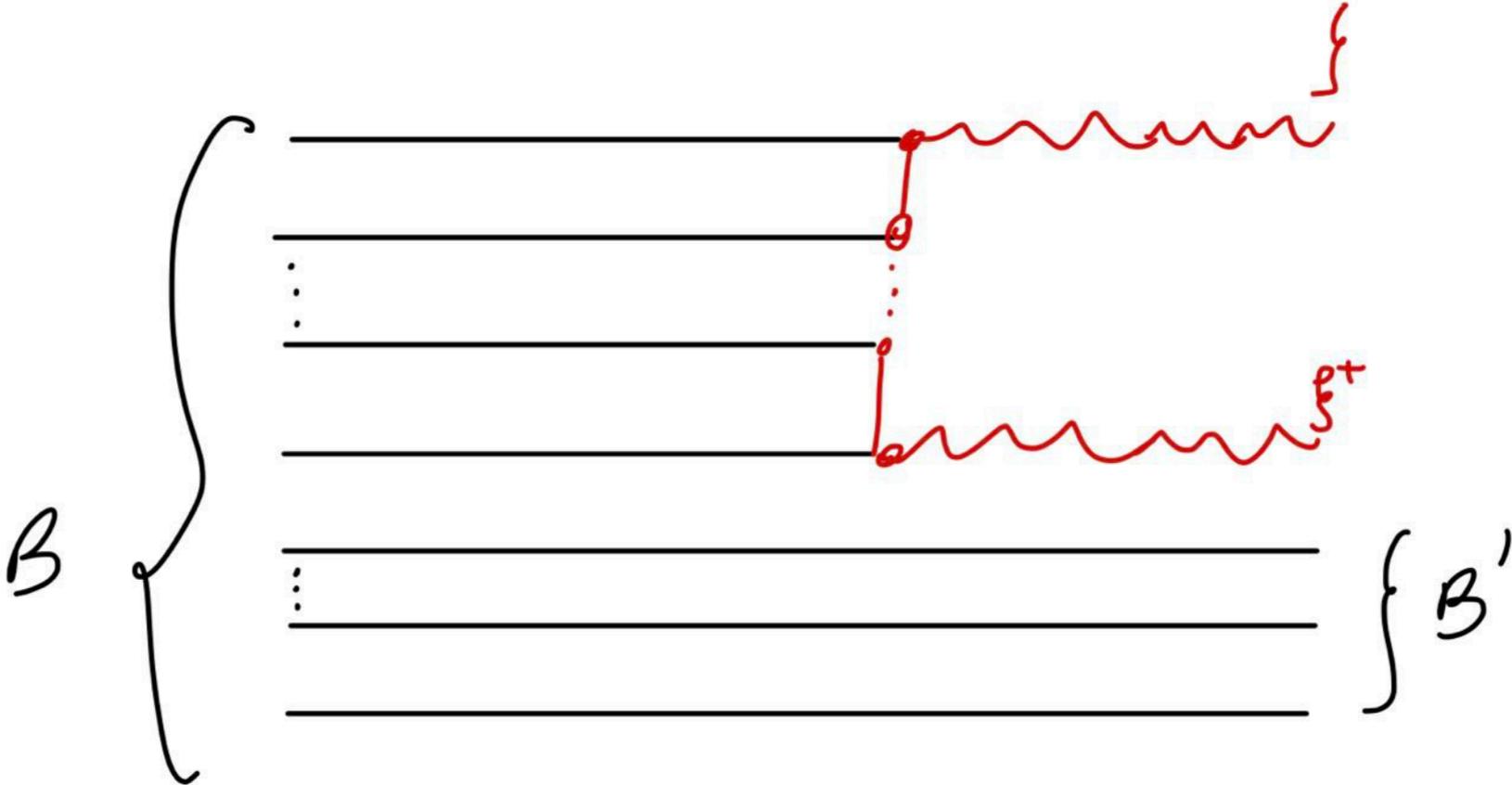


$$\Gamma \sim (\alpha N)^2 \frac{1}{R} \sim \frac{1}{R}$$

$\alpha N \simeq 1$ at saturation

Hawking emission

Temperature emerge as rescattering of multiple constituents



$$E_\xi \gg 1/R$$

$$\Gamma \propto e^{-E_\xi/T} \text{ Boltzmann suppression}$$

More on millicharged DM

Charged cold dark matter generate photon mass

$$m_\gamma = eq\sqrt{\rho_{\text{local}}}/m_{\text{DM}}$$

On galactic scales photon mass is constrained by galaxy magnetic field

$$m_\gamma < \sqrt{eqB_{\text{galaxy}}}$$

Example from before ($eq \sim 10^{-39}$):

$$m_\gamma \sim m_{\text{DM}} \sim 10^{-18}\text{eV}, \quad \rho_{\text{local}} \sim 10^4\text{eV}^4$$

Close to fuzzy-like DM

Maximal flux

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