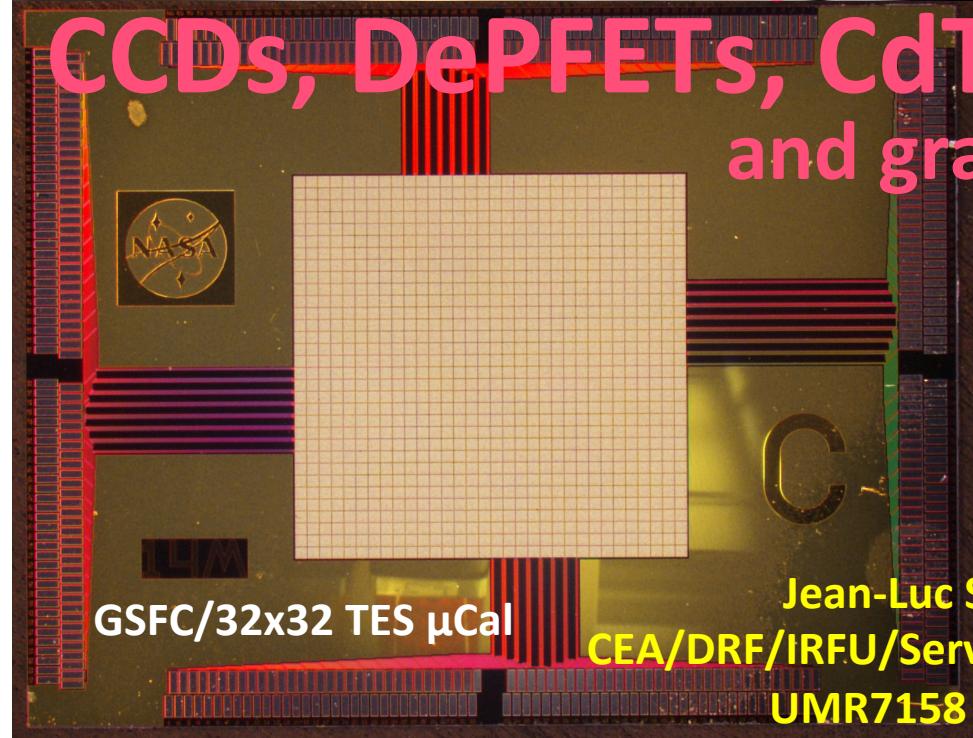


XMM/EPIC/CCDs

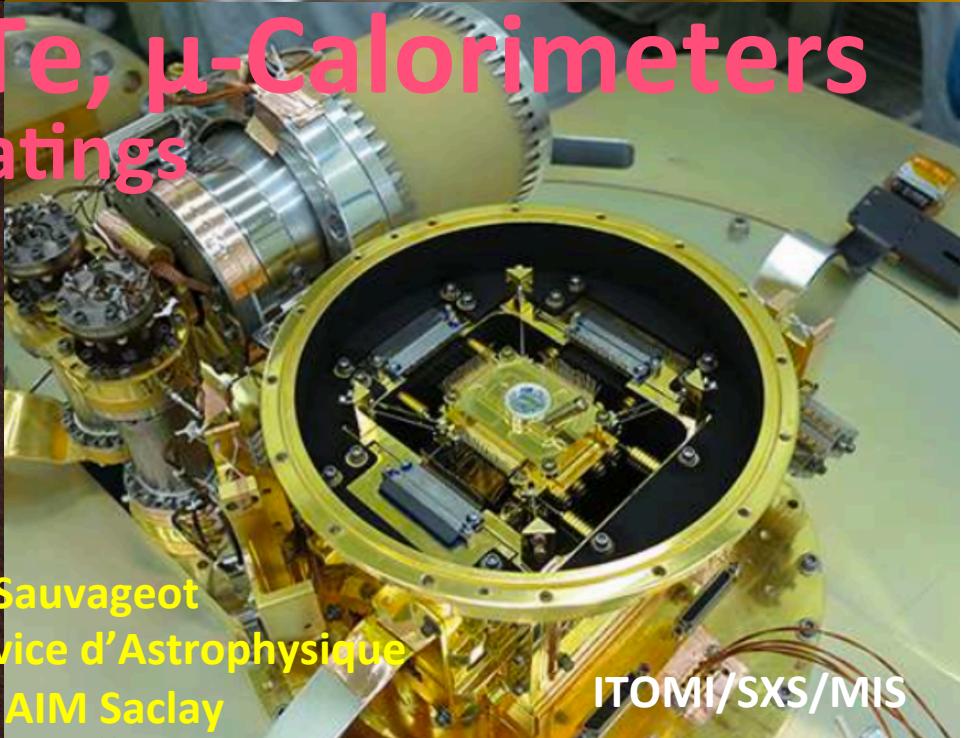


ATHENA/WFI/DePFet



GSFC/32x32 TES μ Cal

Jean-Luc Sauvageot
CEA/DRF/IRFU/Service d'Astrophysique
UMR7158 AIM Saclay

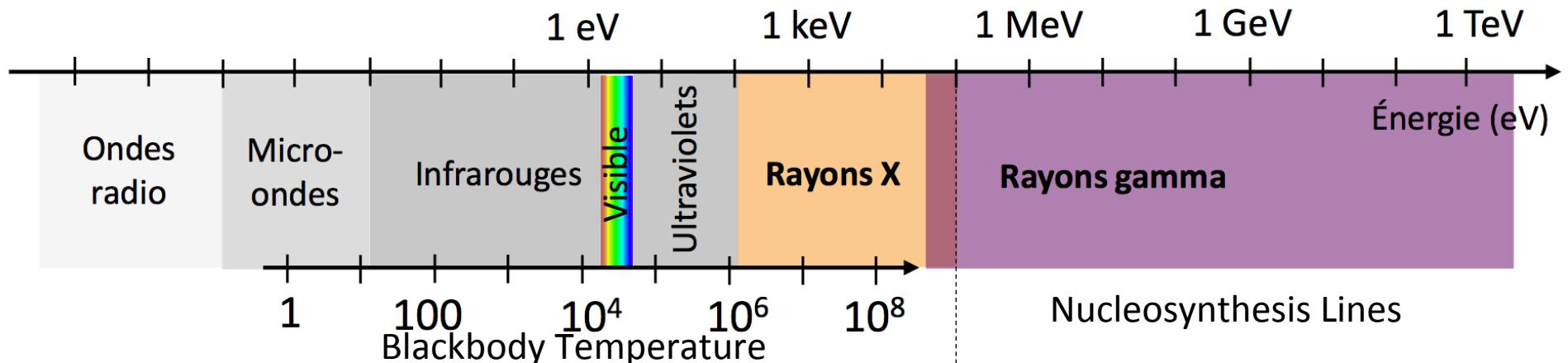


ITOMI/SXS/MIS

- I. [CCDs](#) in X-Rays or how do we use CCDs to obtain X-ray Photon List?
- II. How [DePFETs](#) array improve onto the CCDs observations?
- III. The [CdTe](#) at higher Energies.
- IV. [Gratings](#): Wavelength dispersive spectroscopy against Fano factor.
- V. The [μ-Calorimeters](#) & their exquisite spectral resolution.

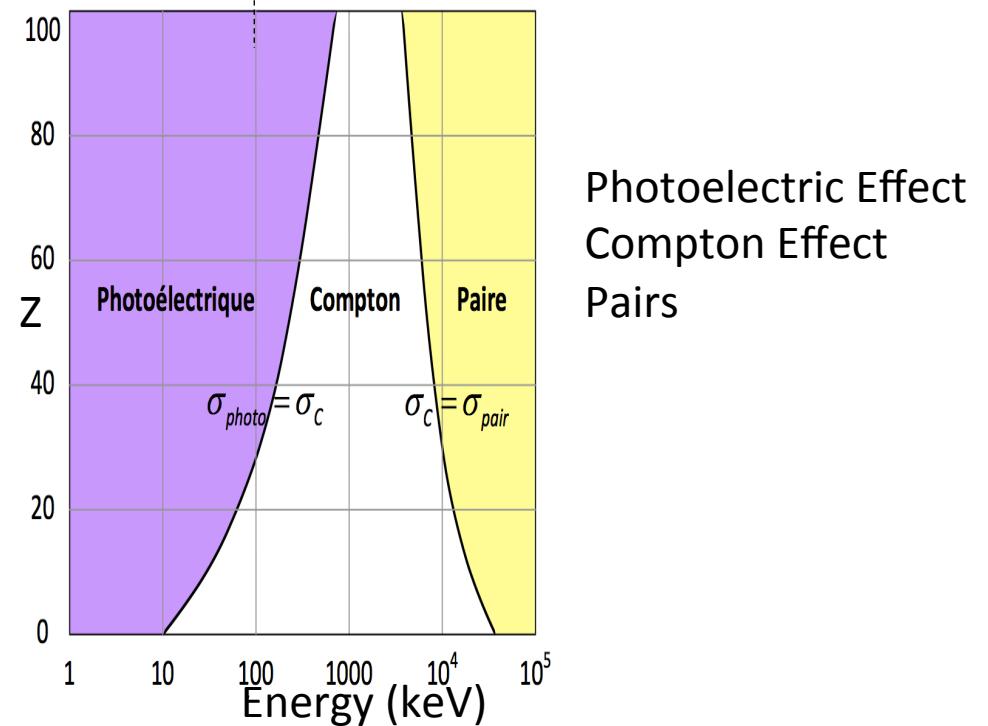
I + II ~50min, III 10min, IV 2min & V 50min

X-ray and Matter



X-ray range
 $200\text{eV} < E < 100\text{keV}$

dominated by **PhotoElectric Effect**



X-rays & Astronomy

- Large Mirror ?

ATHENA : $S_{\text{eff}} \leq 1\text{m}^2 @ 1\text{keV}$ & $S_{\text{eff}} \leq 0.1\text{m}^2 @ 10\text{keV}$

Not so « Large »

- Celestial Sources are relatively weak in X-rays photon flux...
- Thus, as long as the detector is fast enough ($\leq 1\text{ms}$) X-ray observations are always in photon by photon Mode

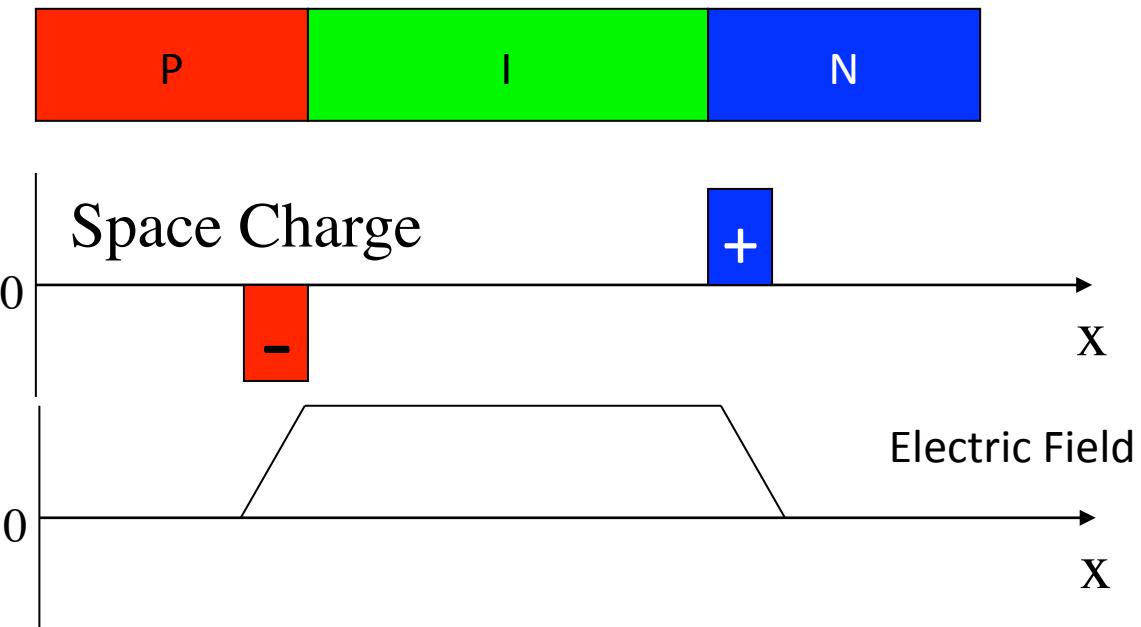
⇒ A spectro-imagery X-ray observation in Astronomy is a photon list with X,Y,En, Arrival Time.

« Historic » X-rays detector

X-ray Astronomy began in the 60 's ...

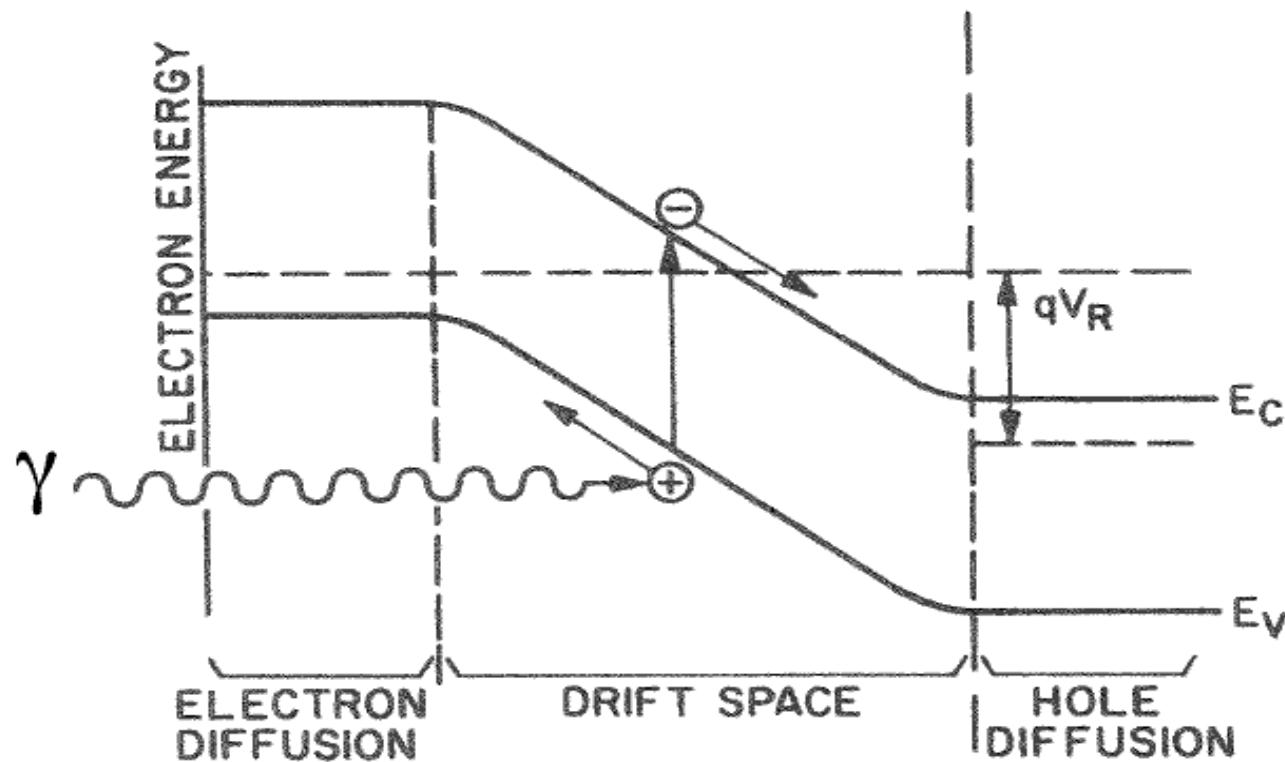
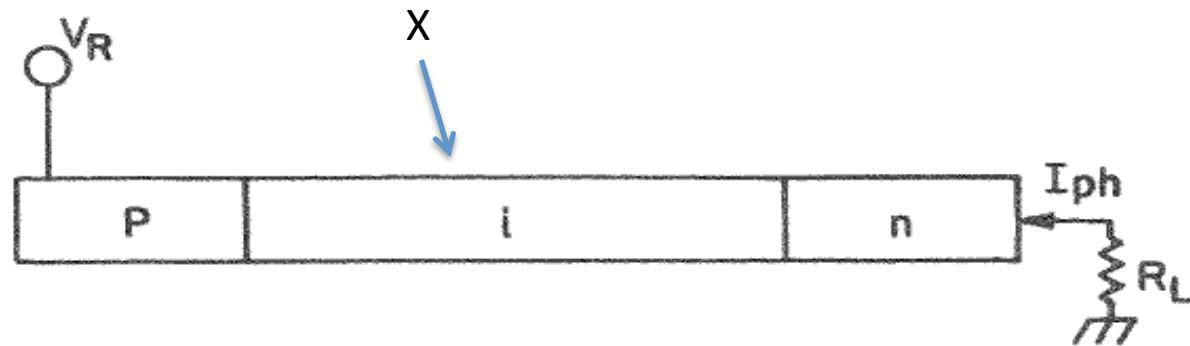
- Gaz Counter (~1960)
- Chanel Multiplier Array
- Gas Scintillating Proportional Counter (1970-1990)
- Solid State Spectrometer (1985)
- Charge Coupled Device (1990-2016)

I.PN-CCDs are p-i-n reversed biased diode



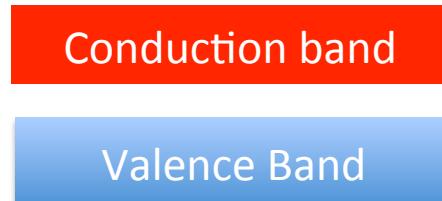
Reverse biased diode & Low Temperature
allow to have very small leakage current.
Leakage current $\ll 1\text{pA}$

I. CCDs e^- -hole drift



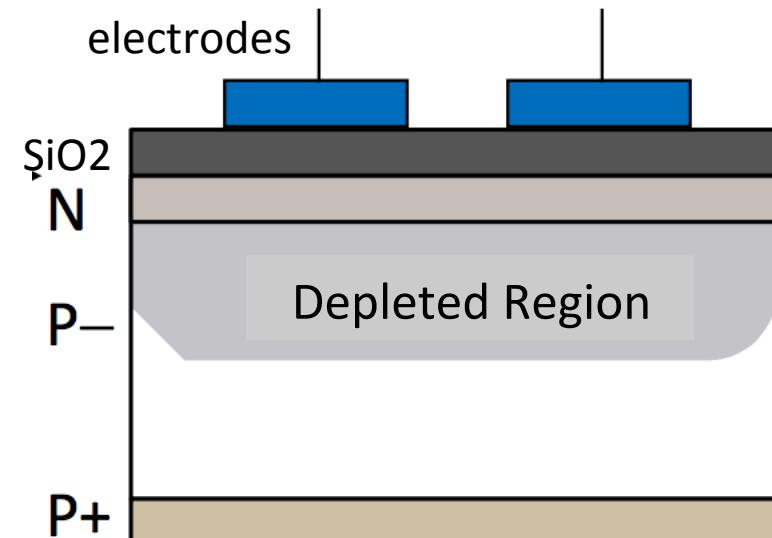
I. CCDs Internal Structure

- X-ray detection by Semi-Conductors
- Band-Structure Conduction & Valence Band

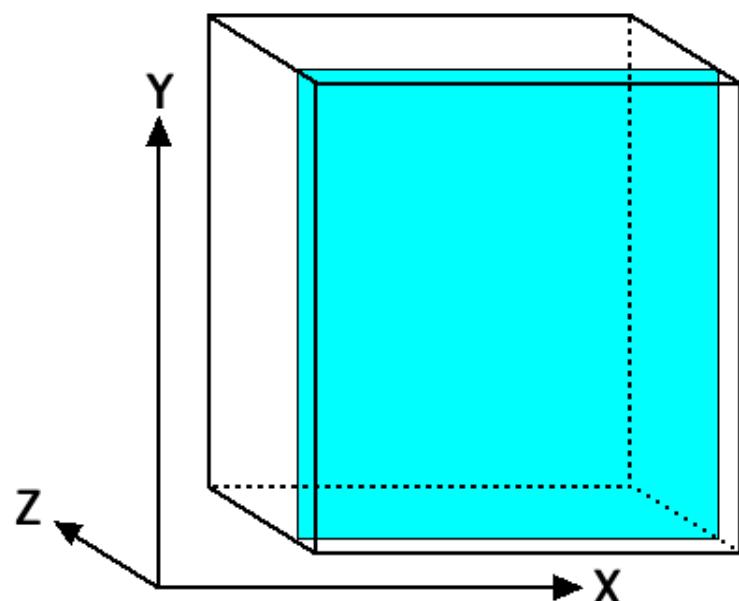


1. Photons → PhotoElectric Effect
2. Charges accumulated in the depleted region

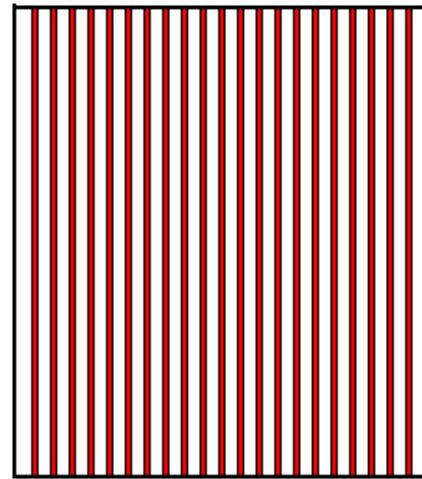
Each X interacting in the depleted region will create e^- -hole pairs that will be separated by the electric field



I. CCDs in 3 steps !

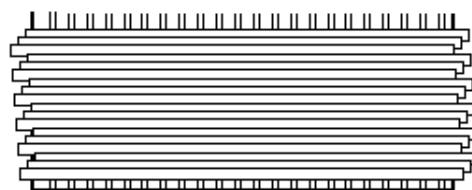


1
n+ (P)
doped layer
Homogeneous
potential
well



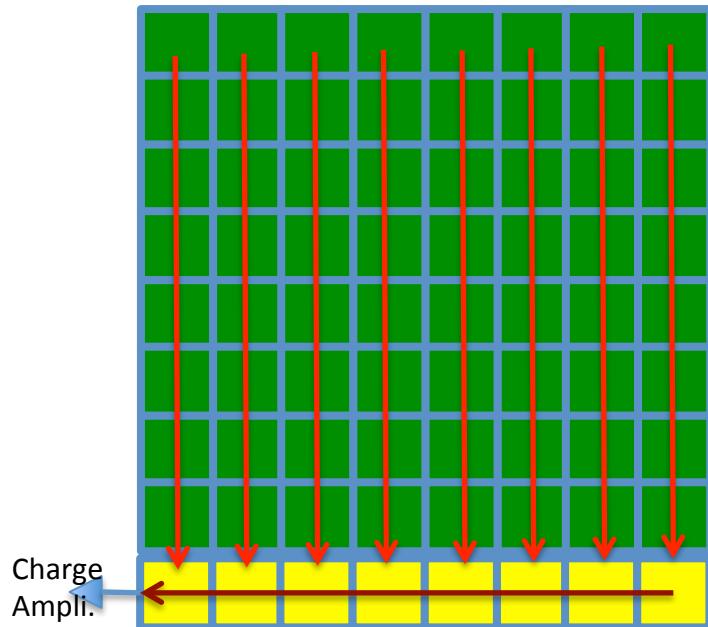
2
p+ (B)
doped thin
strips
Minimum
potential
well defining
column

3
Metallic
Surface
Electrode.
Potential
well
underneath



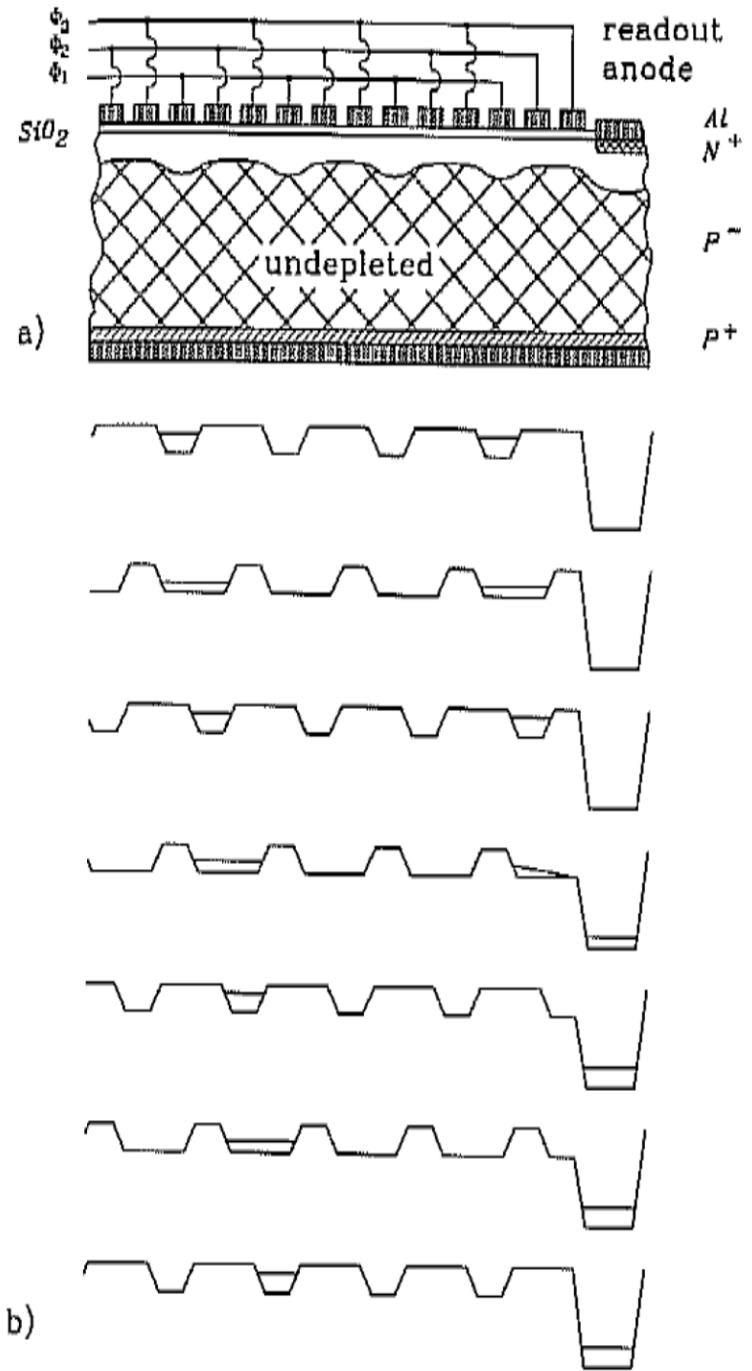
I. CCDs Readout

3/4 electrodes / pixels
(1 collection and 2 fences
+ 2 channel stops (P doped) perpendicular

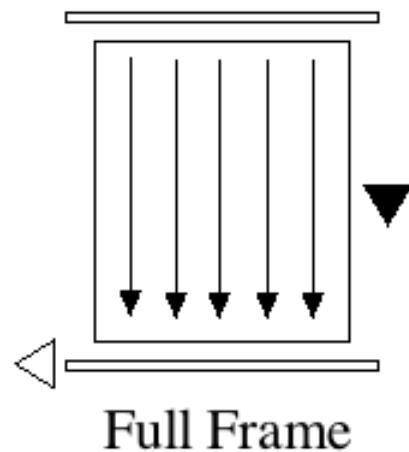


1. Readout by transfer from one pixel to its neighbour.
2. Each Charge is measured by a charge amplifier

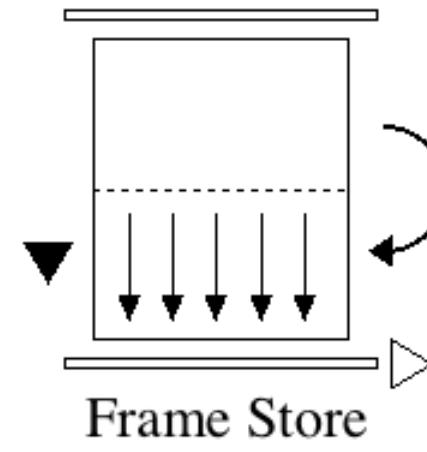
$$\text{Transfer Efficiency} = 1 - ((N_{n+1} - N_n)/N_n) = 0.999999 !$$



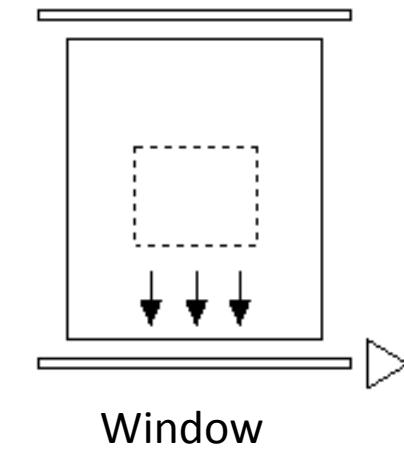
I. CCDs Modes



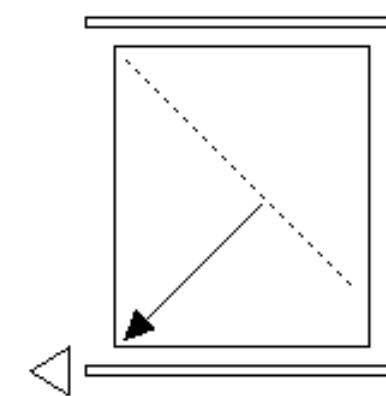
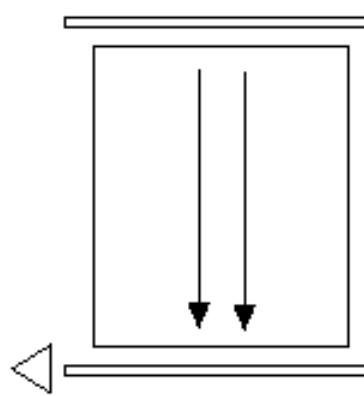
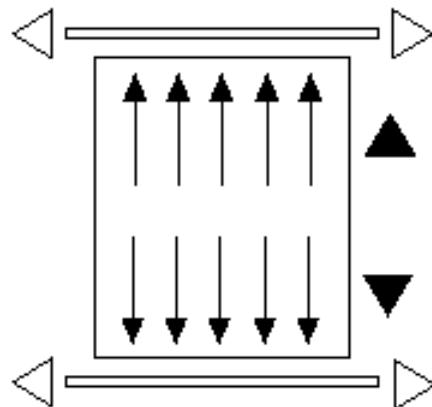
Full Frame



Frame Store



Window



I. CCDs in X-Rays

Xrays interaction → one photo-electron produced → ionisation
→ e⁻-hole pairs produced by photoelectric Effect.

CCDs are pixelated → X,Y information

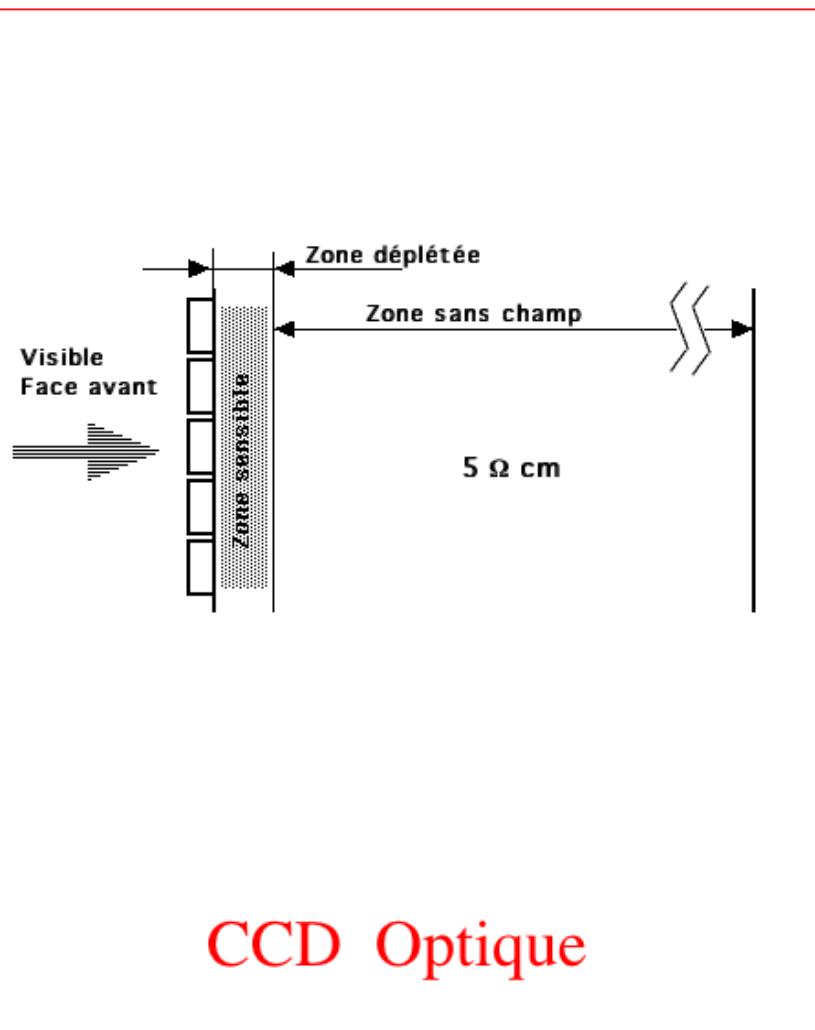
Arrival Time is known through Frame time

Spectral resolution is obtain by counting the holes or e⁻

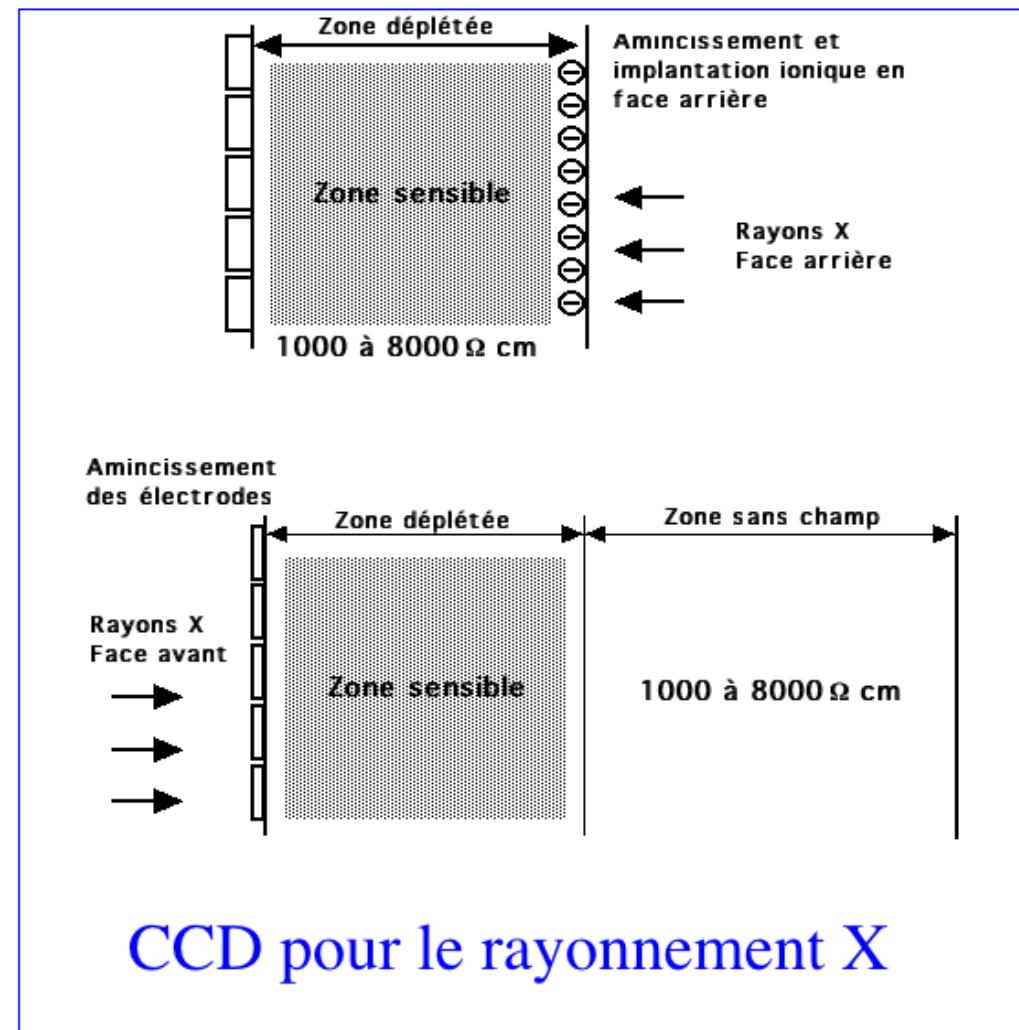
To have spectral information, one needs **not** to have
2 photons on the same pixel in the same frame

→ CCDs are read as fast as possible to avoid Pile-up

Visible vs X-rays CCDs



CCD Optique

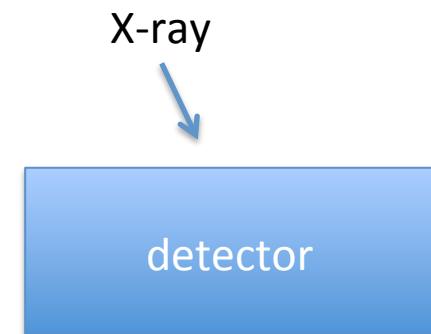


CCD pour le rayonnement X

I. CCDs Spectral Resolution



For monochromatic X-rays, the exact number of electrons produced depend on the way the primary photo-electron loose its energy...



Each e^- in the conduction band carry $\sim E_{\text{gap}}$ energy

$$\Delta E (\text{FWHM}) = 2.36 \varepsilon (\sigma_{\text{readout}}^2 + \sigma_{\text{coll}}^2 + F_{\text{SI}} E_x / \varepsilon)^{1/2}$$

σ_{readout} : Readout Noise (in e^-)

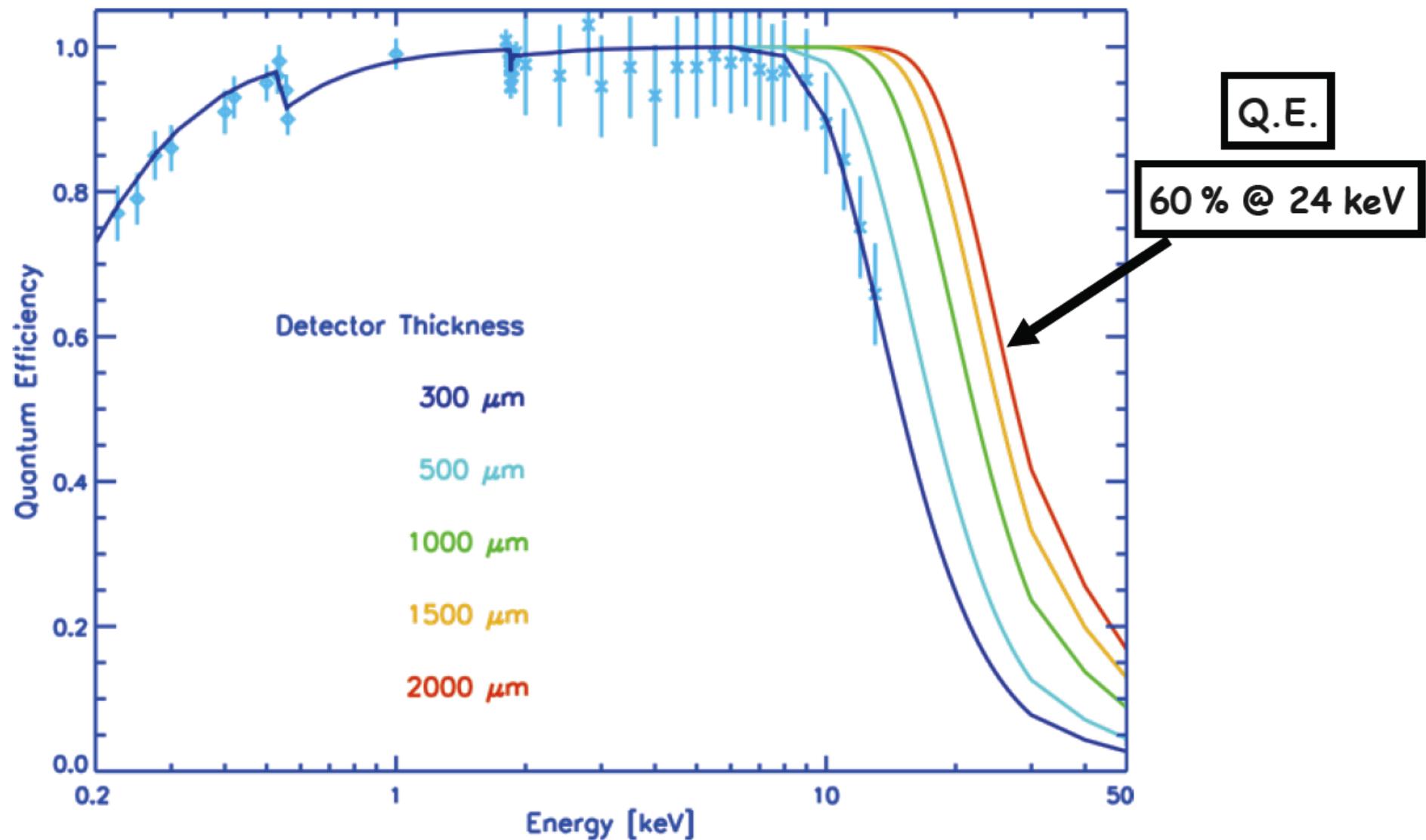
σ_{coll} : Noise from incomplete collection (in e^-)

F_{SI} : Fano Factor (~ 0.11)

ε : Energy to create a pair ($\sim 3.65 \text{ eV}$)

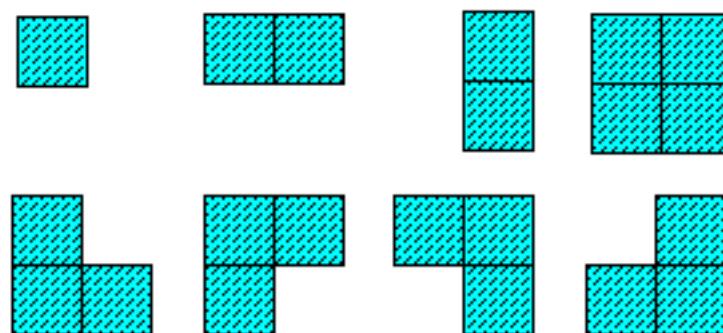
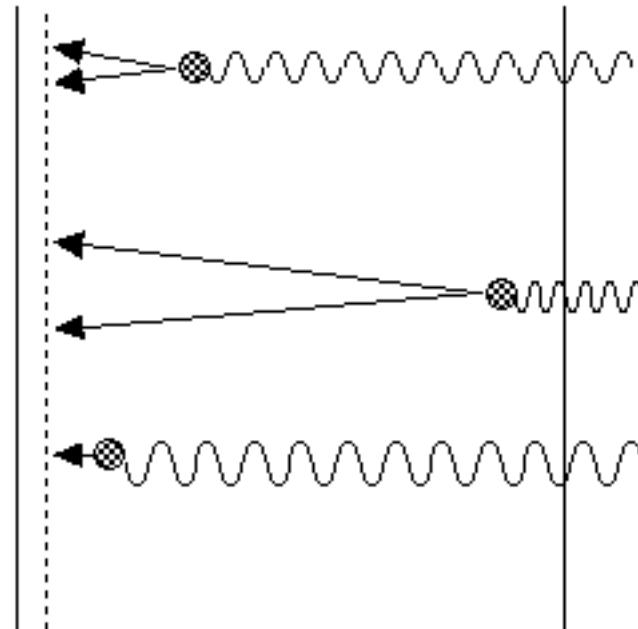
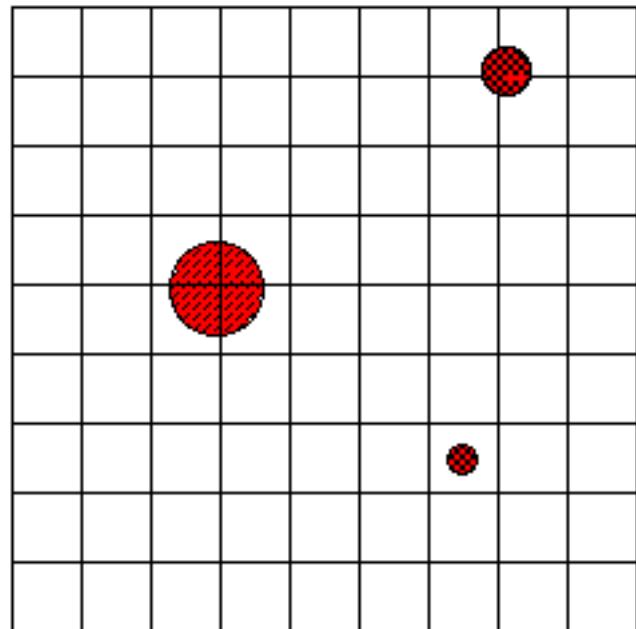
$\rightarrow \Delta E = 120 \text{ eV} @ 6 \text{ keV}$

I. CCDs Quantum Efficiency

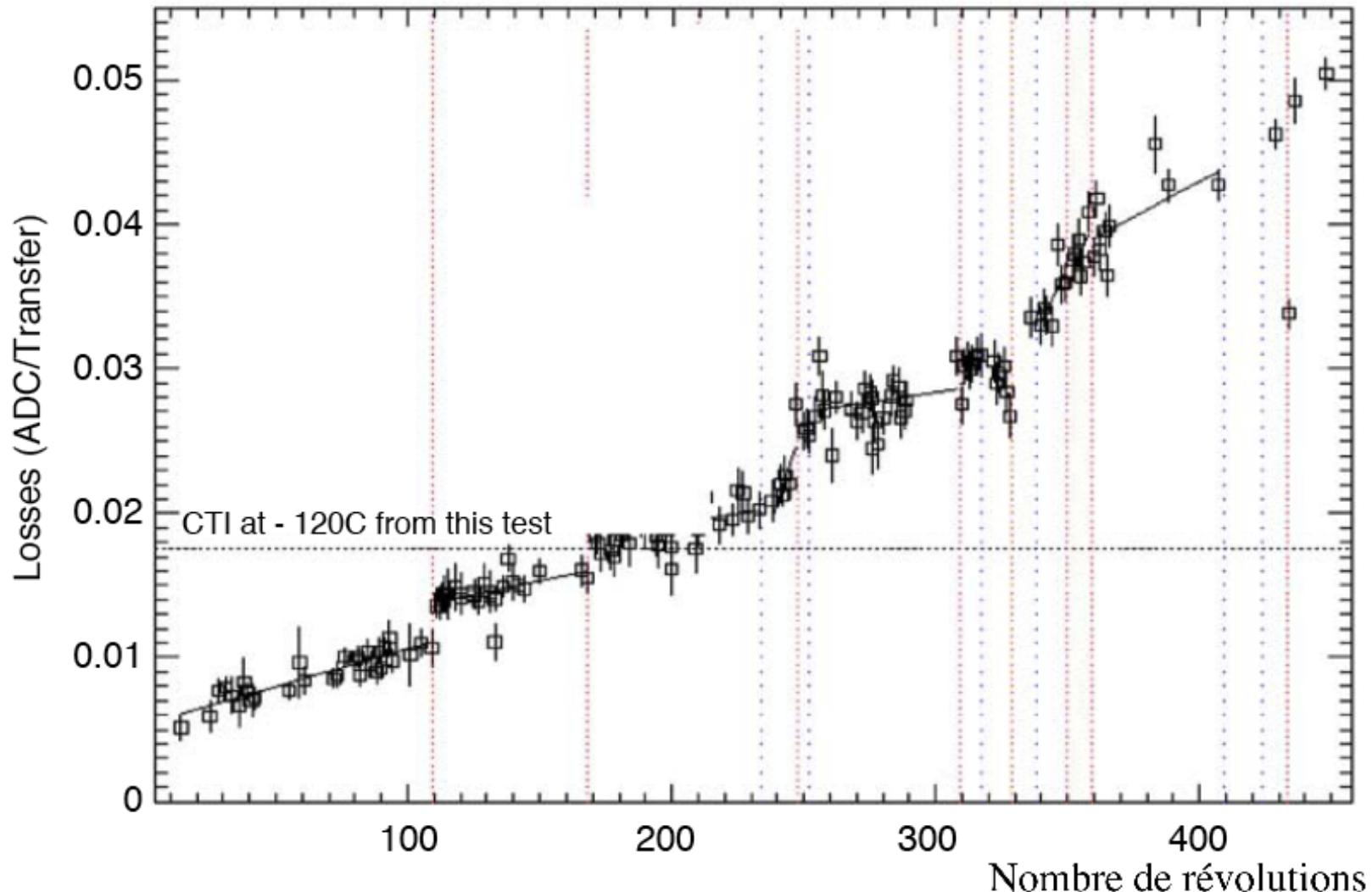


Spatial data flow optimized for X-rays

CCDs: The Event Detection Unit

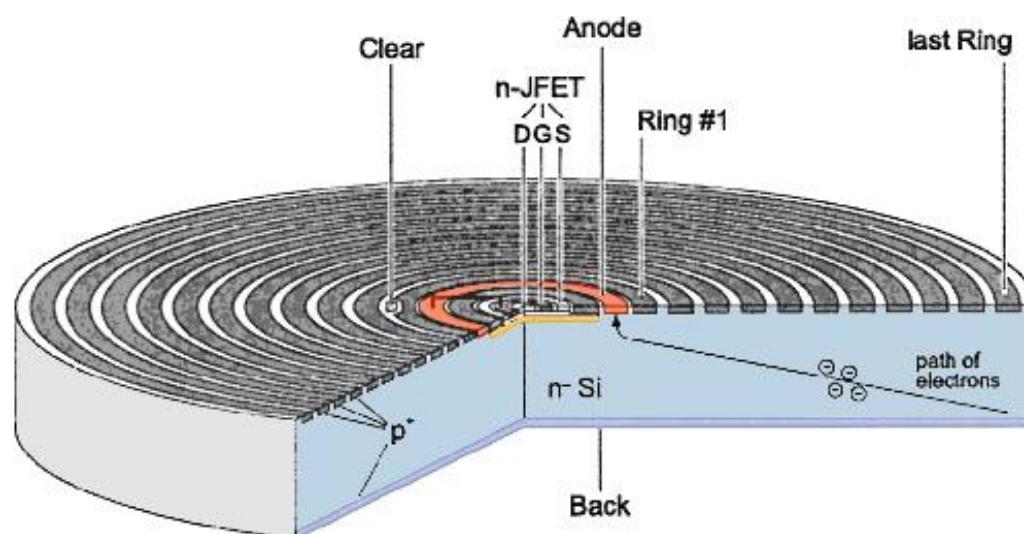
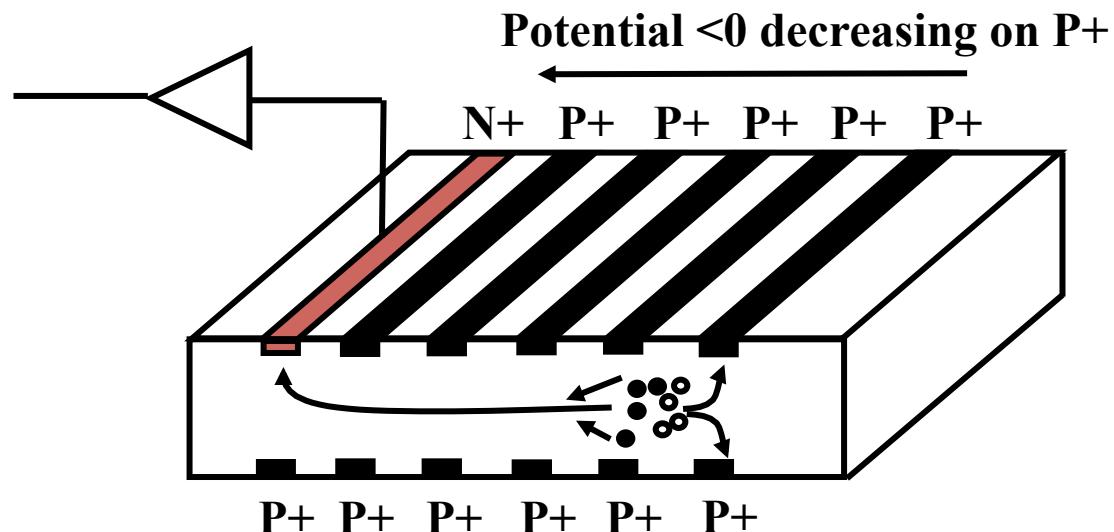


Charge Transfert inefficiency



Degradation of the charge transfert efficiency of the XMM/EPIC_MOS CCDs
Onboard XMM-Newton

II. Silicon Drift Detector (SDD)



II. SDD DEPFET

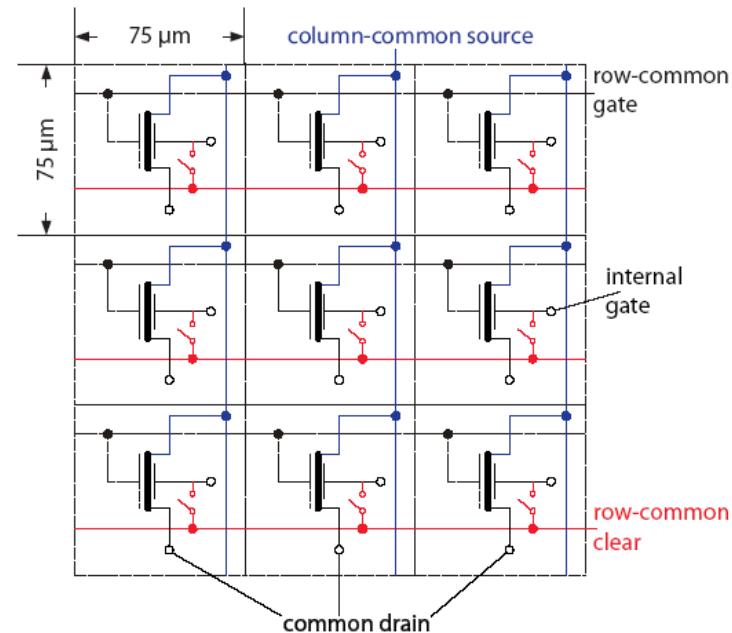
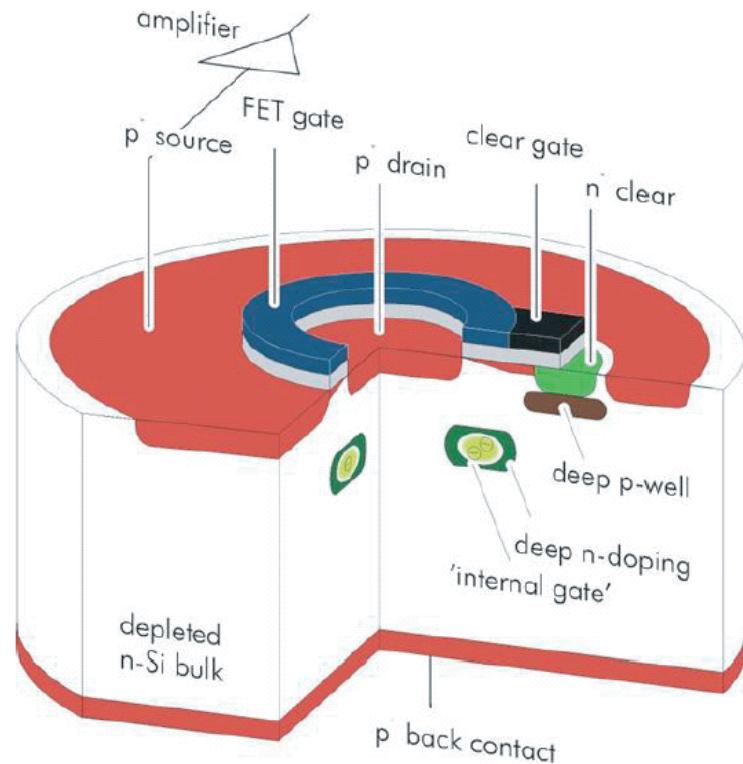
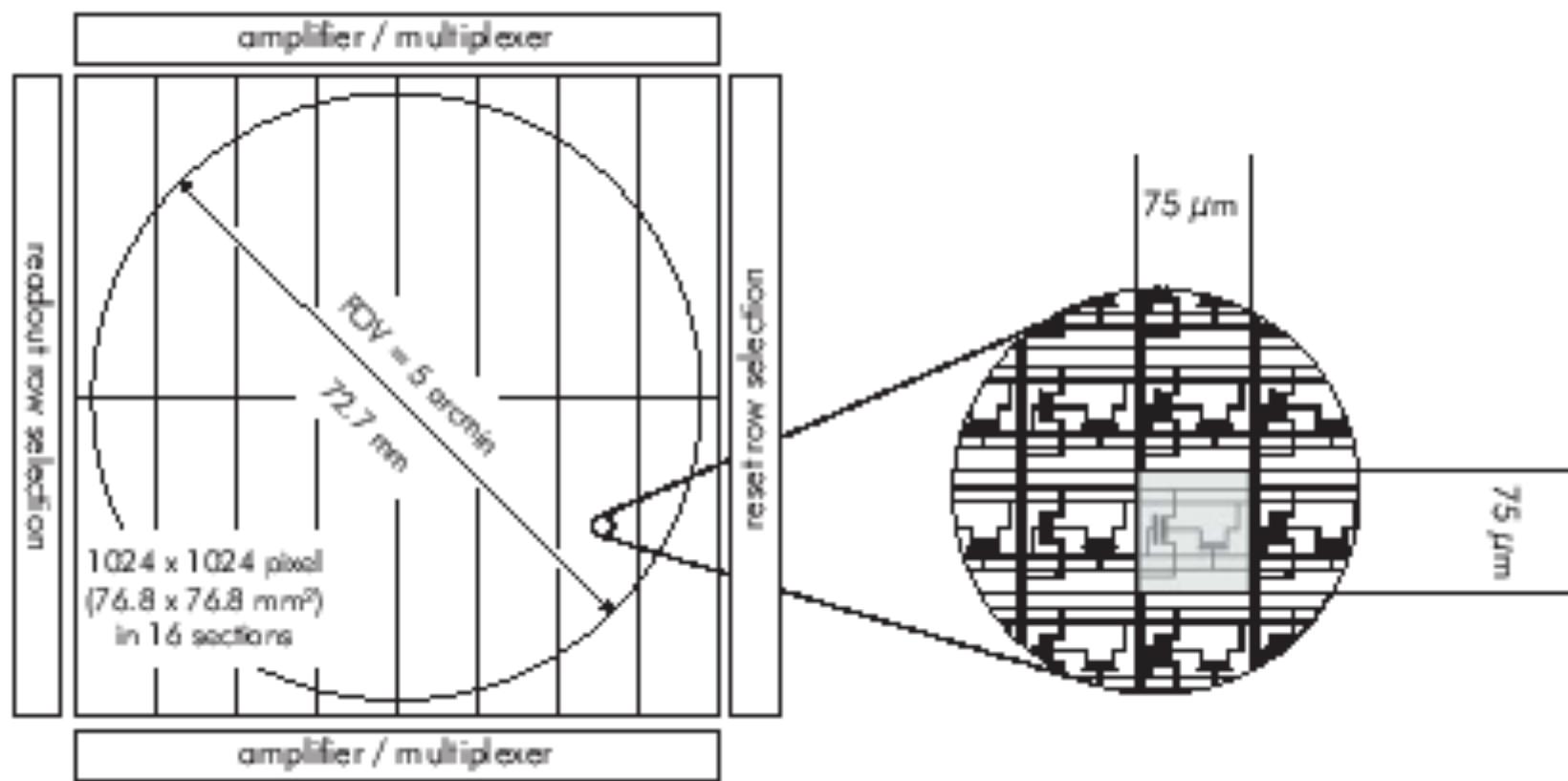


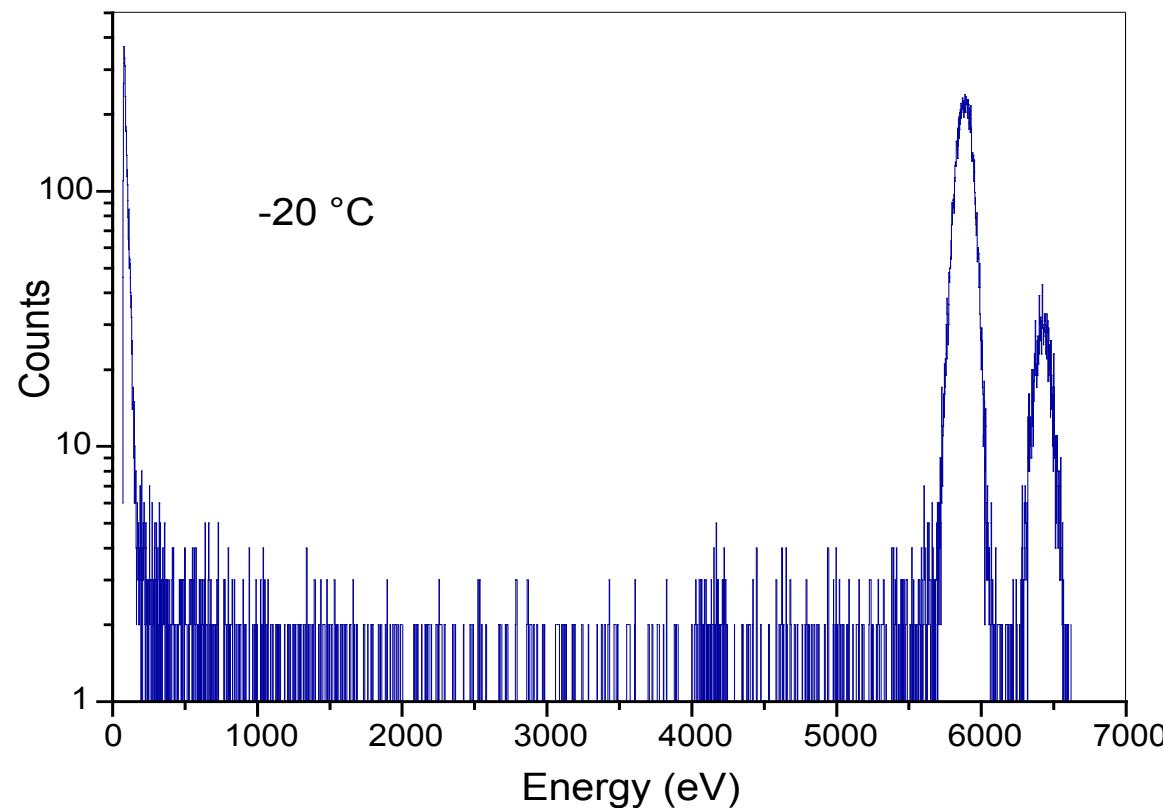
Figure 1, Left: Schematic view of a circular DEPFET pixel. X-ray photons enter through the p back contact. Signal charges are stored close to the opposite (top) side. Right: A 3×3 pixel circuitry visualizes how DEPFET pixels are interconnected in an APS matrix.

II. SDD XEUS / IXO/ATHENA



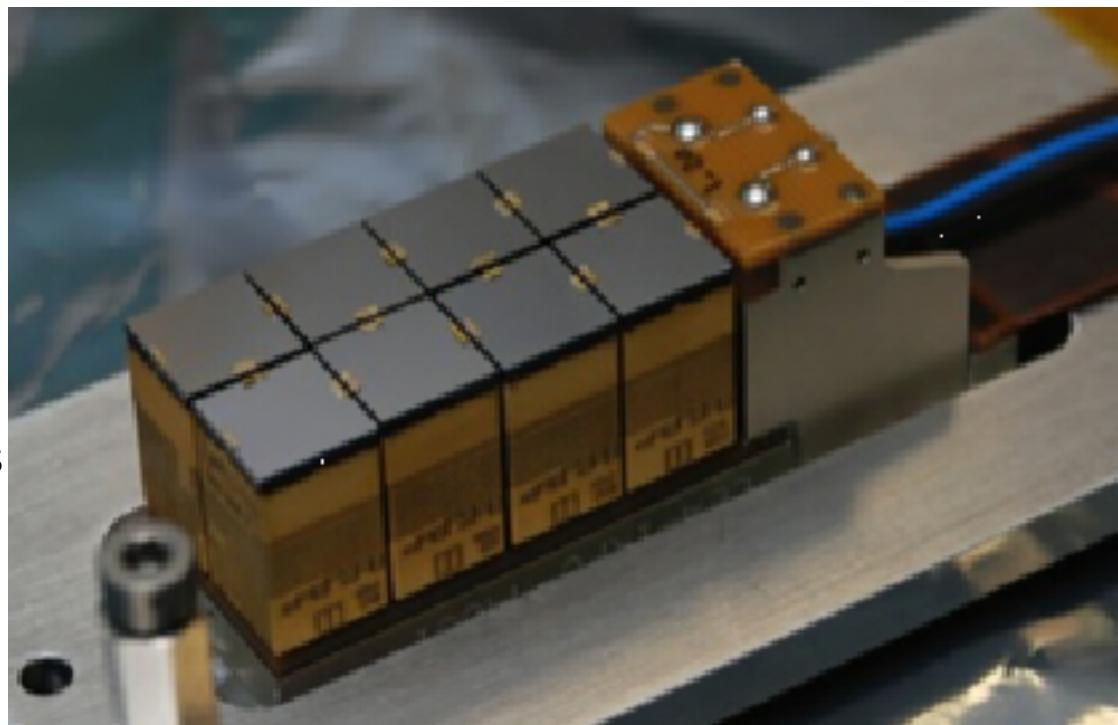
II. SSD Matrix with integrated DEPFet

- Active Pixel Sensor Matrix Low Power Consumption
- working @ room Temperature
- Filling Factor ~100 %
- Fast Readout (4 μ s/pixel, @ ms / matrice)
- NO transfert -> Robustness under irradiation
- Silicon Spectral Resolution



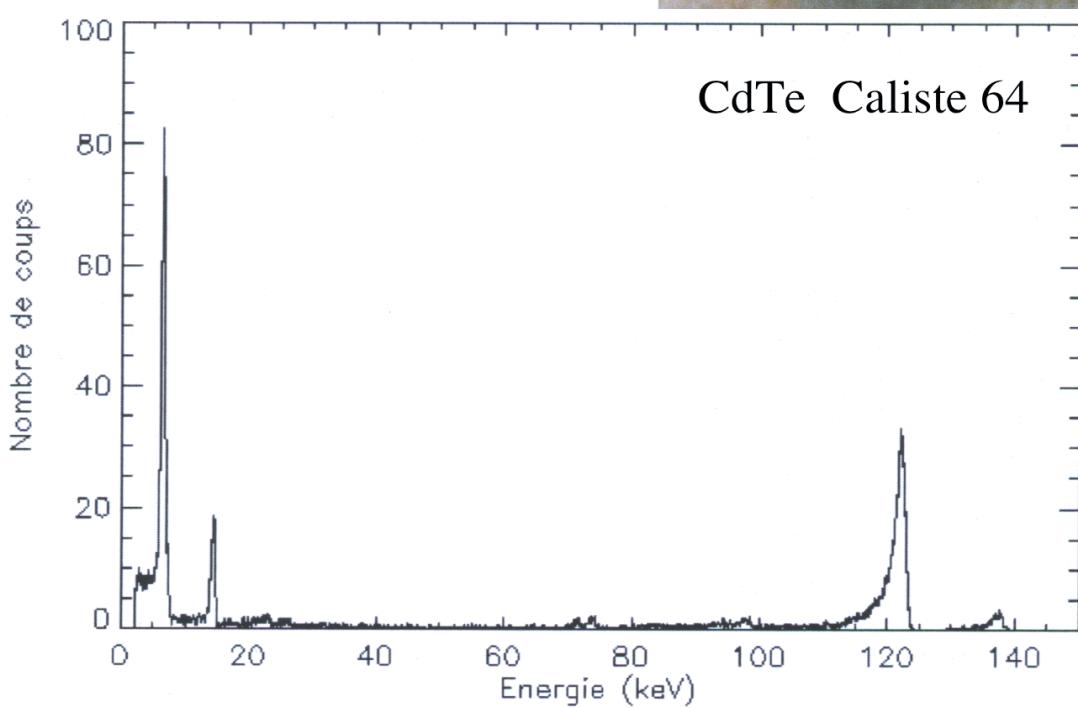
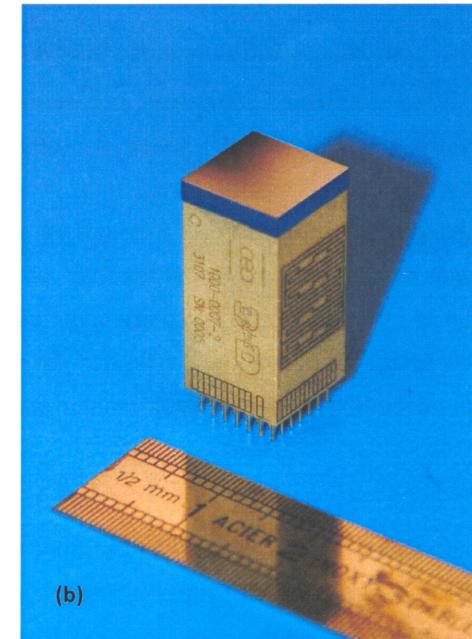
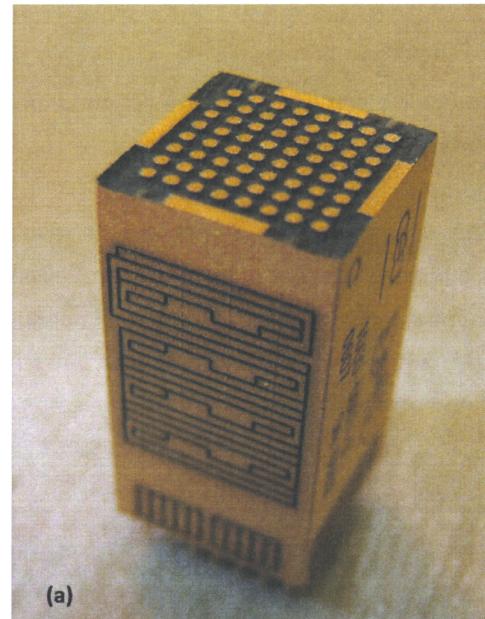
III. CdTe

- CdTe & CdZnTe have a larger gap & higher Z than Si
- Thus work @300K and absorb X-rays up to higher E



MACSi: 2048 pixels
8 HD-modules
With 1cm^2 CdTe

III. Cd Te Caliste 64



Spectre (Co57), résolution spectrale: 820 eV à 14.4 keV

(b) Détecteur intégré au module électrique

**Detector with
integrated electronic**

Comparison of Semi-Conductors for X-rays

	Silicon	Germanium	CdTe
Atomic Number	14	32	48,52
En. forbidden band(eV)	1.12	0.67	1.42
En. Pair Creation (eV)	3.62	2.96 (@77K)	4.42
Density	2.33	5.33	5.85
Q.E. (@100keV, 5mm)	19%	77%	99%
ρ ($\Omega \cdot \text{cm}$)	$2.3 \cdot 10^5$	47	10^9
Spatial Inst.	EPIC(XMM)	SPI(INTEGRAL)	ISGRI(INTEGRAL)
e ⁻ Mobility($\text{cm}^2 \text{V}^{-1} \text{s}$)	1400	3900	1350
Hole mobility($\text{cm}^2 \text{V}^{-1} \text{s}$)	1900	1900	120
e ⁻ Lifetime	>1ms	>1ms	1μs
Hole Lifetime	1ms	2ms	0.05 μs

IV. Gratings

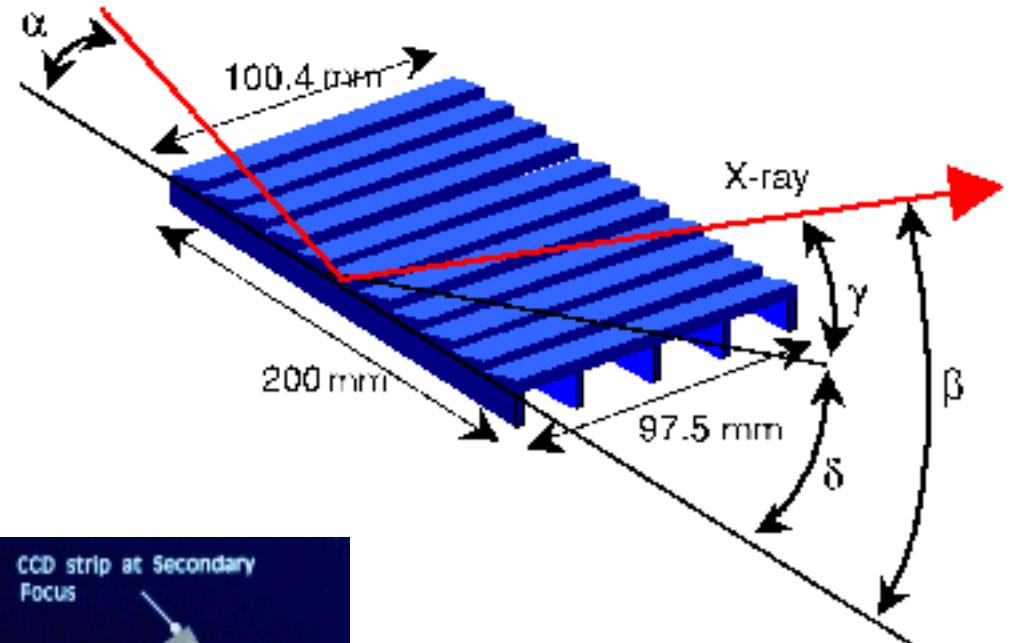
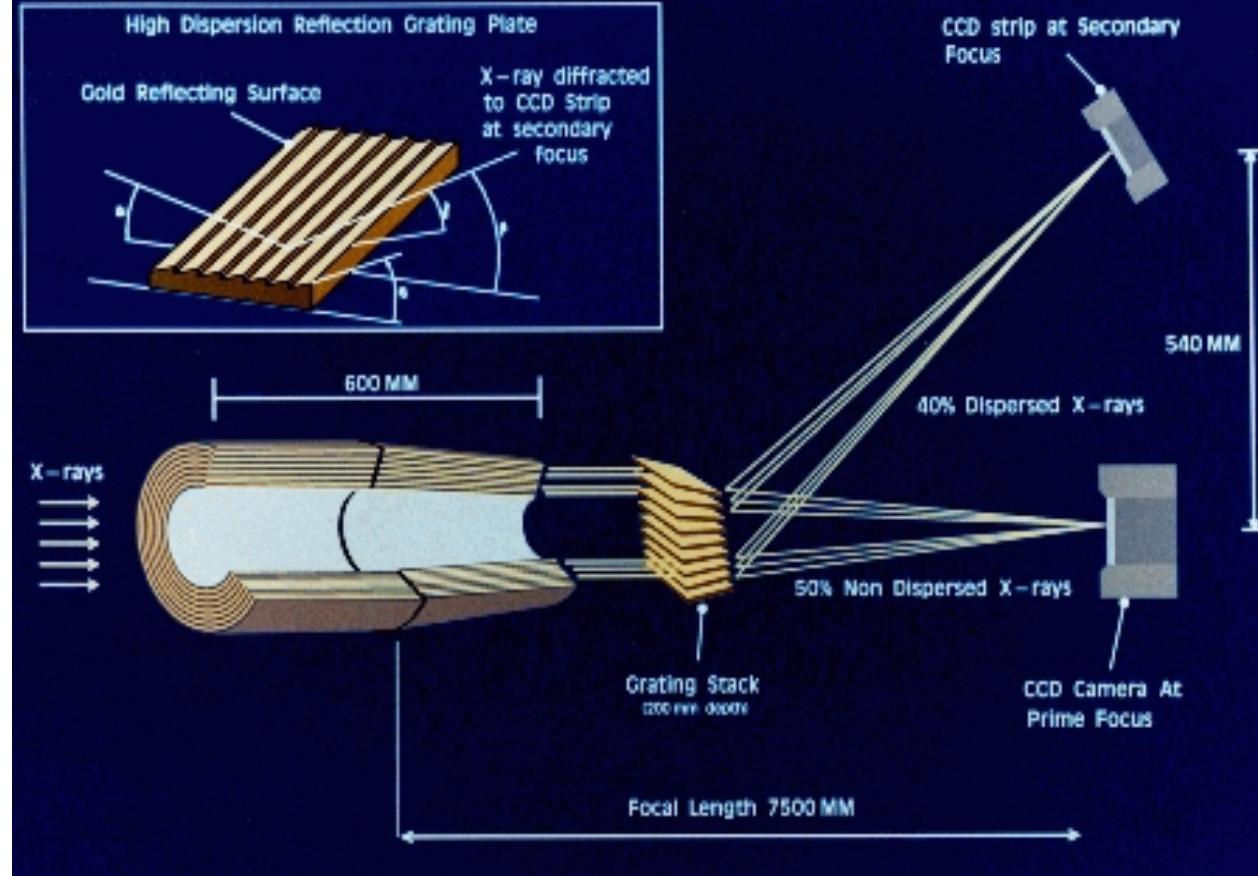
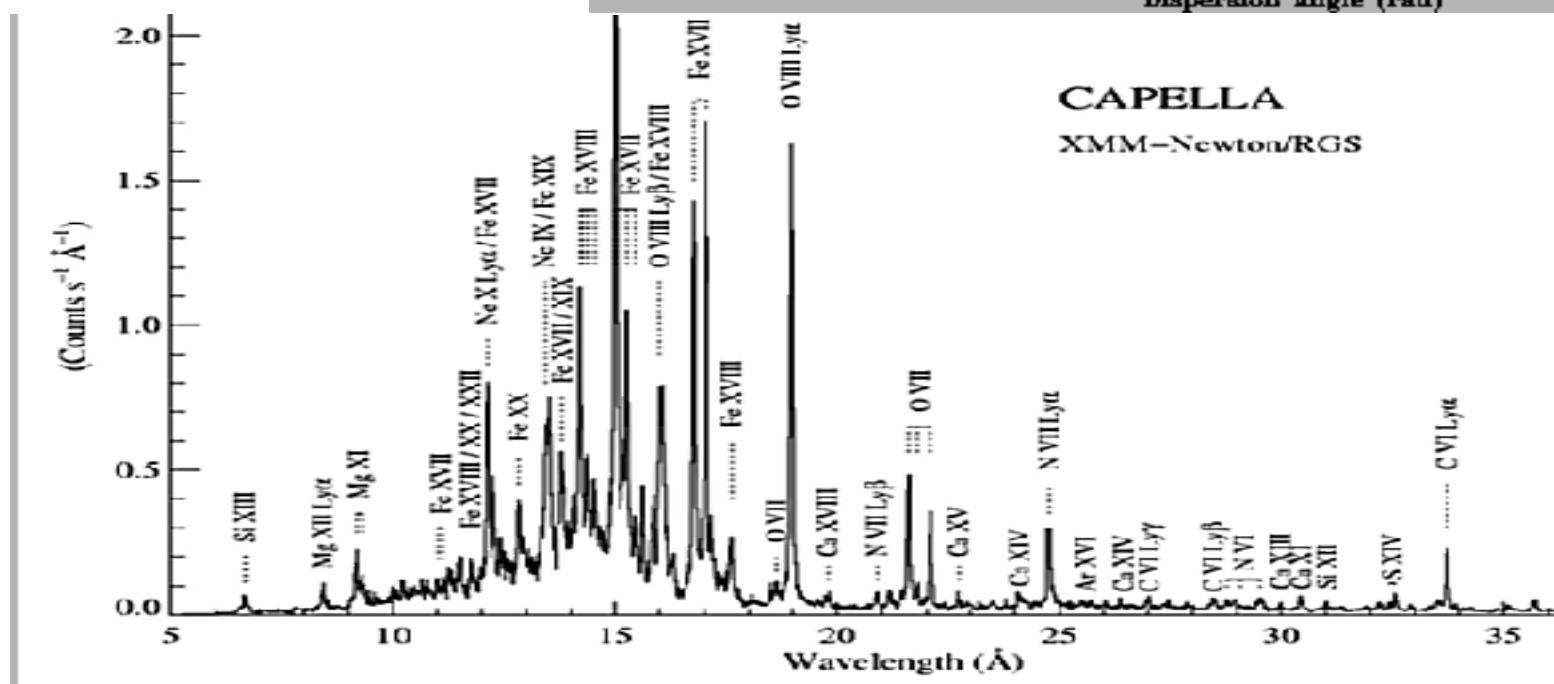
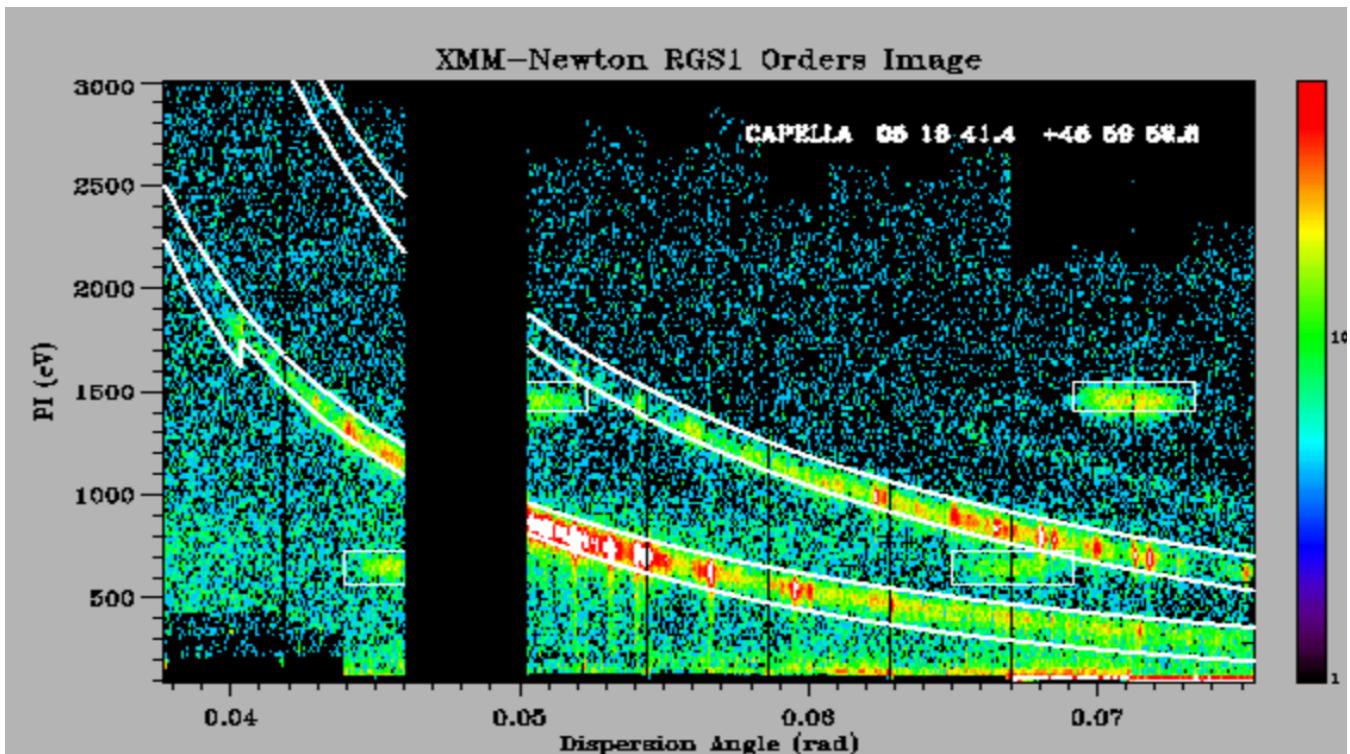


Table 8: RGS In-orbit Performance

		RGS1			RGS2						
		10 Å	15 Å	35 Å	10 Å	15 Å	35 Å				
Effective area (cm ²)	1st order	51	61	21	53	68	25				
	2nd order	29	15	-	31	19	-				
Resolution (km s ⁻¹)	1st order	1700	1200	600	1900	1400	700				
	2nd order	1000	700	-	1200	800	-				
Wavelength range	1st order	5 - 38 Å (0.35 - 2.5 keV)									
	2nd order	5 - 20 Å (0.62 - 2.5 keV)									
Wavelength accuracy	1st order	$\pm 6 \text{ m}\text{\AA}$		$\pm 6 \text{ m}\text{\AA}$		$\pm 6 \text{ m}\text{\AA}$					
	2nd order	$\pm 5 \text{ m}\text{\AA}$		$\pm 4 \text{ m}\text{\AA}$		$\pm 4 \text{ m}\text{\AA}$					
Bin size [3x3 (27 μ) ² pixels]	2.5 arcsec (cross dispersion direction)										
	7 - 14 m \AA (dispersion direction, first order)										

IV. Gratings



Semi-conductors

Energy Resolution Limitation



Some energy coming from the X-ray is degraded in heat which is NOT measured in ionisation detector such as CCDs

Since Heat formation is inescapable

⇒ Let's all the X-ray energy convert in Heat and measure it

Thus build a micro-calorimeter ...

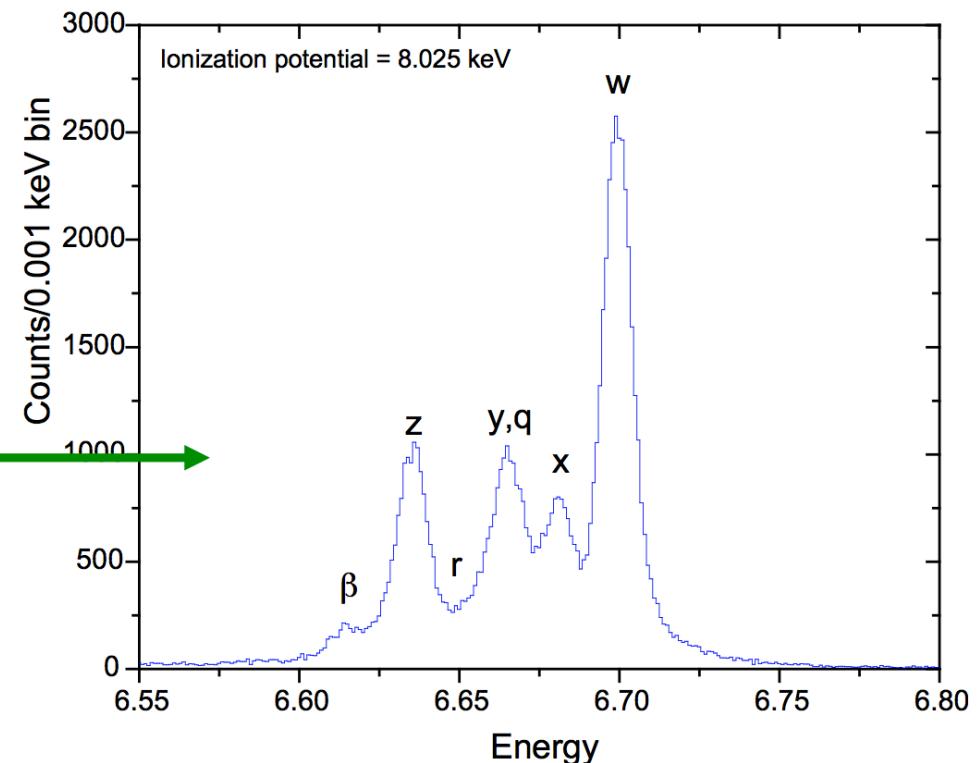
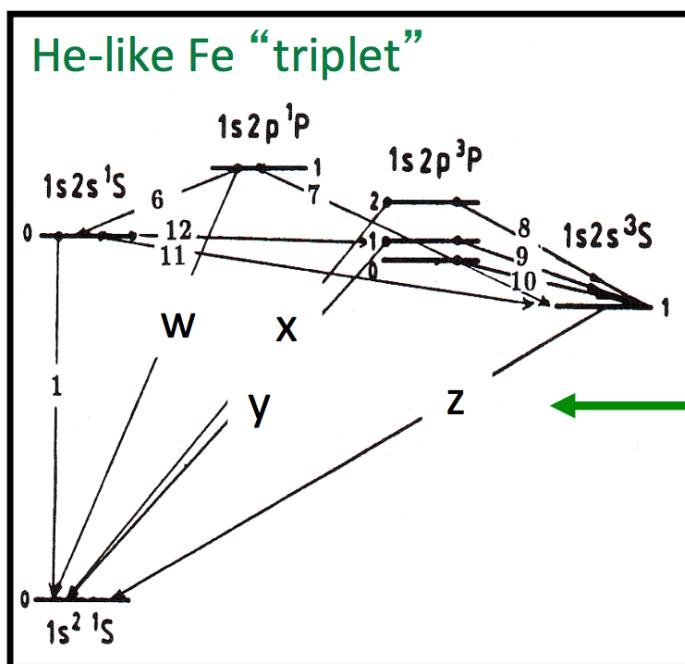
What to learn with very high spectral resolution?

- Emission line ratios provide density & temperature diagnostics
- Emission and absorption line energies identify ions
- Line shapes can be used to study velocities
 - e.g. within galaxy clusters or environment around a supermassive black hole

R: Resonance line (w)

I: Intercombination line (x+y)

F: Forbidden line (z)



V. μ -Calorimeters

- Principle
- Sensors
 - Transition Edge Sensors
 - Metal-Insulator Sensors
- Readout Electronic
- Conclusion

μ -Calorimeters Principle (1)

Two definitions:

- μ -Calorimeters measure Energy pulse issued from the interaction of a high energy photon with an absorber.
- μ -bolometers are designed to be sensitive to an Energy flux as in sub-mm Astronomy

Calorimeters Basis:

-I- The Heat capacity decrease to small value when temperature decrease

$$C = \alpha T + \beta(T/\theta_D)^3 \quad (\theta_D : \text{Debye temperature})$$

Small deposit of energy will enhance the temperature of the system:

$$\Delta T = E_x / mC$$

-II- There are sensors able to measure such temperature @ very low T.

μ -Calorimeters Principle (2)

Main Elements of a μ -Calorimeter

An **Absorber** of the x-ray photons linked through a **thermal impedance** to a **sensor** (thermometer). This sensor is itself linked to a cold bath through **an another thermal impedance**.

The thermal impedance Sensor-Cold Bath should be adapted to have enough time to measure the temperature elevation but not too much so that the pixel is ready for a new interaction.

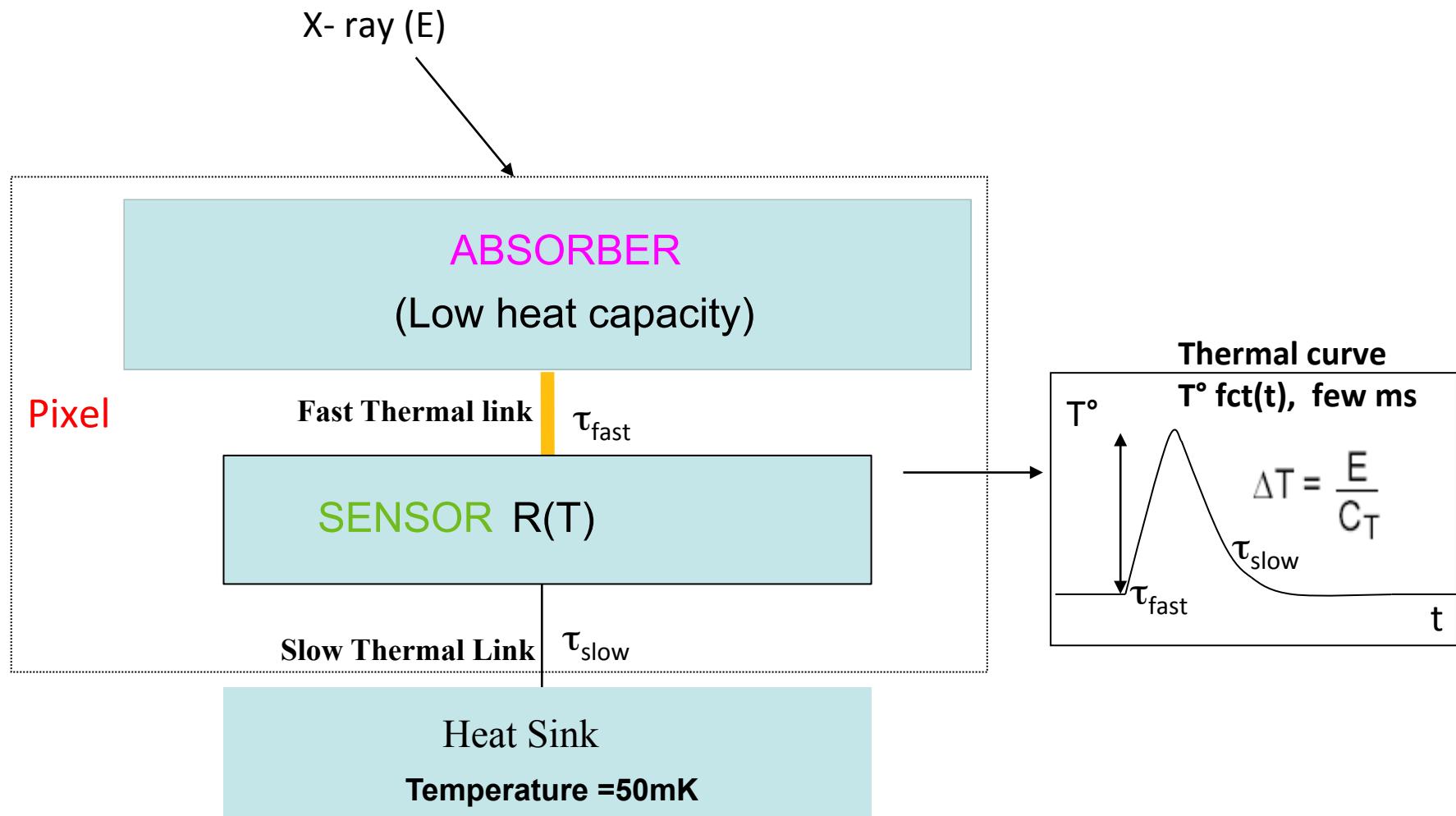
The signal will be a thermal curve vs Time.

From the curve parameter, one can deduce the X-ray energy.

Since we measure a temperature, (i.e. a phonon bath in thermal equilibrium), there is no partition in energy (like e^- -hole pair and heat a in semi-conductor)...

Thus no Fano Factor Limitation

μ -Calorimeters Principle (3)



Rule of Thumb: $\tau_{\text{slow}} \geq 10 * \tau_{\text{fast}}$

μ -Calorimeters Principle (4)

$$\Delta T = E_x / M C(T)$$

E_x : Energy of incoming photon

M : Pixel Mass

$C(T)$: heat Capacity @ T

ΔT : Temperature Elevation

$$C = \alpha T + \beta (T/\theta_{\text{Debye}})^3$$

e^- contribution

Phonon contribution

To have good μ -Cal \Rightarrow Enhance ΔT

thus minimize M : smaller pixels

for given X-ray : minimize C & T

for a given T : minimize α & β

To kill α \Rightarrow choose material with NO free e^-

To minimize phonon Contrib., maximize θ_{debye}

Minimizing M \Rightarrow smaller pixel size (what about FOV or complexity ...)

Minimizing T \Rightarrow Good but STRONG Spatial CryoCooler Limitations (1 μ W@50mK)

Minimizing C \Rightarrow efficiency will depend on the pixel technology & choices

Signal to Noise Ratio

C. Enss and D. McCammon

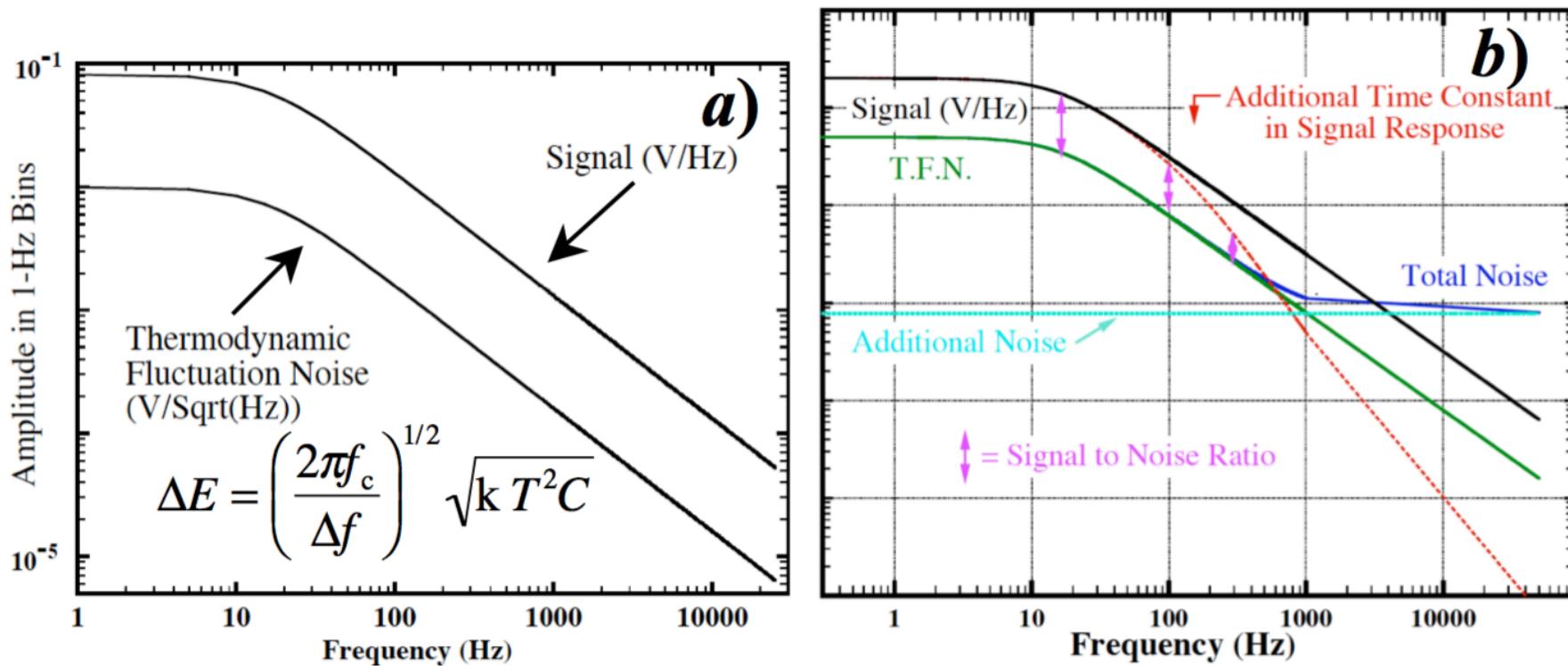


Fig. 4. *a*) Constant ratio of signal to TFN allows signal amplitude to be measured with arbitrarily high accuracy as the bandwidth Δf goes to infinity. *b*) Usual limitation is that either an additional noise source or an additional pole in the signal response causes the signal to total noise ratio to drop above some frequency.

μ -Calorimeters Absorber (1)

Good Qualities for X-ray Absorber:

- High Stopping Power \Rightarrow High Z material
- Heat Capacity as low as possible
- Fast Thermalization

Those 3 properties are more or less antonymous

μ -Calorimeters Absorber (2)

Resistive Material (Au Bi,...):

Very Good Thermalization (no Energy Trapping)

But Heat Capacity $\propto T \Rightarrow$ Dominant Component @ Low T

$$C = \alpha T + \beta (T/\theta_{\text{Debye}})^3$$

Superconducting Materials (Sn, Ta):

Heat Capacity depends only of phonon Contribution when $T \ll T_c$. (very small).

X-ray break Cooper Pairs, producing quasi-particles that trap Energy.

If these QPs are slow to recombine \Rightarrow Slow Thermalization

Insulators & Large Gap SemiConductors (Si, Ge):

Heat Capacity depend only of phonon Contribution & θ_{Debye} is High (\Rightarrow Very Small Heat Capacity) but X-ray produce e^- -Hole pairs that trap Energy for a while in impurities or defects thus Slow Thermalization

Small Gap SemiConductors: HgTe HgZnTe) :

Small Gap \Rightarrow Good Thermalization but θ_{Debye} is low thus Heat Capacity not so small

Sensors (1)

There are 3 main sensor types:

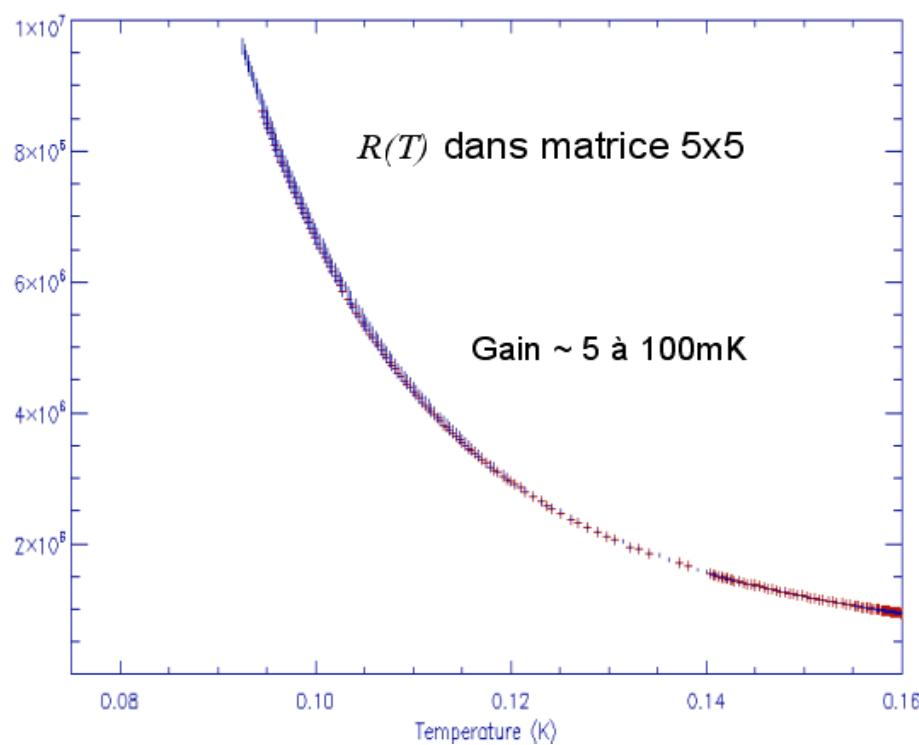
- **TES** (Transition Edge Sensor)
- **MIS** (Metal-Insulator Sensor)
- MMC (Metallic Magnetic Calorimeter)

α is the log derivative of R/T

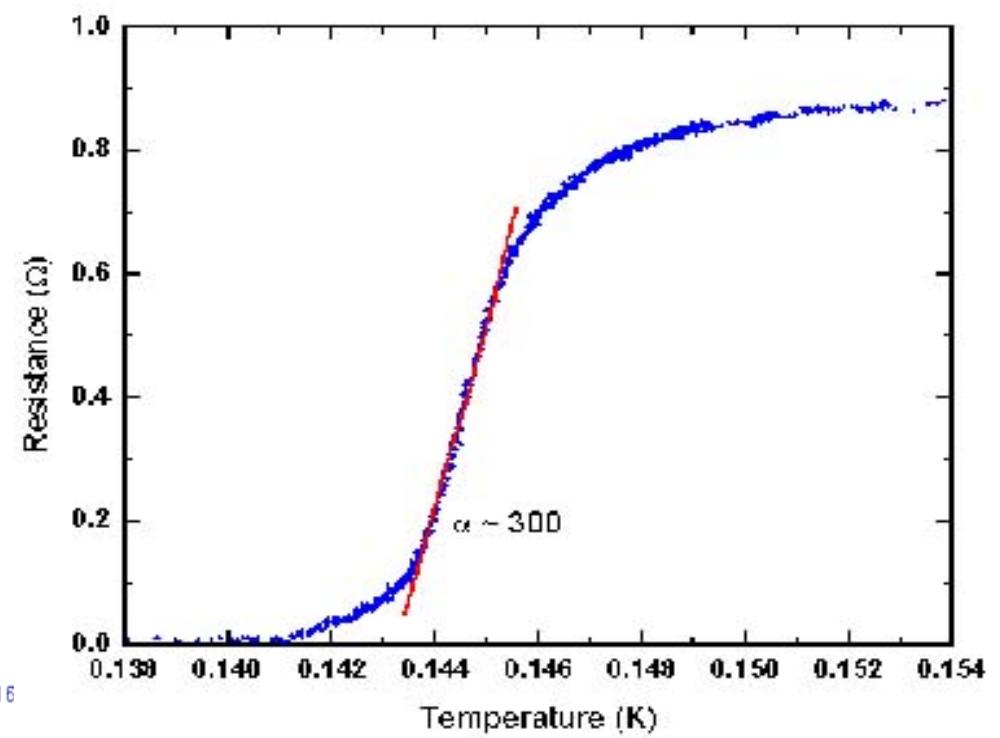
$$\alpha = d(\log R)/d(\log T) = T/R \frac{dR}{dT}$$

Sensors (2)

MIS



TES



Sensors (3)

- Sensors polarised in order to create an electro-thermal feedback:
- MIS polarized in current \Rightarrow Energy dissipation ($P=R I^2$) decreases when T increase (since R decrease)
- TES polarized in voltage \Rightarrow Energy dissipation ($P=U^2/R$) decreases when T increase (since R increase)
- In both cases, the Electro-Thermal Feedback speed up the return to thermal equilibrium and thus the energy dynamic.

Sensors (4)

MIS Sensors follow $R=R_0 e^{\nu(T_0/T)}$ Efros & Shklovskii Law

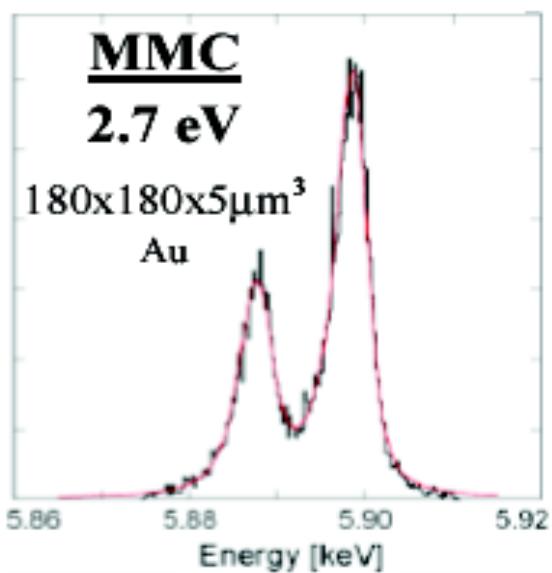
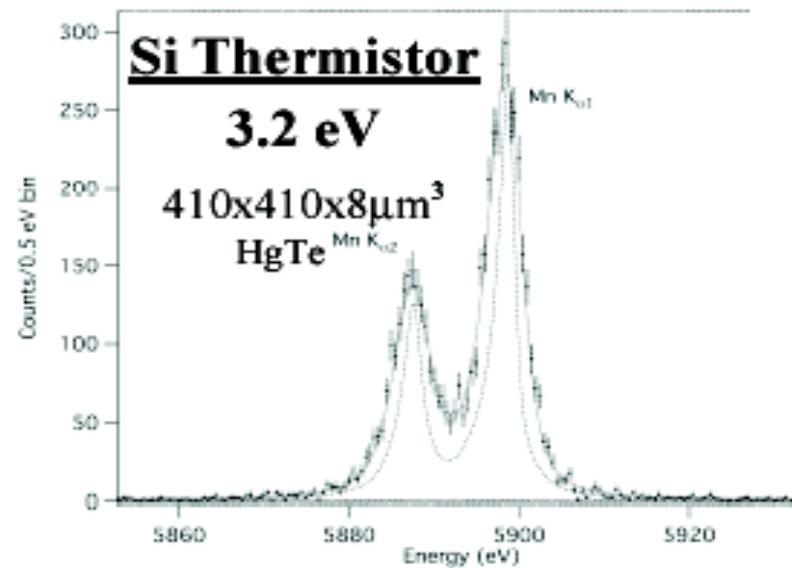
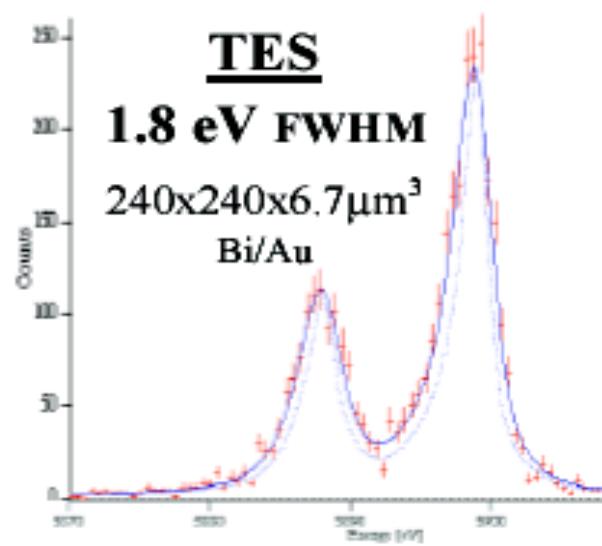
T_0 & R_0 are doping dependant. Small $\alpha \approx 5-10$ & $R \approx 1 - 10 \text{ M}\Omega$

R quite high \Rightarrow Relatively slow sensors

But MIS are sensitive over a large range of E (even if not strictly linear)

TES : Superconductor goes from ZERO ohm to normal resistor ($R \approx 1 \text{ m}\Omega$) in a mK interval thus sensitivity is excellent ($\alpha \geq 100$) \Rightarrow One can use metallic absorber with excellent thermalization & high stopping power (Gold). Fast system but suffer saturation effect.

Spectral Resolution



$$\Delta E_{FWHM} = 2.35\xi\sqrt{KT^2C}$$

Spécifications
ATHENA+ :
2eV @ 6 keV

Readout Electronics

- Need of cryo-electronics

TES : $m\Omega$ impedance \Rightarrow very weak current

MIS : $M\Omega$ signal \Rightarrow Short Connections

- Both have to minimize the thermal load onto the coldest part (50mK)
- If Large Matrices \Rightarrow Multiplexing is Mandatory

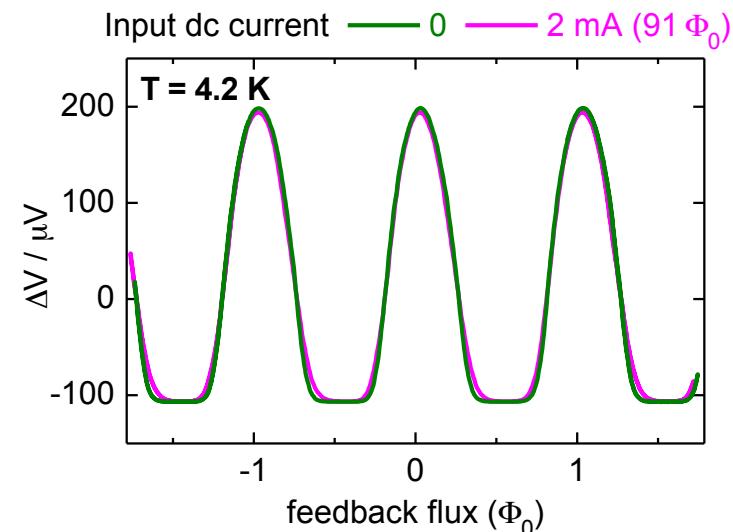
What kind of Electronics could be used ?

TES Readout Electronics (1)

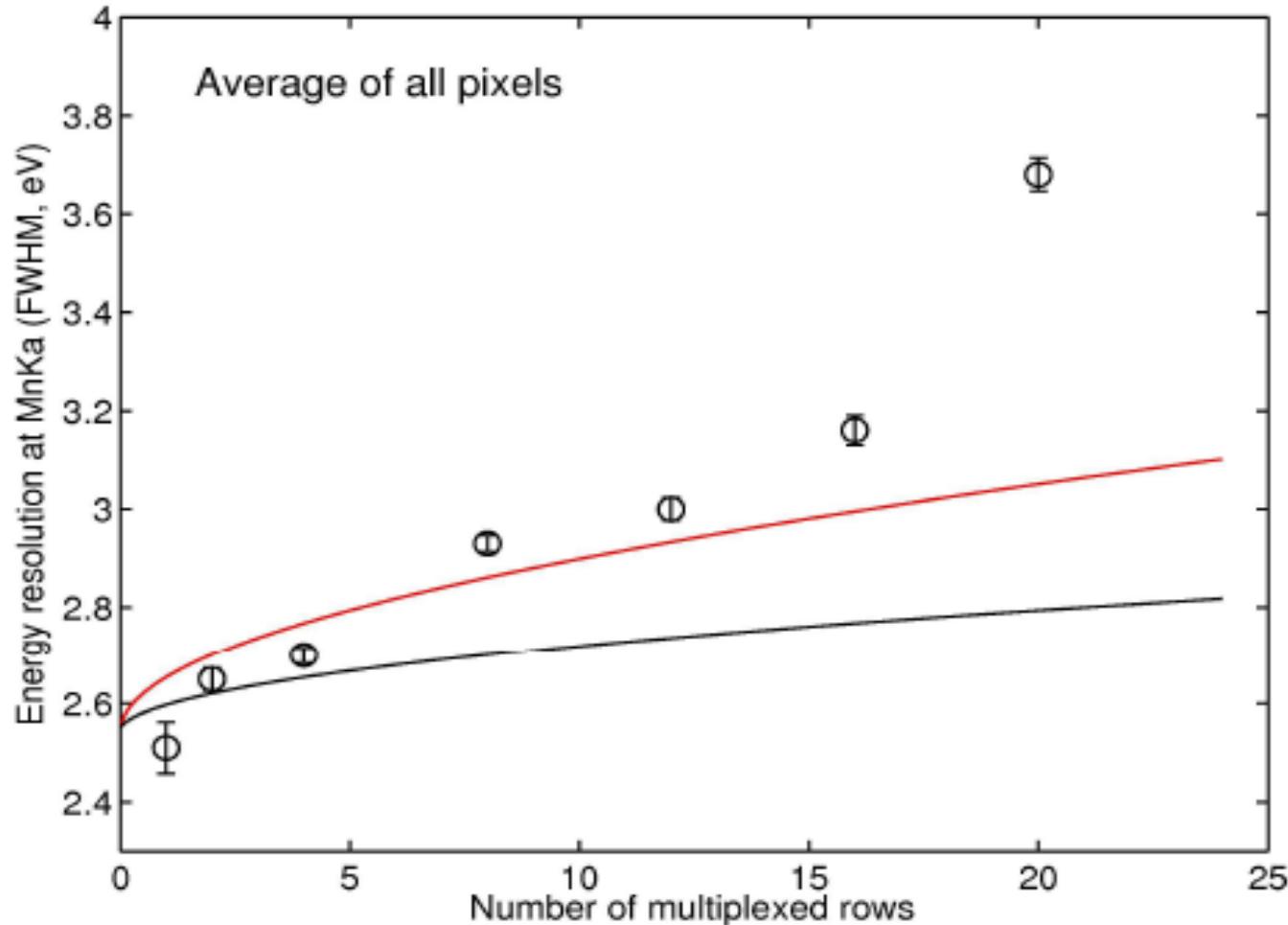
- TES : $m\Omega$ impedance \Rightarrow very weak current
 \Rightarrow SQUID @ the lowest temperature.

OK for one detector but in Large Matrix, SQUID power consumption will become the dominant contribution.

Due to cyclic response, SQUID are **very difficult to multiplex**



TES Readout Electronics (2)



Energy resolution degradation with the number of multiplexed pixels

GSFC TES arrays

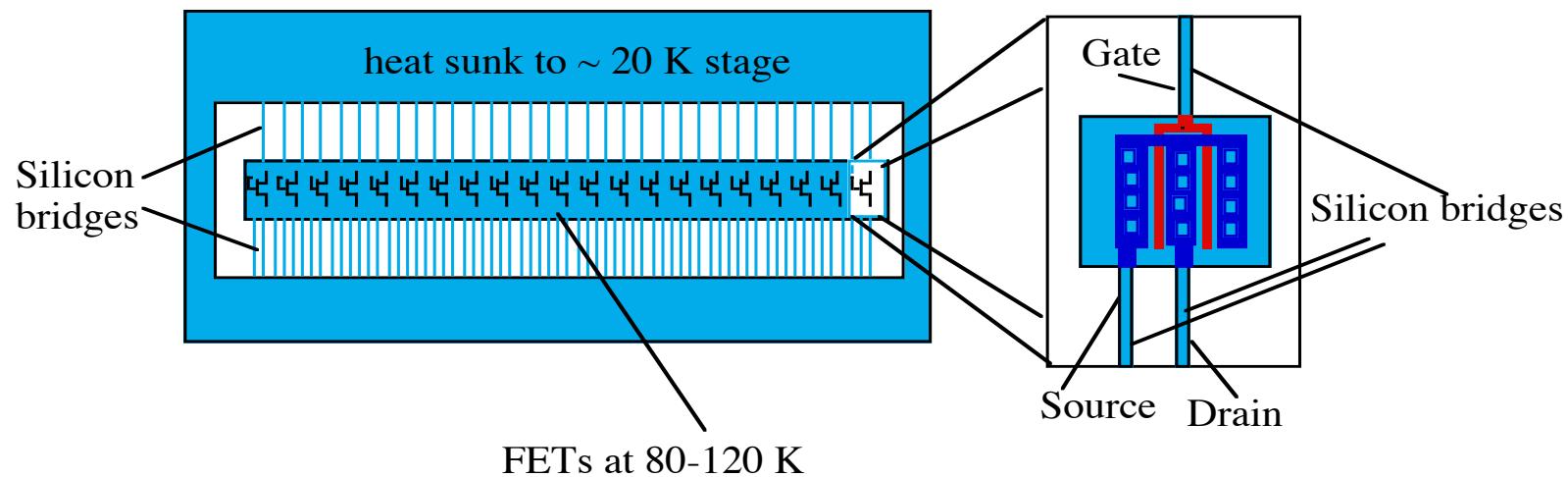
MIS Readout Electronics (1)

Astro-H/ITOMI MIS Matrices (GSFC)

60 μ W/ JFET. JFETs @ 20 K.

JFETs have optimal noise @130 K \Rightarrow Self-Heating !

Single JFET



Si-MOSFET OK for very high impedances : $100 \text{ G}\Omega$

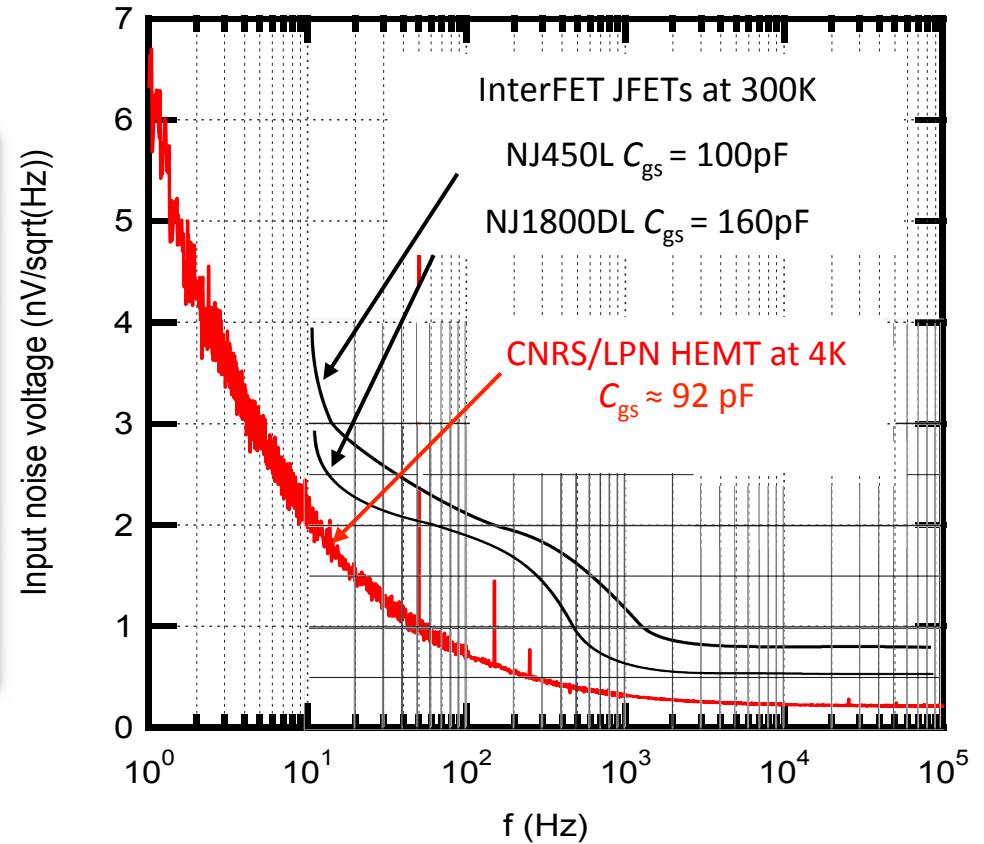
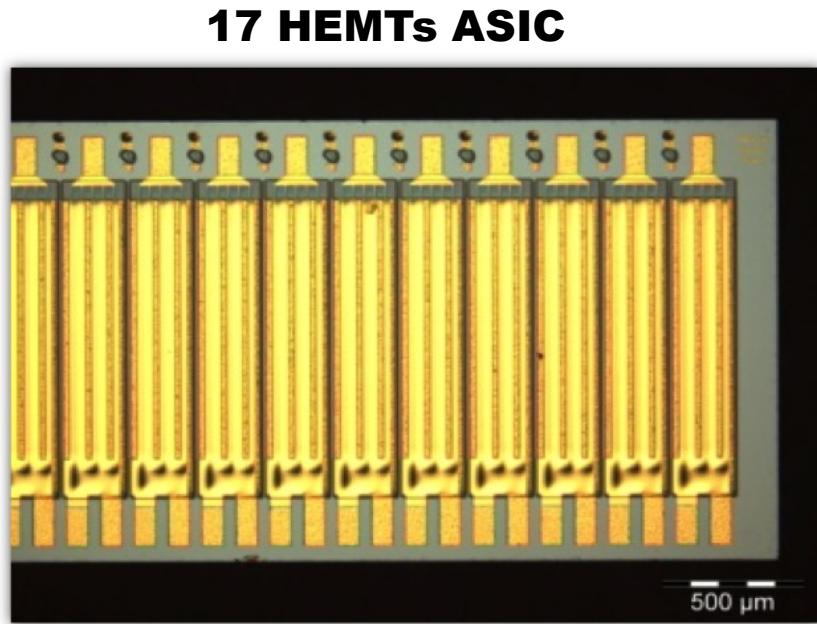
For intermediate sensor impedances ($100 \text{ k}\Omega - 100 \text{ M}\Omega$) :

GaAs/GaAlAs HEMT possible (work @ very low T)

Bipolar Electronic SiGe @2- 4 K for Low($1\text{k}\Omega-10\text{k}\Omega$) impedances

MIS Readout Electronics (2)

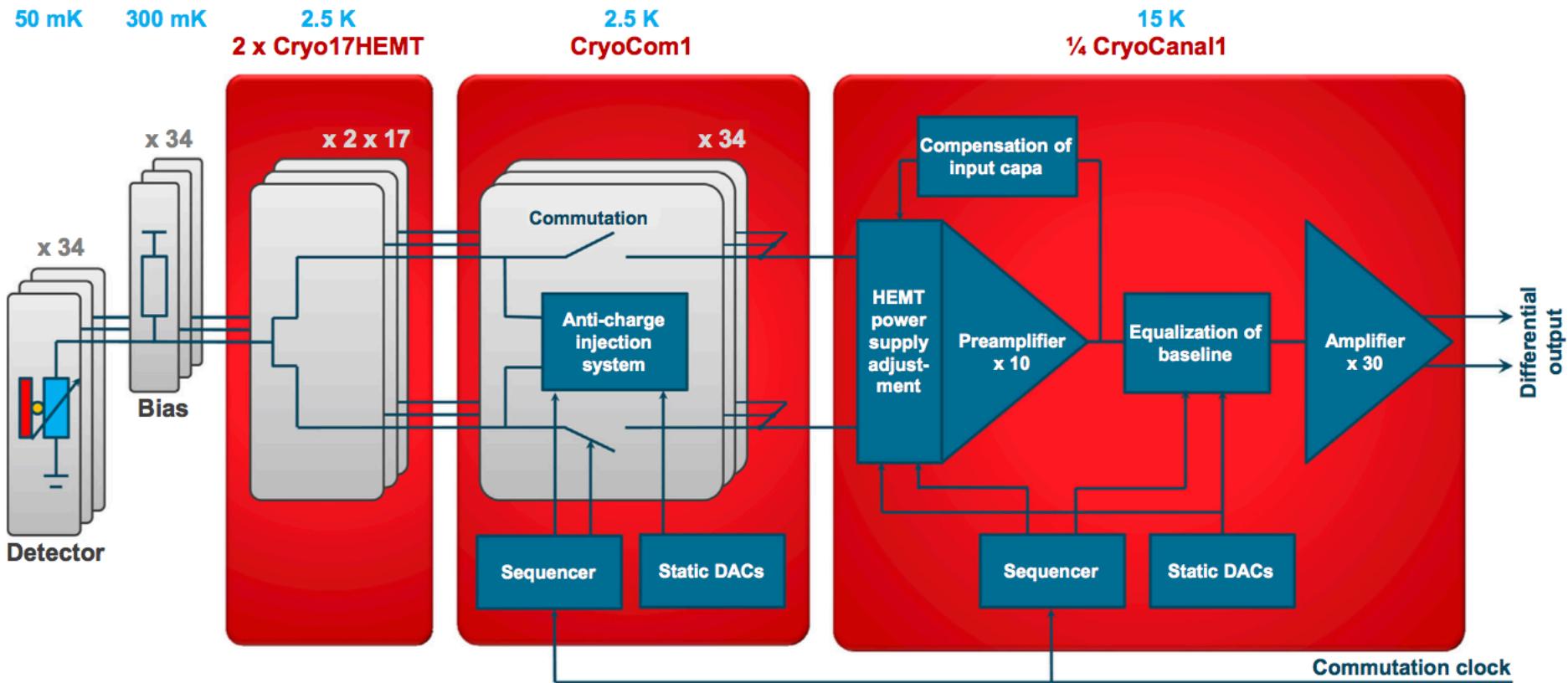
GaAs/GaAlAs HEMTs
LPN (Laboratoire de Photonique et nanostructures)



Noise voltage: comparison with Si JFET

HEMT with $C_{gs} = 92 \text{ pF}$

MIS Readout Electronics (3)

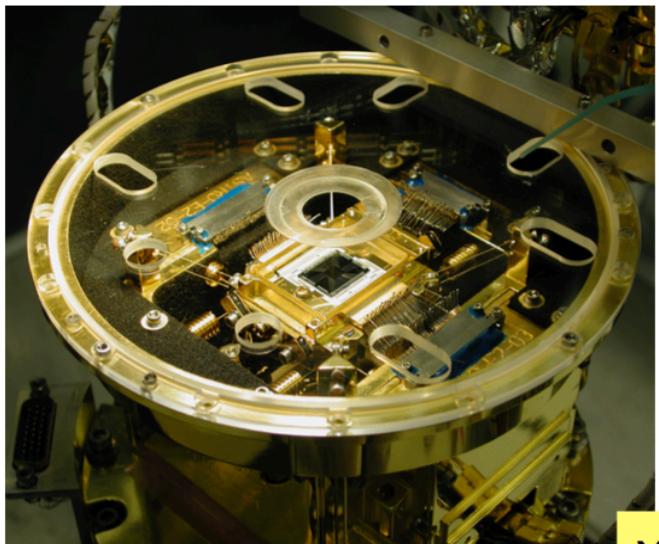


Complete schematic of MIS large array Cryo-Electronics
With HEMTs, MUX and Amplification

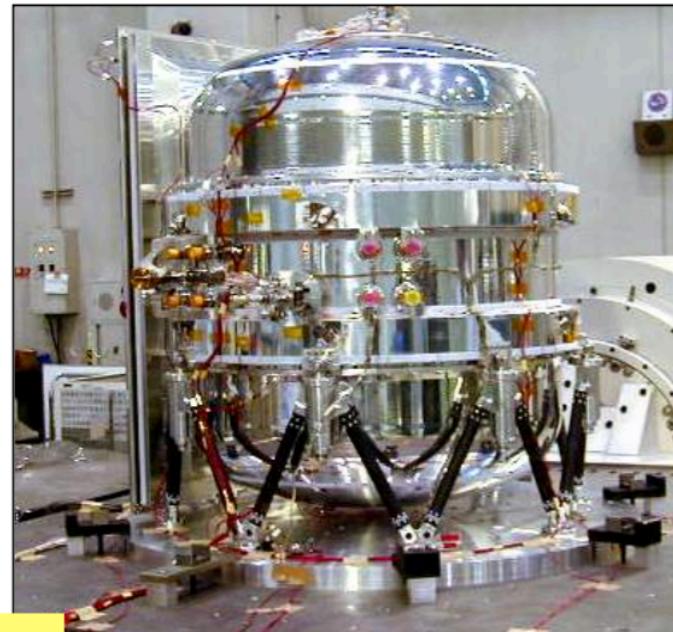
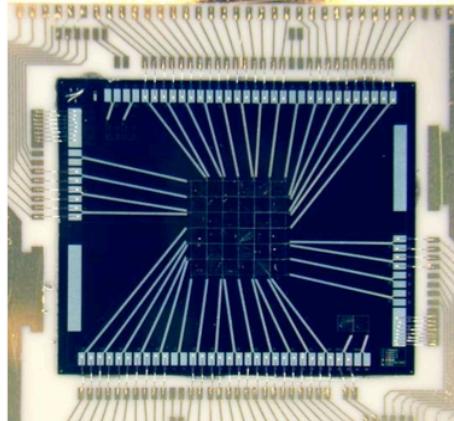
TES/MIS Comparizon

Red=Difficult Blue=Easy	TES	MIS
Fabrication	Few Steps(~10)	Difficult (~120 steps)
Heat Capacity per Pixel	0.8 pJ/K	0.04 pJ/K
Sensitivity(sensor)	$\alpha=75$ (1/F noise)	$\alpha= -5 -10$
Speed	25 cp/s (goal)	~3-4 cp/s
Saturation in E	Yes (Limited by ETF & stiffness of transition)	No ; $R=R_0 \exp(T_0/T)^{0.5}$
Magn. Field Sensitivity	Very sensitive TES & SQUIDs	NO
Sensor Impedance	$m\Omega$	$M\Omega$
Cryogenic Electronic	SQUIDS (difficult)	HEMT/SiGe (~Classical)
MUX	VERY Difficult	~classical
Filling Factor	Small (Deposited Absorber)	Very good (sticked Absorber)
P (4000p+Readout) @ 50mK	SQUIDS $\Rightarrow >1\mu W$	Very Low @50mK

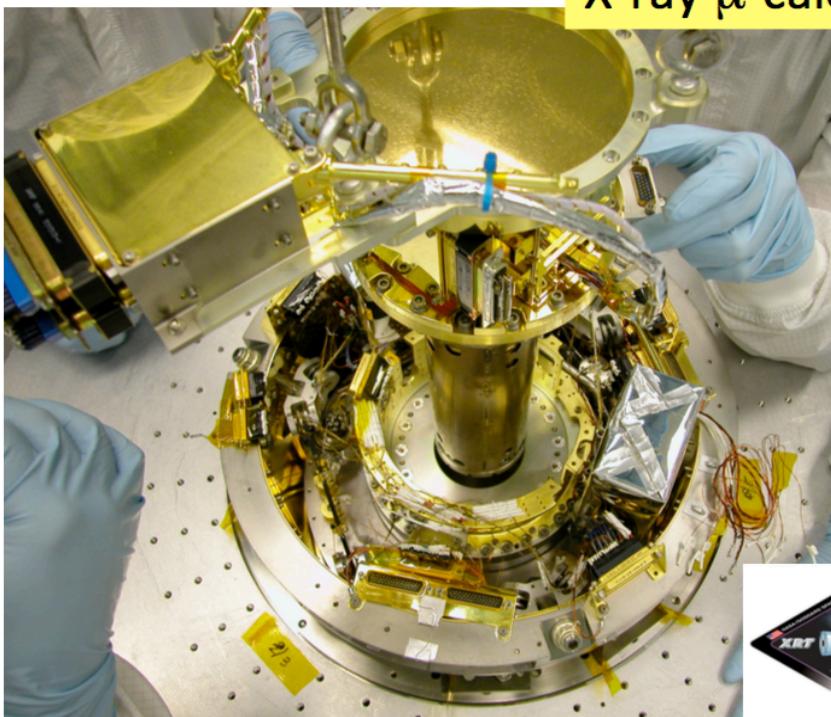
Itomi: MIS State of the art



SXS 6x6 matrix

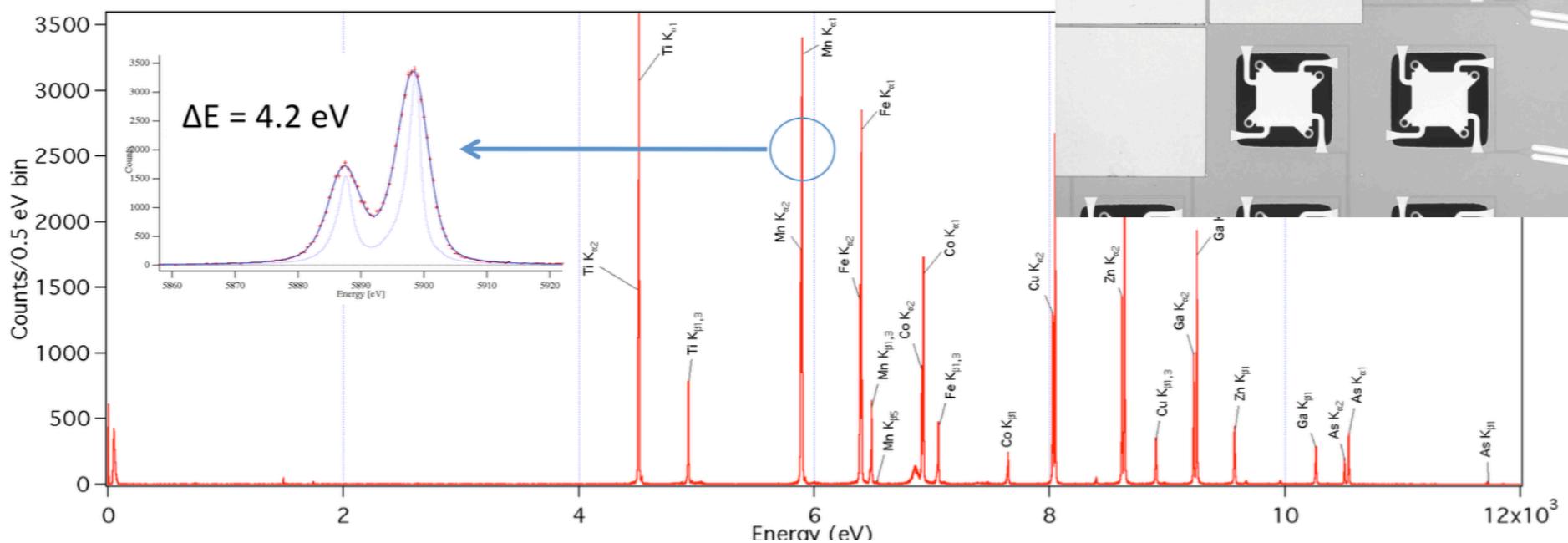
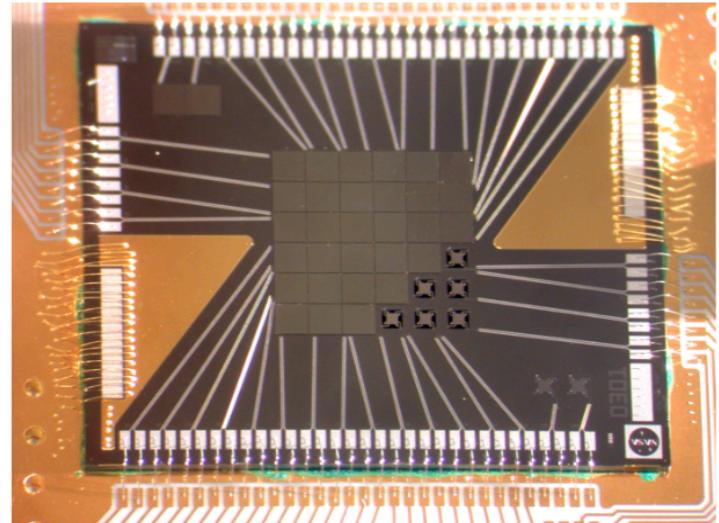


X-ray μ -calorimeter array



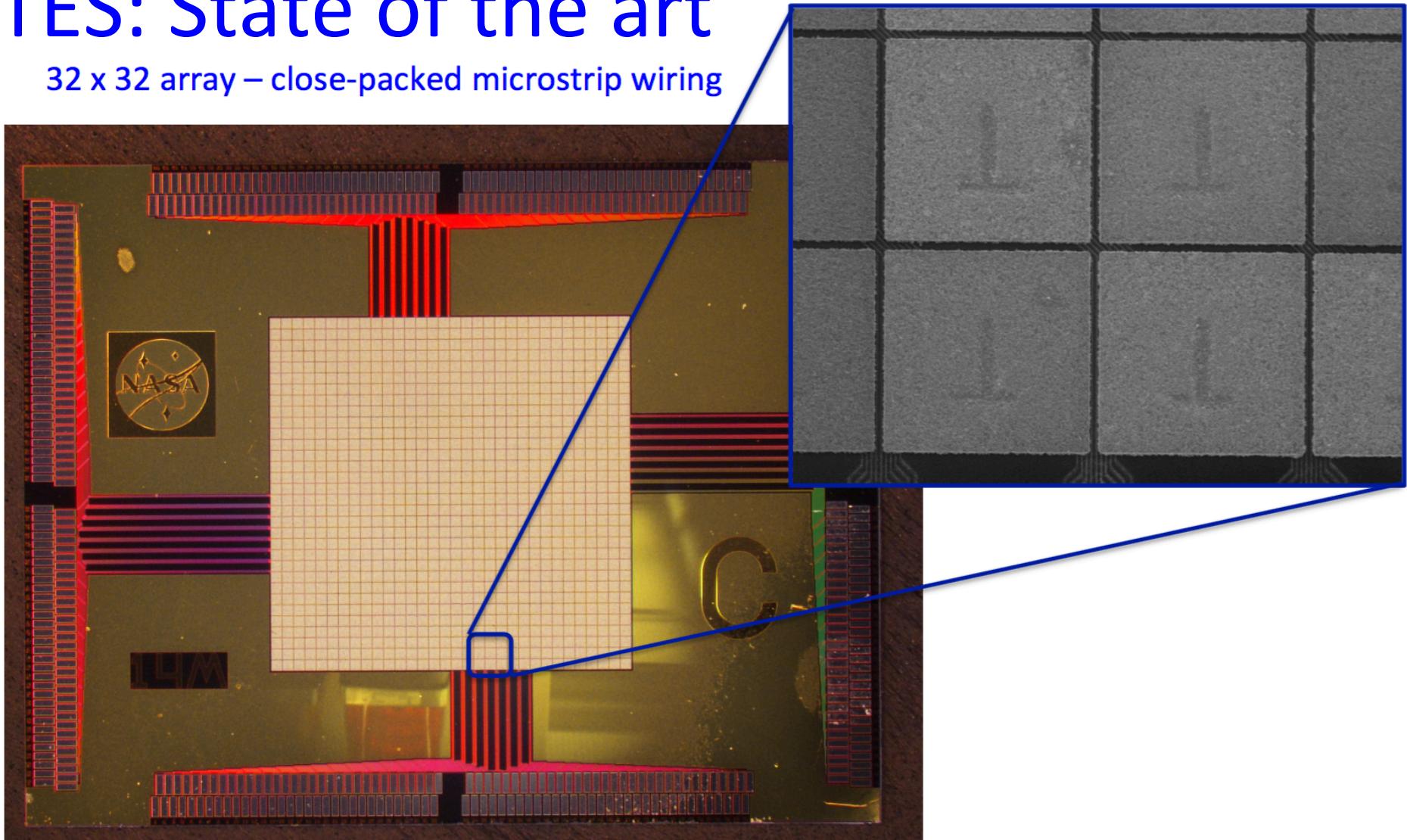
Itomi: MIS State of the art

- 4 eV resolution at 6 keV.
- 6 x 6 array of $820 \mu\text{m}$ pixels (30" pixels, 3' FOV)
- Ion-implanted thermistors; manually attached absorbers
- 5.6 m focal length, HPD better than 1.7'
- Technology (eg. JFET amplifiers) limits number of pixels & count rate capability



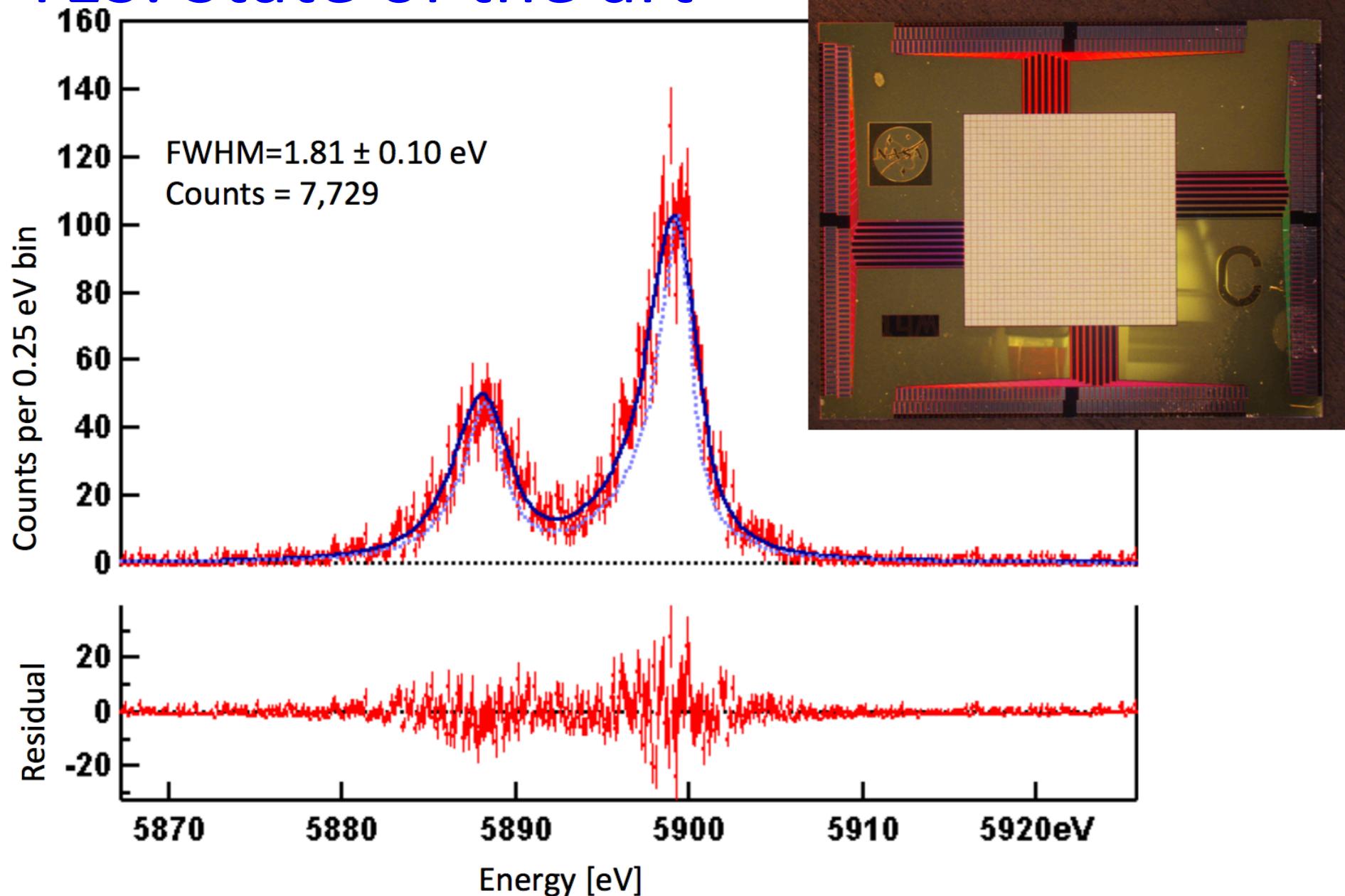
TES: State of the art

32 x 32 array – close-packed microstrip wiring

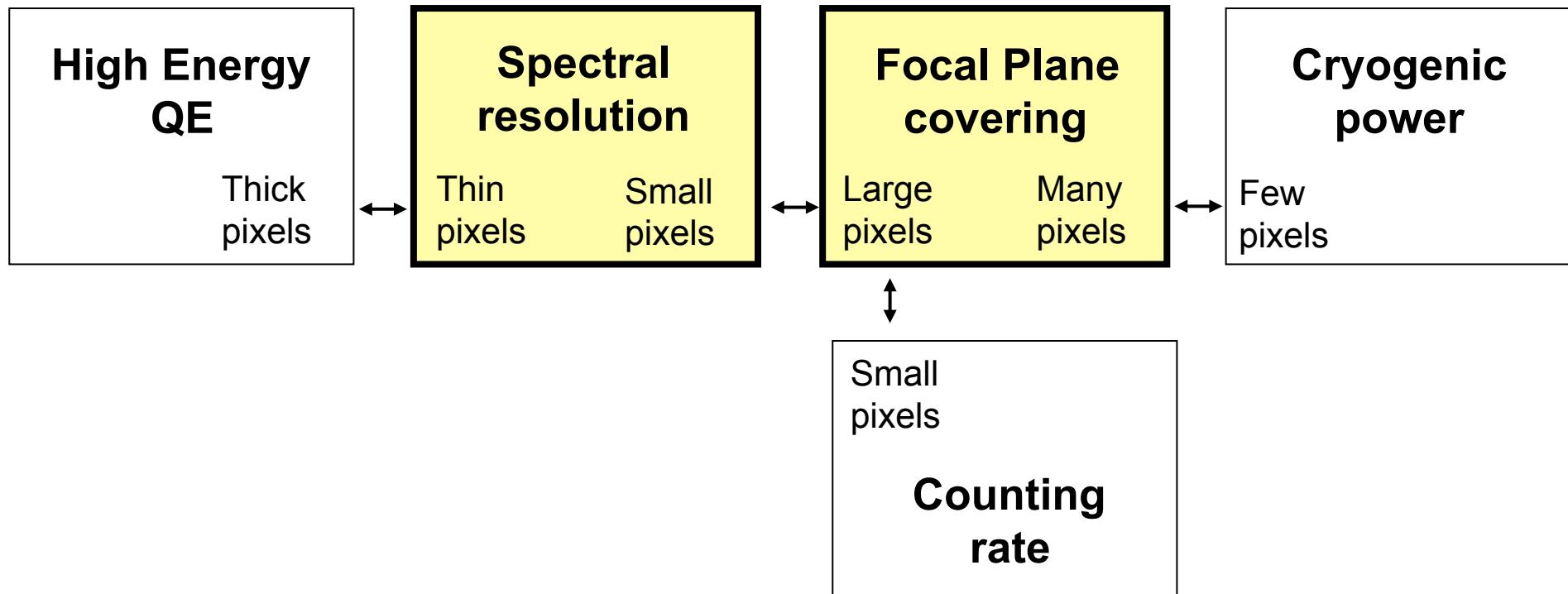


- Absorbers:
- 242 x 242 μm on 250 μm pitch
 - Au: 1.5 μm , Bi: 3.0 μm – electroplated
 - >90% QE at 6 keV

TES: State of the art



Microcalorimeters Trade-offs



Conclusion

- **CCDs** were the first real spectro-imagers in X-rays Astronomy (2000-2016). **DEPFets** present excellent improvement.
- **CdTe** are THE semiconductors for high energies
- **μ -Calorimeters** will be THE next generation Low-Energy instruments ... Very demanding in Spatial instrument budgets (Thermic, Power, weight...)