#### INSTITUT D'ASTROPHYSIQUE DE PARIS



PHYSICS and ASTROPHYSICS of COSMIC RAYS

OHP Saint Michel l'Observatoire, France

# EFFECTS OF REACCELERATION AND SOURCE GRAMMAGE ON SECONDARY COSMIC RAYS SPECTRA

28/11/19 - Virginia Bresci



$$-\frac{\partial}{\partial z} \left[ D_{\alpha} \frac{I_{\alpha}}{\partial z} \right] + v_{A} \frac{\partial I_{\alpha}}{\partial z} + 2h_{d} n_{d} v_{\alpha} \sigma_{\alpha} I_{\alpha} \,\delta(z) - \frac{2}{3} v_{A} A_{\alpha} p^{3} \frac{\partial F_{\alpha}}{\partial p} \delta(z) + 2h_{d} \delta(z) \frac{\partial}{\partial p} \left[ \left( \frac{dp}{dt} \right)_{\alpha, \text{Ion.}} I_{\alpha} \right] =$$

 $= A_{\alpha} p^2 Q_{\alpha}(p) \delta(z)$ 

...What experiments really measure:

$$I_{\alpha}(z, E_k)dE_k = v_{\alpha}(p)F_{\alpha}(z, p)p^2dp$$



**Spatial Diffusion** 



$$-\frac{2}{3}v_{A}A_{\alpha}p^{3}\frac{\partial F_{\alpha}}{\partial p}\delta(z) + 2h_{d}\delta(z)\frac{\partial}{\partial p}\left[\left(\frac{dp}{dt}\right)_{\alpha,\text{Ion.}}I_{\alpha}\right] =$$

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Advection

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Adiabatic expansion

Ionization of the medium

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Source term

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#### **SOURCE TERM**

#### Primary nuclei



Diffusion

Advection

$$+\frac{1}{3}\left(\frac{du}{dz}\right)p\frac{\partial f(z,p)}{\partial p}+\eta\frac{n_1u_1}{4\pi p_{\rm inj}^2}\,\delta(p-p_{\rm inj})\,\delta(z)=0$$

Compression

Injection at the shock surface



$$f_0(p) = s \frac{\eta n_1}{4\pi p_{\rm inj}^3} \left(\frac{p}{p_{\rm inj}}\right)^{-s}$$

 Power law in momentum, slope s=3r/ (r-1) depends ONLY on the compression ratio r=u1/u2 -> 4
 No dependence upon diffusion

#### Secondary CRs

Created in interactions between primaries - ISM:

Nuclei
$$A_{\alpha}p^{2}Q_{\alpha}(p) = \sum_{\alpha' > \alpha} 2h_{d}n_{d}v_{\alpha'}(E_{k})\sigma_{\alpha'\alpha}I_{\alpha',0}(E_{k})$$







- Transition from this regime to Galactic turbulence generates a break in the diffusion coefficient just around  $\sim 200 \text{ GV}$ :



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**DIFFERENT BOUNDARY CONDITION:** 



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• Distribution of SNRs that CRs found in the Galactic disk as a function of their radius in the S-T phase

$$P(r_{\rm SN}) dr_{\rm SN} = K_P \frac{dt(r_{\rm SN})}{T_{\rm Max}}; \qquad T_{Max} \approx 3 \times 10^4 \,\mathrm{yr}$$

$$E_{\text{Max}}(t) \approx 100 \left(\frac{t}{t_{\text{ST}}}\right)^{-\frac{4}{5}} \text{TeV} = 100 \left(\frac{r_{\text{SN}}}{r_{\text{ST}}}\right)^{-2} \text{TeV}$$

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$$e^{-p/p_{Max}(r_{SN})}$$

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AVERAGE VOLUME

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$$AVERAGE VOLUME$$

Particles up to  $E_k \sim \text{TeV/n}$  are typically confined for a time  $T_{\text{SN}} \approx 3 \times 10^4$  yr inside the sources

A non-negligible production of secondaries might come from interactions occurring inside the SNR **BEFORE** the escape of primaries:

CONTRIBUTION FROM PRIMARY PARTICLES THAT ARE STILL LOCATED INSIDE THE SOURCES

$$Q_{src,\alpha} = v_{\alpha} r(s) n_{src,j} \frac{sV_{SN}T_{SN}\mathscr{R}}{\pi R_d^2} \times A_{\alpha'}K_{\alpha'} \int_{E_{th}}^{+\infty} dE'_{k,\alpha'} \frac{d\sigma_{\alpha',j}}{dE_{k,\alpha}} \left(\frac{p'(E'_{k,\alpha'})}{p_{inj,\alpha'}}\right)^{2-s}$$

**NB:** for spallation processes  $\frac{d\sigma_{\alpha',j}}{dE_{k,\alpha}} \equiv$ 

$$\frac{d\sigma_{\alpha',j}}{dE_{k,\alpha'}} \equiv \sigma_{\alpha,\alpha'}\,\delta(E'_{k,\alpha'} - E_{k,\alpha})$$

## **RESULTS**





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